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2 **Ocean Mesoscale and Frontal-scale Ocean-Atmosphere Interactions and**
3 **Influence on Large-scale Climate: A Review**
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46 **Abstract**

47 Two decades of high-resolution satellite observations and climate modeling studies have
48 indicated strong ocean-atmosphere coupled feedback mediated by ocean mesoscale processes,
49 including semi-permanent and meandering SST fronts, mesoscale eddies, and filaments. The air-
50 sea exchanges in latent heat, sensible heat, momentum, and carbon dioxide associated with this
51 so-called mesoscale air-sea interaction are robust near the major western boundary currents,
52 Southern Ocean fronts, and equatorial and coastal upwelling zones, but they are also ubiquitous
53 over the global oceans wherever ocean mesoscale processes are active. Current theories,
54 informed by rapidly advancing observational and modeling capabilities, have established the
55 importance of mesoscale and frontal-scale air-sea interaction processes for understanding large-
56 scale ocean circulation, biogeochemistry, and weather and climate variability. However,
57 numerous challenges remain to accurately diagnose, observe, and simulate mesoscale air-sea
58 interaction to quantify its impacts on large-scale processes. This article provides a
59 comprehensive review of key aspects pertinent to mesoscale air-sea interaction, synthesizes
60 current understanding with remaining gaps and uncertainties, and provides recommendations on
61 theoretical, observational, and modeling strategies for future air-sea interaction research.

62
63 **Significance Statement**

64 Recent high-resolution satellite observations and climate models have shown a significant impact
65 of coupled ocean-atmosphere interactions mediated by small-scale (mesoscale) ocean processes,
66 including ocean eddies and fronts, on Earth's climate. Ocean mesoscale-induced spatial
67 temperature and current variability modulate the air-sea exchanges in heat, momentum, and mass
68 (e.g., gases such as water vapor and carbon dioxide), altering coupled boundary layer processes.
69 Studies suggest that skillful simulations and predictions of ocean circulation, biogeochemistry,
70 and weather events and climate variability depend on accurate representation of the eddy-
71 mediated air-sea interaction. However, numerous challenges remain in accurately diagnosing,
72 observing, and simulating mesoscale air-sea interaction to quantify its large-scale impacts. This
73 article synthesizes the latest understanding of mesoscale air-sea interaction, identifies remaining
74 gaps and uncertainties, and provides recommendations on strategies for future ocean-weather-
75 climate research.

76 **1. Introduction**

77 Decades of observational and modeling analysis have broadly identified two fundamental
78 regimes of ocean-atmosphere coupling dependent on the spatial scale of ocean surface
79 variability. The first regime involves the ocean response to large-scale (>1000 km) internal
80 atmospheric variability, which drives a response in sea surface temperature (SST) through the
81 mediation of surface turbulent heat fluxes and upper-ocean turbulent mixing (e.g., Frankignoul et
82 al. 1985; Alexander and Scott 1997). The large-scale ocean response feeds back onto the
83 incipient atmospheric circulation anomaly to reinforce or erode it (e.g., Bladé 1997). In this
84 framework, the ocean is viewed as relatively passive, mainly advecting anomalies, storing heat,
85 and integrating white noise atmospheric forcing.

86
87 The second regime, the focus of this paper, involves an atmospheric response driven by ocean
88 mesoscale eddy-induced spatial SST and current variability. Here, the term “mesoscale eddies
89 and fronts” broadly refers to all forms of oceanic processes with horizontal length-scales smaller
90 than the first regime of air-sea interaction (>1000 km) but larger than oceanic submesoscale (~1-
91 10 km), although several outstanding issues regarding the submesoscale air-sea interactions will
92 be discussed in Sections 5 and 6. These processes include coherent, swirling, and transient ocean
93 circulations with length-scales near the Rossby radius of deformation (Chelton et al. 2011),
94 filamentary eddy structures that are widely observed in coastal upwelling systems, and semi-
95 permanent fronts and undulations near the midlatitude western boundary currents (WBCs) and
96 their extensions, and SST fronts along the equatorial tongue in the Pacific and Atlantic oceans.

