WHOI Geophysical Fluids Program Summer, 2009

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Nonlinear waves

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Introduction to water waves Lecture 1





Q: What's so special about water waves?

Why study water waves?

- a) Woods Hole Oceanographic Institute
- b) For the Program on Nonlinear Waves:

Water waves provide a concrete physical example of a dynamical system rich enough to exhibit many of the mathematical concepts that have been developed in recent years:

linear stability, nonlinear stability

solitons, complete integrability

chaos, sensitive dependence on initial data

singularities, blow-up in finite time

deterministic vs. probabilistic models

 Water waves evolve on a "human" time-scale, so we can observe many of these concepts in physical experiments

Introduction to water waves





- Q: What are "water waves"?
- A (for my lectures): Waves in the water that you see or feel at the beach or in a boat

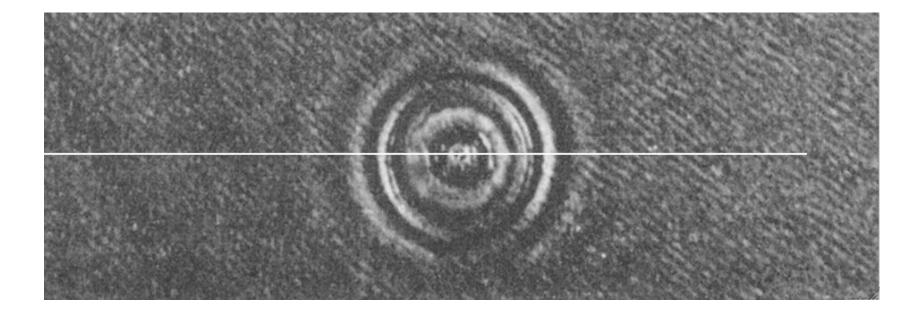
(sometimes called "surface water waves")

Properties of water waves

- Surface water waves have their maximum displacement at the free surface
- Waves propagate, with little dissipation
 - ask a baby in a bath-tub
 - Snodgrass *et al* (1966)
- Approximately periodic: 0.1 s < T < 25 s.
- Approximate maximum speed: $c = \sqrt{gh}$
- Water waves are "dispersive":

Long waves travel faster than short waves (for gravity-induced waves)

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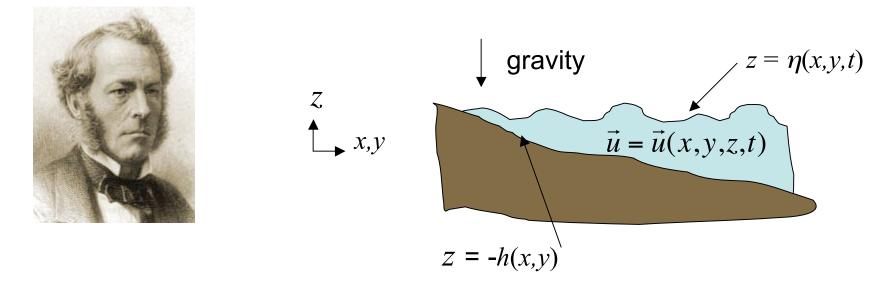
from Stoker's Water Waves (1957)

Ocean waves I am ignoring

- Sound waves (pressure waves) in water
 Speed of sound in water (at 10° C): 1450 m/sec
 Speed of 2004 tsunami: < 200 m/sec
 Pressure waves create initial conditions for surface waves
- Internal waves

Due to variations in fluid density Period of surface waves: seconds Period of internal waves: hours

Inertial waves (including Rossby waves)
 Due to rotation of earth
 Period of inertial waves ≥ 12 hours



Find: $\eta(x,y,t)$ – position of free surface, and $\vec{u}(x,y,z,t)$ – fluid velocity, for $-h(x,y) < z < \eta(x,y,t)$, for t > 0and for all $\{x,y\}$ with $\eta+h > 0$

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- [one more coming...]

x, y

- 2. Coordinate system(s)
- [*x*,*y*,*z*] denote fixed, laboratory coordinates
- Denote the current location of a fixed fluid particle by

 $\vec{x}(t) = \{x(t), y(t), z(t)\}$

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- $\left(\frac{D}{Dt}\right)$ means "following the fluid particle"
- acceleration of the fluid particle (in *z*-direction) is

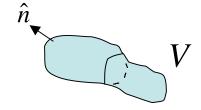
$$\frac{Dw}{Dt}(t,\vec{x}(t)) = \frac{\partial w}{\partial t} + \frac{\partial w}{\partial x} \cdot \frac{Dx}{Dt} + \frac{\partial w}{\partial y} \cdot \frac{Dy}{Dt} + \frac{\partial w}{\partial z} \cdot \frac{Dz}{Dt},$$
$$\frac{Dw}{Dt} = \frac{\partial w}{\partial t} + u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z}.$$

