WHOI Geophysical Fluids Program Summer, 2009

富藏手六金

滚

いれいきの?

Nonlinear waves

Roger Grimshaw Loughborough U., UK Harvey Segur U. of Colorado, USA

Introduction to water waves Lecture 1

Q: What's so special about water waves?

Why study water waves?

- a) Woods Hole Oceanographic Institute
- b) For the Program on Nonlinear Waves:

Water waves provide a concrete physical example of a dynamical system rich enough to exhibit many of the mathematical concepts that have been developed in recent years:

linear stability, nonlinear stability

solitons, complete integrability

chaos, sensitive dependence on initial data

singularities, blow-up in finite time

deterministic vs. probabilistic models

c) Water waves evolve on a "human" time-scale, so we can observe many of these concepts in physical experiments

Introduction to water waves

Q: What are "water waves"?

A (for my lectures): Waves in the water that you see or feel at the beach or in a boat

(sometimes called "surface water waves")

Properties of water waves

- Surface water waves have their maximum displacement at the free surface
- Waves propagate, with little dissipation
	- ask a baby in a bath-tub
	- Snodgrass *et al* (1966)
- Approximately periodic: 0.1 s < *T* < 25 s.
- Approximate maximum speed: $c = \sqrt{gh}$
- Water waves are "dispersive":

Long waves travel faster than short waves (for gravity-induced waves) !

Long waves travel faster than short waves (for gravity-induced waves)

from Stoker's *Water Waves* (1957)

Ocean waves I am ignoring

- Sound waves (pressure waves) in water Speed of sound in water (at 10° C): 1450 m/sec Speed of 2004 tsunami: < 200 m/sec Pressure waves create initial conditions for surface waves
- Internal waves

Due to variations in fluid density Period of surface waves: seconds Period of internal waves: hours

• Inertial waves (including Rossby waves) Due to rotation of earth Period of inertial waves ≥ 12 hours

Find: $\eta(x, y, t)$ – position of free surface, and $\vec{u}(x,y,z,t)$ – fluid velocity, for $-h(x,y) < z < \eta(x,y,t)$, for $t > 0$ and for all $\{x,y\}$ with $\eta+h \geq 0$

- 1. Assumptions
- Water is incompressible $(•)$ no sound waves)

- 1. Assumptions
- Water is incompressible $(•\ \text{no sound waves})$
- Water has uniform density \leftrightarrow no internal waves)

- 1. Assumptions
- Water is incompressible $(•\ \text{no sound waves})$
- Water has uniform density \leftrightarrow no internal waves)
- Neglect rotation of earth \rightarrow no inertial waves)

- 1. Assumptions
- Water is incompressible $(•)$ no sound waves)
- Water has uniform density (4) no internal waves)
- Neglect rotation of earth \rightarrow no inertial waves)
- Gravity is a constant, downward force

- 1. Assumptions
- Water is incompressible $(•\ \text{no sound waves})$
- Water has uniform density (4) no internal waves)
- Neglect rotation of earth \rightarrow no inertial waves)
- Gravity is a constant, downward force
- Neglect effects of viscosity (sometimes wrong)

- 1. Assumptions
- Water is incompressible $(•\ \text{no sound waves})$
- Water has uniform density (4) no internal waves)
- Neglect rotation of earth \rightarrow no inertial waves)
- Gravity is a constant, downward force
- Neglect effects of viscosity (sometimes wrong)
- Effects of wind can be added or not

- 1. Assumptions
- Water is incompressible $(•\ \text{no sound waves})$
- Water has uniform density $($ \rightarrow no internal waves)
- Neglect rotation of earth \rightarrow no inertial waves)
- Gravity is a constant, downward force
- Neglect effects of viscosity (sometimes wrong)
- Effects of wind can be added or not
- [one more coming...]

z

x, *y*

- 2. Coordinate system(s)
- $[x, y, z]$ denote fixed, laboratory coordinates
- Denote the current location of a fixed fluid particle by

 $\ddot{\cdot}$ \rightarrow $\vec{x}(t) = \{x(t), y(t), z(t)\}$

z

x, *y*

- 2. Coordinate system(s)
- $[x, y, z]$ denote fixed, laboratory coordinates
- Denote the current location of a fixed fluid particle by