97
98 The SST signature from these ocean mesoscale processes modifies surface turbulent heat and
99 momentum fluxes, driving local responses in marine atmospheric boundary layer (MABL)
100 processes (Small et al. 2008), inducing responses in winds, clouds, and rainfall (e.g., Deser et al.
101 1993; Tokinaga et al. 2009; Frenger et al. 2013; Miyamoto et al. 2018; 2022; Takahashi et al.
102 2020, 2021). The MABL responses then drive non-local responses in the path and activity of
103 storm tracks in the extratropics (e.g., Czaja et al. 2019) and deep moist convection in the tropics
104 (e.g., Li and Carbone 2012; Skillingstad et al. 2019; de Szoeke and Maloney 2020). The
105 atmospheric response to ocean mesoscales feeds back onto eddy activity and SST, altering the
106 large-scale ocean circulation, further influencing these atmospheric processes (e.g., Nakamura et

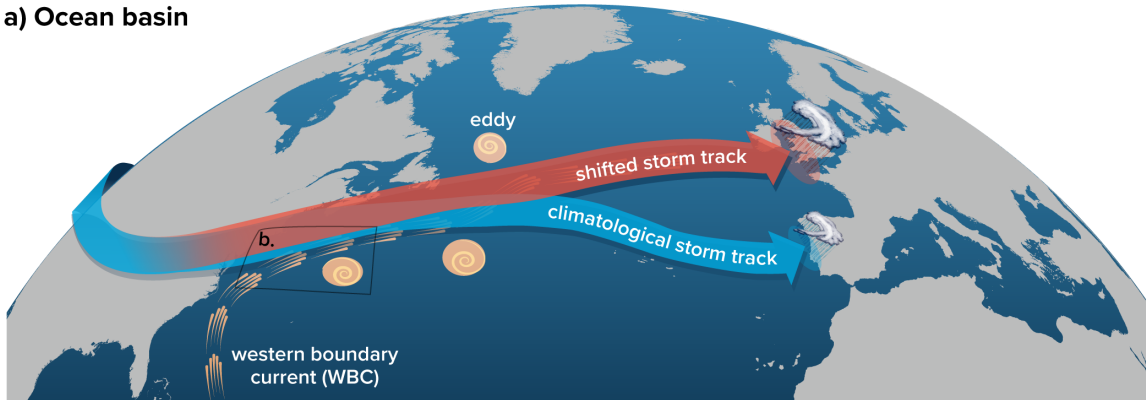
107 al. 2008; Hogg et al. 2009; Frankignoul et al. 2011; Taguchi et al. 2012). Mesoscale ocean
108 surface currents also affect the wind stress and heat fluxes as well as the kinematic profiles in the
109 MABL, which influence ocean circulation, including the stability and strength of the WBCs and
110 their meanders (Renault et al. 2016b, 2019b) and the basin-scale coupled climate variability such
111 as ENSO (e.g., Luo et al. 2005). The ocean drives the SST variability more strongly than the
112 atmosphere at longer time-scales and shorter spatial-scales (Bishop et al. 2017), suggesting the
113 need to include rectified coupled effects of ocean mesoscale eddies in high-resolution coupled
114 climate models (Bryan et al. 2010; Kirtman et al. 2012; Roberts et al. 2016; Hewitt et al. 2020).