3. Physics

(a) Mass conservation of fluid with constant density

Total mass flow rate from V

$$= \oint_{\partial V} [(\rho)\vec{u} \cdot \hat{n}] ds = 0$$



3. Physics

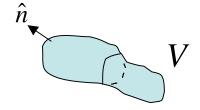
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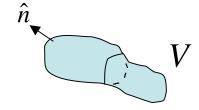
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Valid for all choices of V \twoheadrightarrow

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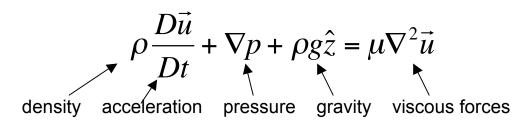
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Q: Is there any reason to believe $\vec{\omega} = 0$?

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(c) Back to reality

Navier-Stokes equations:



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$$\rho \frac{D\vec{u}}{Dt} + \nabla p + \rho g\hat{z} = \mu \nabla^2 \vec{u}$$

Neglect viscous forces \rightarrow Euler's equations

$$\frac{D\vec{u}}{Dt} + \frac{1}{\rho}\nabla p + g\hat{z} = 0$$

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If $\vec{\omega} \equiv 0$ at t = 0, then $\vec{\omega} \equiv 0$ for all t.

3. Physics

(d) Last assumption: irrotational flow

$$\vec{\omega} \equiv 0$$
 at $t = 0$.

$$\Rightarrow \quad \vec{u} = \nabla \phi, \quad \nabla^2 \phi = 0 \qquad \text{everywhere in fluid.}$$

This gives the simplest mathematical model, and the one studied the most. Dropping the assumption gives a more general model. (For internal waves, this is not valid.)

3. Physics

(e) Bernoulli's Law:

Show
$$\frac{D\vec{u}}{Dt} = \partial_t \vec{u} + \nabla(\frac{1}{2} |\vec{u}|^2) - \vec{u} \times \vec{\omega}$$

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$$\partial_t \phi + \frac{1}{2} |\nabla \phi|^2 + \frac{p}{\rho} + gz = F(t) = 0 \text{ in fluid}$$

3. Physics

- (f) What happens on the bottom boundary of fluid?
 - No flow through the (solid) bottom boundary:

→ on z = -h(x,y),

$$\vec{u} \cdot \nabla \{z + h(x, y)\} = 0$$
 or
 $\partial_z \phi + \nabla \phi \cdot \nabla h = 0$ or
 $\partial_n \phi = 0$

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(g) What happens at the free surface?

- With no surface tension and no wind

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$$\Rightarrow \qquad p = 0 - \sigma \left[\nabla \cdot (\frac{\nabla \eta}{\sqrt{1 + |\nabla \eta|^2}}) \right] \quad \text{on} \quad z = \eta(x, y, t).$$

3. Physics

(g) What happens at the free surface?

- 2 boundary conditions on $z = \eta(x,y,t)$:

* Kinematic: a particle on the surface remains there $\frac{D}{Dt} \{\eta(x(t), y(t), t) - z\} = 0$

$$\partial_t \eta + \nabla \phi \cdot \nabla \eta = \partial_z \phi$$

3. Physics

(g) What happens at the free surface?

- 2 boundary conditions on $z = \eta(x,y,t)$:

* Kinematic: a particle on the surface remains there

$$\frac{D}{Dt}\{z(t) - \eta(x(t), y(t), t)\} = 0$$

$$\Rightarrow \qquad \partial_t \eta + \nabla \phi \cdot \nabla \eta = \partial_z \phi$$

* Dynamic: Use surface pressure in Bernoulli's Law

$$\partial_{t}\phi + \frac{1}{2} |\nabla\phi|^{2} + g\eta = \frac{\sigma}{\rho} \nabla \cdot \{\frac{\nabla\eta}{\sqrt{1 + |\nabla\eta|^{2}}}\} \quad (+ \text{ wind ?})$$

Equations of motion

for an irrotational flow, with no forcing from wind:

$$\partial_t \eta + \nabla \phi \cdot \nabla \eta = \partial_z \phi,$$
 on $z = \eta(x, y, t)$

$$\partial_t \phi + \frac{1}{2} |\nabla \phi|^2 + g\eta = \frac{\sigma}{\rho} \nabla \cdot \{ \frac{\nabla \eta}{\sqrt{1 + |\nabla \eta|^2}} \}, \quad \text{on } z = \eta(x, y, t),$$

$$\nabla^2 \phi = 0 \qquad -h(x,y) < z < \eta(x,y,t),$$

$$\partial_z \phi + \nabla \phi \cdot \nabla h = 0,$$
 on $z = -h(x, y).$

Extra lecture -optional

Monday afternoon (when?):

Start with the nonlinear equations of water waves, linearize them, and explore some fundamental concepts that emerge from the linearized equations.