 $\ddot{\cdot}$ \rightarrow $\vec{x}(t) = \{x(t), y(t), z(t)\}$

– Velocity of this fluid particle is $-$ so $\left(\frac{D}{Dt}\right)$ means "following the fluid particle" *Dt* $\sqrt{ }$ \setminus $\left(\frac{D}{D_t}\right)$ $\overline{ }$ ' \rightarrow \vec{u} (\vec{x}, t = $(u,v,w) = \left(\frac{Dx}{D} \right)$ $\frac{Dx}{Dt}$, *Dy* $\frac{dy}{Dt}$, *Dz Dt* $\sqrt{ }$ \setminus $\left(\frac{Dx}{D} \cdot \frac{Dy}{D} \cdot \frac{Dz}{D} \right)$ \int

z

x, *y*

- 2. Coordinate system(s)
- [x,y,z] denote fixed, laboratory coordinates
- Denote the current location of a fixed fluid particle by

 $\ddot{\cdot}$ \rightarrow $\vec{x}(t) = \{x(t), y(t), z(t)\}$

- Velocity of this fluid particle is \rightarrow \vec{u} (\vec{x}, t = $(u,v,w) = \left(\frac{Dx}{D} \right)$ $\frac{Dx}{Dt}$, *Dy* $\frac{dy}{Dt}$, *Dz Dt* $\sqrt{ }$ \setminus $\left(\frac{Dx}{D} \cdot \frac{Dy}{D} \cdot \frac{Dz}{D} \right)$ \int
- $-\left(\frac{D}{Dt}\right)$ means "following the fluid particle" *Dt* $\sqrt{ }$ \setminus $\left(\frac{D}{D_t}\right)$ $\overline{ }$ '
- acceleration of the fluid particle (in *z*-direction) is !

$$
\frac{Dw}{Dt}(t,\vec{x}(t)) = \frac{\partial w}{\partial t} + \frac{\partial w}{\partial x} \cdot \frac{Dx}{Dt} + \frac{\partial w}{\partial y} \cdot \frac{Dy}{Dt} + \frac{\partial w}{\partial z} \cdot \frac{Dz}{Dt},
$$

$$
\frac{Dw}{Dt} = \frac{\partial w}{\partial t} + u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z}.
$$

3. Physics

(a) Mass conservation of fluid with constant density

Total mass flow rate from V

$$
= \oint_{\partial V} [(\rho) \vec{u} \cdot \hat{n}] ds = 0
$$

3. Physics

(a) Mass conservation of fluid with constant density

Total mass flow rate from V

$$
= \oint_{\partial V} [(\rho) \vec{u} \cdot \hat{n}] ds = 0
$$

Divergence theorem \rightarrow

$$
= \rho \iiint_V [\nabla \cdot \vec{u}] dv = 0
$$

3. Physics

(a) Mass conservation of fluid with constant density

Total mass flow rate from V

$$
= \oint_{\partial V} [(\rho) \vec{u} \cdot \hat{n}] ds = 0
$$

Divergence theorem \rightarrow

$$
= \rho \iiint_V [\nabla \cdot \vec{u}] dv = 0
$$

Valid for all choices of $V \rightarrow$

$$
\nabla \cdot \vec{u} = 0
$$

- 3. Physics
	- (b) Aside:

Define vorticity: $\vec{\omega}$ = curl(\vec{u})

- 3. Physics
	- (b) Aside:

Define vorticity: $\vec{\omega}$ = curl(\vec{u})

Suppose (for some reason) $\vec{\omega} = 0$

A there is a velocity potential, $\phi(\vec{x},t)$, with $\vec{u} = \nabla \phi$

3. Physics

(b) Aside:

Define vorticity: $\vec{\omega}$ = curl(\vec{u})

Suppose (for some reason) $\vec{\omega} = 0$

A there is a velocity potential, $\phi(\vec{x},t)$, with $\vec{u} = \nabla \phi$

 $\nabla \cdot \vec{u} = 0 \Leftrightarrow \nabla^2 \phi = 0$ Then

The velocity of the fluid is found by solving LaPlace' eq'n! What could be simpler?

3. Physics

(b) Aside:

Define vorticity: $\vec{\omega}$ = curl(\vec{u})

Suppose (for some reason) $\vec{\omega} = 0$

A there is a velocity potential, $\phi(\vec{x},t)$, with $\vec{u} = \nabla \phi$

 $\nabla \cdot \vec{u} = 0 \Leftrightarrow \nabla^2 \phi = 0$ Then

The velocity of the fluid is found by solving LaPlace' eq'n! What could be simpler?