115
116 Aside from earlier limited observational studies showing evidence of the MABL response to
117 mesoscale SSTs (e.g., Sweet et al. 1981), the first observational global-scale surveys of the
118 MABL and surface wind responses based on satellite observations were provided by Chelton et
119 al. (2004) and Xie (2004), followed by comprehensive review papers by Small et al. (2008) and
120 Kelly et al. (2010). The number of publications that include aspects of mesoscale air-sea
121 interaction has grown exponentially in the last decade or so (see Robinson et al. 2018, 2020),
122 which also emphasizes a strong cross-disciplinary nature of the research subject (e.g., AMS
123 Special Collection on [Climate Implications of Frontal Scale Air-Sea Interaction](#), and the *J.*
124 *Oceanography* Special Collection on “Hot Spots” in the climate system, Nakamura et al. 2015).
125 Notwithstanding the existing review papers, no comprehensive synthesis papers exist that
126 consolidate the exponential increase in scientific understanding of mesoscale air-sea interaction.
127 This forms the key motivation of this review, which mainly focuses on a synthesis of the studies
128 since Small et al. (2008).

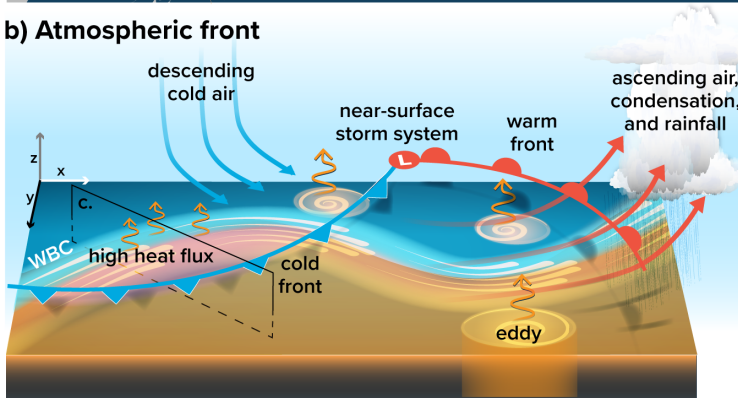
129
130 The paper is organized in the following logical order. Section 2 discusses the air-sea flux
131 responses to mesoscale SST and surface currents, along with theories and analytical studies of
132 MABL dynamics describing the flux responses. The subsequent two sections review critical
133 aspects of large-scale atmospheric and ocean circulation responses resulting from the
134 atmospheric boundary layer processes. That is, Section 3 discusses the tropospheric responses
135 emphasizing the modulation of local and downstream adjustments of extratropical weather
136 systems and their aspects related to climate change. Section 4 probes into the oceanic responses
137 due to thermal and mechanical feedback processes. The chapter emphasizes the need to develop

138 new theories and parameterizations to account for rectified effects of eddy-atmosphere
 139 interaction. Section 5 explores the emerging observational platforms critical for accurate in situ
 140 and remote-sensing characterization of air-sea interaction at small spatial scales in the coming
 141 decade. Section 6 provides a summary and synthesis.
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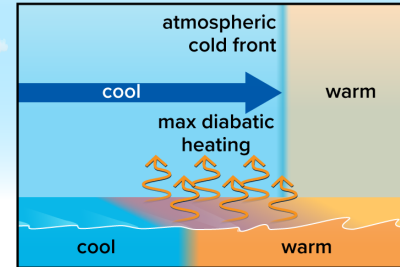
a) Ocean basin



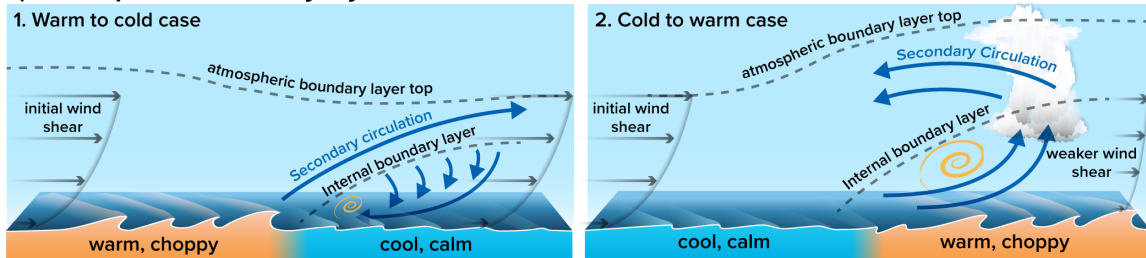
b) Atmospheric front



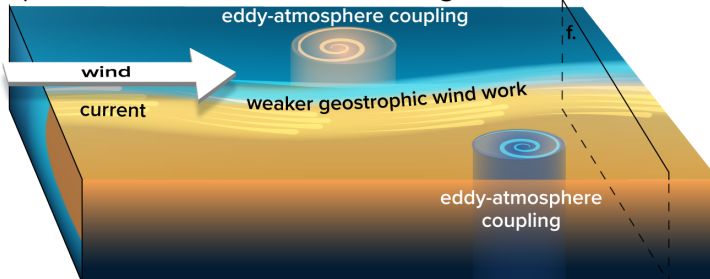
c) Atmospheric front cross-section



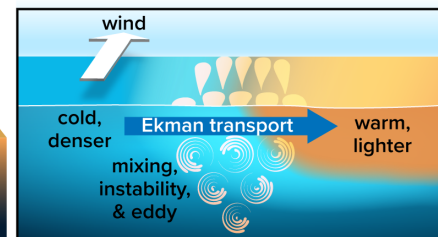
d) Atmospheric boundary layer



e) Ocean fronts & eddies interacting with wind



f) Stratification, instability & turbulence at fronts



143

144 **Figure 1:** Schematic illustrations of the coupled ocean-atmosphere feedback processes in the Northern
145 Hemisphere. (a) On the basin scale, the storm track affected by the WBCs leads to anomalous rainfall patterns
146 downstream. (b) A zoom-in view over the black box in (a) illustrates cold and warm fronts within a low-
147 pressure system traversing the semi-permanent SST front. On the trailing edge of the cold front (purple), the
148 cold/dry air mass over the warm ocean water induces large diabatic heating of the storms, strengthening the
149 storm. A similar process might occur over the transient mesoscale eddies. The modified air mass ascends over
150 the warm front, leading to deep cumulus clouds and heavy precipitation. (c) A 2-D view of the cross-section in
151 (b), where the cold front translates eastward over the SST front. When the cold front is east of the SST front,
152 the large air-sea temperature and humidity differences (purple) cause the maximum upward turbulent heat flux,
153 facilitating the diabatic frontogenesis. (d) A 2-D view of the MABL with the cross-frontal winds. For the
154 warm-to-cold case, the warm air blowing over cold water downwind of the SST front leads to a stable internal
155 boundary layer with a capping inversion and a shallow clockwise secondary circulation. Due to weaker vertical
156 mixing, the surface wind slows down, reinforcing the initial wind shear. The weak wind over cold SST yields a
157 reduced surface drag. For the cold-to-warm case, MABL and internal boundary layers deepen quickly, with the
158 counter-clockwise secondary circulation developing downstream. The increased turbulent mixing accelerates
159 the surface wind, leading to a well-mixed wind profile. The choppy surface waves on the warm side due to
160 higher winds enhance surface drag. Wind direction also changes across the front as wind speed adjusts to local
161 stability (not featured in this schematic). The surface currents near the ocean front (also not shown) modulate
162 the wave slopes and surface roughness via wave-current interaction and the wind stress via current-wind
163 interaction. (e) Meandering eastward currents and mesoscale eddies under a uniform westerly wind. On a large
164 scale, because surface currents are oriented downwind, the relative wind leads to weaker geostrophic wind
165 work than the absolute wind, stabilizing the large-scale circulation but stimulating submesoscale instabilities.
166 Over the eddies, eddy-atmosphere coupling induces the diabatic dissipation of eddy potential energy (thermal
167 feedback) and the negative geostrophic eddy wind work via current-wind interaction (mechanical feedback),
168 weakening the eddy energy. The eddies' swirling currents manifest reversely in the wind stress, leading to
169 current-induced wind stress curls and the up/downwelling in the ocean. (f) The cross-section across the
170 front/jet in (e). The down-front wind drives an eastward Ekman transport of cold/dense water over warm/light
171 water, reducing stratification near the front. The unstable front leads to enhanced turbulence and submesoscale
172 activity, with the induced secondary circulation accelerating the jet. The oceanic frontogenesis influenced by
173 the surface waves is not featured in this schematic but illustrated in Figure 9.