Q: Is there any reason to believe $\vec{\omega} = 0$?

3. Physics

(c) Back to reality

Navier-Stokes equations:

3. Physics

(c) Back to reality

Navier-Stokes equations:

$$
\rho \frac{D\vec{u}}{Dt} + \nabla p + \rho g \hat{z} = \mu \hat{X}^2 \vec{u}
$$

Neglect viscous forces \rightarrow Euler's equations

$$
\frac{D\vec{u}}{Dt} + \frac{1}{\rho} \nabla p + g\hat{z} = 0
$$

3. Physics

(c) Back to reality

Navier-Stokes equations:

$$
\rho \frac{D\vec{u}}{Dt} + \nabla p + \rho g \hat{z} = \mu \nabla^2 \vec{u}
$$

Neglect viscous forces \rightarrow Euler's equations

$$
\frac{D\vec{u}}{Dt} + \frac{1}{\rho} \nabla p + g\hat{z} = 0
$$

Take curl
$$
\rightarrow \frac{D\vec{\omega}}{Dt} = \partial_t \vec{\omega} + \vec{u} \cdot \nabla \vec{\omega} = \vec{\omega} \cdot \nabla \vec{u}
$$

3. Physics

(c) Back to reality

Navier-Stokes equations:

$$
\rho \frac{D\vec{u}}{Dt} + \nabla p + \rho g \hat{z} = \mu \nabla^2 \vec{u}
$$

Neglect viscous forces \rightarrow Euler's equations

$$
\frac{D\vec{u}}{Dt} + \frac{1}{\rho} \nabla p + g\hat{z} = 0
$$

Take curl
$$
\rightarrow \frac{D\vec{\omega}}{Dt} = \partial_t \vec{\omega} + \vec{u} \cdot \nabla \vec{\omega} = \vec{\omega} \cdot \nabla \vec{u}
$$

If $\vec{\omega} \equiv 0$ at $t = 0$, then $\vec{\omega} \equiv 0$ for all t.

3. Physics

(d) Last assumption: irrotational flow

$$
\vec{\omega} \equiv 0 \qquad \text{at } t = 0.
$$

$$
\vec{u} = \nabla \phi, \quad \nabla^2 \phi = 0
$$
 everywhere in fluid.

This gives the simplest mathematical model, and the one studied the most. Dropping the assumption gives a more general model. (For internal waves, this is not valid.)

3. Physics

(e) Bernoulli's Law:

Show
$$
\frac{D\vec{u}}{Dt} = \partial_t \vec{u} + \nabla (\frac{1}{2} |\vec{u}|^2) - \vec{u} \times \vec{\omega}
$$

3. Physics

(e) Bernoulli's Law:

Show
$$
\frac{D\vec{u}}{Dt} = \partial_t \vec{u} + \nabla (\frac{1}{2} | \vec{u} |^2) - \vec{u} \times \vec{w}
$$

For an irrotational flow,

$$
\frac{D\vec{u}}{Dt} = \partial_t \vec{u} + \nabla(\frac{1}{2}|\vec{u}|^2) = \nabla\{\partial_t \phi + \frac{1}{2}|\nabla \phi|^2\}
$$

3. Physics

(e) Bernoulli's Law:

Show
$$
\frac{D\vec{u}}{Dt} = \partial_t \vec{u} + \nabla (\frac{1}{2} | \vec{u} |^2) - \vec{u} \times \vec{w}
$$

For an irrotational flow,

$$
\frac{D\vec{u}}{Dt} = \partial_t \vec{u} + \nabla (\frac{1}{2}|\vec{u}|^2) = \nabla {\partial_t \phi} + \frac{1}{2} |\nabla \phi|^2
$$

Euler \rightarrow ∇ { $\partial_t \phi + \frac{1}{2} |\nabla \phi|^2 + \frac{p}{2}$ ρ $+ gz$ } = 0

$$
\rightarrow \qquad \left| \partial_t \phi + \frac{1}{2} |\nabla \phi|^2 + \frac{p}{\rho} + gz = F(t) \right|
$$

3. Physics

(e) Bernoulli's Law:

Show
$$
\frac{D\vec{u}}{Dt} = \partial_t \vec{u} + \nabla (\frac{1}{2}|\vec{u}|^2) - \vec{u} \times \vec{w}
$$