174
175 The readers might find it helpful to visualize key feedback mechanisms discussed throughout the
176 paper by referring to the schematic illustrations in Figure 1, which are organized at different
177 characteristic length scales and by processes. The MABL response to a mesoscale SST front
178 (Figure 1d) corresponds to Section 2. The diabatic heat exchanges between the atmospheric
179 fronts and the SST fronts (Figure 1b-c) are elaborated in Section 3b, while a broader view of
180 modulation of the midlatitude storm track by the WBCs and the subsequent downstream rainfall
181 patterns (Figure 1a) is discussed in detail in Sections 3a-c. The discussion about the modulation
182 of wind stress and heat fluxes by the mean and eddy currents and their feedback to oceans
183 (Figure 1e) jibes with Section 4a. The resulting fine-scale near-surface instability and turbulence
184 (Figure 1d) are touched upon in Section 4b-c.

185
186 It is not possible to cover all relevant aspects of mesoscale air-sea interaction with sufficient
187 details. There exist many review articles that might be helpful for readers interested in gaining a

188 more in-depth understanding of specific topics. For Section 2, such papers include Bourassa et
189 al. (2013) on challenges/needs for accurate air-sea flux measurements in high-latitude oceans;
190 Swart et al. (2019) on observational strategies to improve Southern Ocean heat and gas flux
191 estimates; Cronin et al. (2019) on global air-sea flux accuracy requirements; Bourassa et al.
192 (2019) on satellite remote sensing of wind and winds stress; and Deskos et al. (2021) on sea state
193 impacts on surface winds from a wind energy perspective. For Section 3, Kushnir et al. (2002)
194 reviewed the atmospheric responses to extratropical SST anomalies in climate models. Czaja et
195 al. (2019) updated the extratropical air-sea interaction based on high-resolution climate modeling
196 studies, while Kwon et al. (2010) and Kelly et al. (2010) reviewed the impacts of WBC SST
197 anomalies on seasonal to decadal climate variability. For Section 4, more detailed accounts of
198 surface waves, upper ocean mixing, and submesoscale dynamics are provided by Sullivan and
199 McWilliams (2010), D’Asaro (2014), and McWilliams (2016). McGillicuddy et al. (2016) offer a
200 comprehensive review of mechanisms of physical-biological-biogeochemical interactions on the
201 oceanic mesoscale. For Section 5, helpful review papers include Arduin et al. (2019) on
202 observing sea state information; Villas Bôas et al. (2019) on wind-wave-current interaction;
203 Centurioni et al. (2019) on global ocean surface observation networks; and Wanninkhof et al.
204 (2019) on global CO₂ flux measurements. The observational needs for data assimilation, coupled
205 reanalyses, and short-term and extended-range predictions have been discussed by Penny et al.
206 (2019), Domingues et al. (2019), and Subramanian et al. (2019).

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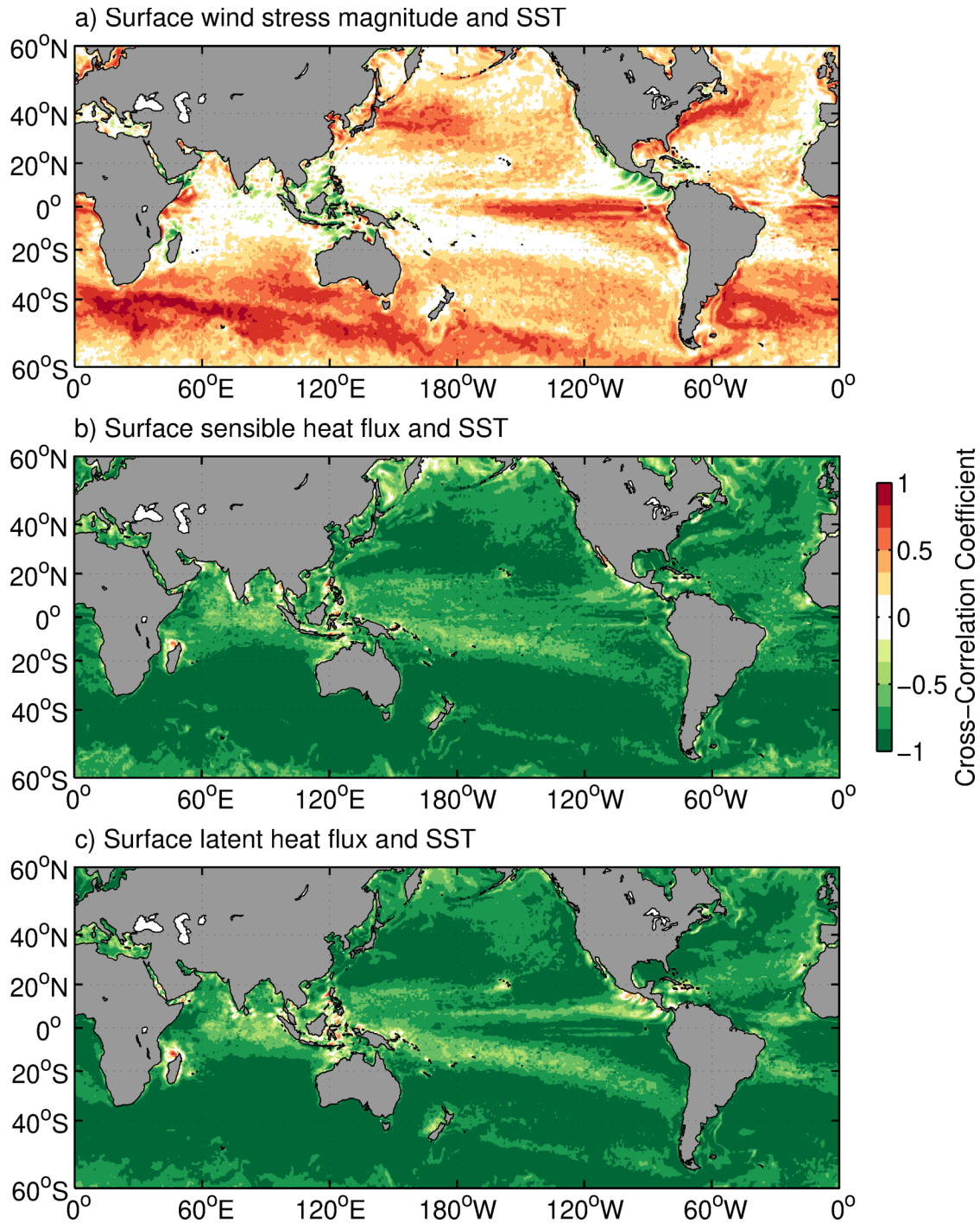
209 **2. Boundary layer and surface heat, momentum, and gas flux responses**

210 Surface fluxes communicate mass and energy between the ocean and atmosphere and are thus
211 vital processes in Earth’s climate system. The ocean is a major reservoir of heat and carbon in
212 the Earth system, and it is increasingly clear that exchanges with the atmosphere occurring on the
213 oceanic mesoscale are significant in shaping Earth’s climate. Recent assessments on projected
214 trends in surface air temperature (SAT) and SST have indicated a need to better understand
215 surface heat fluxes to reconcile conflicting lines of evidence on the projected trends in SAT and
216 SST (e.g., Box TS.1, IPCC 2021). The surface turbulent heat fluxes are composed of sensible
217 and latent heat fluxes, while the surface wind stress represents the turbulent momentum flux
218 between the atmosphere and ocean mediated by surface waves. This section discusses air-sea

219 heat, momentum, and gas flux responses to spatially heterogeneous fields of SST, surface
220 currents, and sea state. We also discuss the local MABL response to ocean-induced mesoscale
221 forcing, given its strong relationship with the surface fluxes. These processes are illustrated in
222 Figure 1d.