For an irrotational flow,

$$
\frac{D\vec{u}}{Dt} = \partial_t \vec{u} + \nabla (\frac{1}{2} |\vec{u}|^2) = \nabla {\partial_t \phi + \frac{1}{2} |\nabla \phi|^2}
$$

Euler \rightarrow ∇ { $\partial_t \phi + \frac{1}{2} |\nabla \phi|^2 + \frac{p}{2}$ ρ $+ gz$ } = 0

$$
\theta_t \phi + \frac{1}{2} |\nabla \phi|^2 + \frac{p}{\rho} + gz = F(t) = 0 \text{ in fluid}
$$

- 3. Physics
	- (f) What happens on the bottom boundary of fluid?
		- No flow through the (solid) bottom boundary:

 \rightarrow on $z = -h(x, y)$,

$$
\vec{u} \cdot \nabla \{z + h(x, y)\} = 0 \quad \text{or}
$$

$$
\partial_z \phi + \nabla \phi \cdot \nabla h = 0 \quad \text{or}
$$

$$
\partial_n \phi = 0
$$

3. Physics

(g) What happens at the free surface?

– With no surface tension and no wind

$$
p = p_{air} (= 0) \qquad \text{on} \quad z = \eta(x, y, t)
$$

3. Physics

(g) What happens at the free surface?

– With no surface tension and no wind

$$
p = p_{air} (= 0) \qquad \text{on} \quad z = \eta(x, y, t)
$$

– With surface tension, usual model:

$$
p = p_{air} + \sigma(\frac{1}{R_1} + \frac{1}{R_2})
$$

const mean curvature

3. Physics

(g) What happens at the free surface?

– With no surface tension and no wind

$$
p = p_{air} (= 0) \qquad \text{on} \quad z = \eta(x, y, t)
$$

– With surface tension, usual model:

$$
p = p_{air} + \sigma(\frac{1}{R_1} + \frac{1}{R_2})
$$

const mean curvature

$$
p = 0 - \sigma \left[\nabla \cdot (\frac{\nabla \eta}{\sqrt{1 + |\nabla \eta|^2}}) \right] \text{ on } z = \eta(x, y, t).
$$

3. Physics

(g) What happens at the free surface?

 $-$ 2 boundary conditions on $z = \eta(x, y, t)$:

 * Kinematic: a particle on the surface remains there *D* $\frac{\partial}{\partial t} \{ \eta(x(t), y(t), t) - z \} = 0$

$$
\rightarrow
$$

$$
\overline{\partial_t \eta + \nabla \phi \cdot \nabla \eta = \partial_z \phi}
$$

3. Physics

(g) What happens at the free surface?

 $-$ 2 boundary conditions on $z = \eta(x, y, t)$:

* Kinematic: a particle on the surface remains there

$$
\frac{D}{Dt}\left\{z(t) - \eta(x(t), y(t), t)\right\} = 0
$$

$$
\rightarrow \qquad \boxed{\partial_t \eta + \nabla \phi \cdot \nabla \eta = \partial_z \phi}
$$

* Dynamic: Use surface pressure in Bernoulli's Law !

$$
\left[\partial_t \phi + \frac{1}{2} |\nabla \phi|^2 + g\eta = \frac{\sigma}{\rho} \nabla \cdot \left\{ \frac{\nabla \eta}{\sqrt{1 + |\nabla \eta|^2}} \right\} \right] \qquad \text{(+ wind?)}
$$

Equations of motion

for an irrotational flow, with no forcing from wind:

$$
\partial_t \eta + \nabla \phi \cdot \nabla \eta = \partial_z \phi, \qquad \text{on } z = \eta(x, y, t)
$$

$$
\partial_t \phi + \frac{1}{2} |\nabla \phi|^2 + g \eta = \frac{\sigma}{\rho} \nabla \cdot \{\frac{\nabla \eta}{\sqrt{1 + |\nabla \eta|^2}} \}, \qquad \text{on } z = \eta(x, y, t),
$$

$$
\nabla^2 \phi = 0 \qquad \qquad -h(x,y) < z < \eta(x,y,t),
$$

$$
\partial_z \phi + \nabla \phi \cdot \nabla h = 0, \qquad \text{on } z = -h(x, y).
$$

Extra lecture -optional

Monday afternoon (when?):

Start with the nonlinear equations of water waves, linearize them, and explore some fundamental concepts that emerge from the linearized equations.