223
224 Spatially heterogeneous SST and surface currents generate localized anomalies in the surface
225 heat and momentum fluxes. The atmospheric and oceanic responses to these flux anomalies are
226 initially confined to the MABL and ocean mixed layer, but the responses to this coupling may
227 spread to the free atmosphere above (Section 3) or the ocean thermocline below (Section 4). The
228 atmospheric boundary layer and the oceanic mixed layer directly mediate responses of the large-
229 scale oceanic and atmospheric circulation to the mesoscale and frontal-scale air-sea coupling.

230
231 Figure 2 shows the strong correlation between monthly mesoscale surface fluxes and ocean
232 mesoscale variability from the ERA5 reanalysis (Hersbach et al. 2020). Here, the turbulent heat
233 flux is defined as positive downward (ocean warming). When the local point-by-point correlation
234 between the turbulent fluxes and SST is strongly negative, the SST variability can be viewed as
235 the ocean forcing the atmosphere (e.g., the warm ocean heats the atmosphere). Similarly, when
236 the correlation between turbulent heat flux and SST *tendency* is positive, the atmosphere is
237 considered to drive ocean variability. Over mesoscale, the wind stress and upward heat fluxes are
238 enhanced over warm SST anomalies (SSTA) and reduced over cool SSTA. The correlations are
239 much stronger for sensible and latent heat flux responses, while the surface stress response on
240 this spatial scale is much more apparent in oceanic frontal boundary regions where mesoscale
241 SST variability is most pronounced. However, it should be noted that the amplitude of
242 correlation represents empirical estimates of the strength of covariability since the atmospheric
243 response to an ocean anomaly modifies the turbulent fluxes and would obscure this simple rule
244 (e.g., Sutton and Mathieu, 2002). The effect of the surface flux on the ocean is discussed in
245 Section 4.



246

247 **Figure 2:** Maps of the cross-correlation coefficients between ERA5 monthly spatially high-pass filtered SST
 248 and (a) wind stress magnitude, (b) surface sensible heat flux, and (c) surface latent heat flux. The spatial high-
 249 pass filter removed variability with spatial scales greater than 1000 km. These maps were averaged over the
 250 30-year period 1991-2020. The ERA5 reanalysis time period used here was 1991-2020. The standard sign
 251 convention for ERA5 surface fluxes is used: positive fluxes mean energy entering the ocean. The high
 252 correlations in these maps correspond to regions of strong mesoscale SST variability, such as in the WBCs and
 253 their extension regions (Kuroshio, Gulf Stream, Brazil Current, and Agulhas Current), along the Antarctic
 254 Circumpolar Current and equatorial fronts, and near the Somali Current. A similar plot to Figure 2(a) can be
 255 found in Small et al. (2008) and Seo (2017).

256 *a. Turbulent heat flux response*

257 On smaller scales encompassed by the oceanic mesoscale and on time-scales longer than
258 synoptic time-scales in the atmosphere (e.g., 2-8 days), spatial variations in the surface turbulent
259 heat fluxes are driven primarily by spatial perturbations of SST, such that negative heat flux
260 anomalies (i.e., atmosphere heat gain) occur over warm SST perturbations and positive heat flux
261 anomalies (i.e., atmosphere heat loss) occur over cool SST perturbations (Figure 2b-c). Over
262 these scales, the ocean forces a response of the atmosphere driven by the surface heat exchange,
263 which is fundamentally distinct from the response over larger spatial scales. Near-surface air
264 temperature and specific humidity adjust slowly to spatially heterogeneous SST as air flows
265 across SST gradients. Ocean mesoscale eddies and SST fronts near the semi-permanent WBCs
266 often generate large air-sea temperature and humidity differences (Figure 1b-c). A dramatic
267 example was observed during the CLIMODE experiment near the Gulf Stream during
268 wintertime, when air-sea temperature differences exceeded 10 °C over 200 km, yielding >1000
269 W/m² surface turbulent heat fluxes into the atmosphere (Marshall et al. 2009).

270
271 Past field experiments captured less extreme but nonetheless strong responses of turbulent heat
272 fluxes and MABL convective turbulence to mesoscale and frontal-scale SSTs. Examples can be
273 found from the Sargasso Sea during the FASINEX experiment (e.g., Friehe et al. 1991), Gulf
274 Stream (e.g., Plagge et al. 2016), Kuroshio (e.g., Tokinaga et al. 2009); Pacific Tropical
275 Instability Waves (Thum et al. 2002), Brazil-Malvinas Confluence system (e.g., Pezzi et al.
276 2005; Villas Bôas et al. 2015; Souza et al. 2021; Cabrera et al. 2022), Agulhas Current (e.g., Jury
277 and Courtney 1991; Messenger and Swart 2016), and the western Arabian Sea (e.g., Vecchi et al.
278 2004).

279
280 The scale dependence of turbulent flux responses to mesoscale SST variations has been
281 quantified primarily from reanalysis-based surface flux and SST datasets (e.g., Li et al. 2017;
282 Sun and Wu 2022). Bishop et al. (2017), in particular, showed that on time-scales longer than
283 one month, the turbulent heat fluxes on the ocean mesoscale and frontal scale are driven by SST
284 variability associated with oceanic internal processes. On shorter time-scales, the variability is
285 driven more by synoptic-scale weather variability, particularly along the storm tracks overlying
286 the WBCs. Based on this simple diagnostic, Kirtman et al. (2012) concluded that eddy-

287 parameterized models grossly underestimate the ocean forcing of the atmosphere in eddy-rich
288 regions (e.g., WBCs and the Southern Ocean) and overestimate the atmospheric forcing of the
289 ocean throughout much of the mid-latitudes compared to the ocean eddy-resolving simulations.

290

291 *b. Turbulent momentum flux and MABL wind responses*

292 The turbulent heat flux response to SST is a crucial process that drives the responses in turbulent
293 momentum flux to SST. The variability in ocean surface currents at mesoscales also affects the
294 wind stress through the relative motion of the surface winds and currents. The most immediate
295 local atmospheric response to SST and surface currents is initially confined to the MABL. The
296 wind and wind stress responses mainly result from a dynamical adjustment of the MABL
297 pressure and vertical turbulent stress profile distinct from simple adjustments of the surface layer
298 logarithmic wind profile (Small et al. 2008; O'Neill 2012; Renault et al. 2016a), the relative
299 importance of which strongly depends upon background wind condition (e.g, Schneider and Qiu
300 2015; Byrne et al. 2015; Section 2c).

301

302 1) Mesoscale SST effects

303 Traditionally, local atmospheric responses to the mesoscale SST have been characterized
304 empirically by linear regressions between collocated mesoscale SSTs and surface winds and
305 surface wind stress, all spatially high-pass filtered to isolate the coupling on scales smaller than
306 about $O(1000 \text{ km})$. Linear regression coefficients, also called coupling coefficients, obtained
307 from satellite-observed wind speed and wind stress indicates ubiquitous increases in their
308 magnitudes over warm SSTs, increases of wind divergence and wind stress divergence co-
309 located with the downwind component of the SST gradient, and wind curl and wind stress curl
310 that scale with crosswind components of SST gradients (Chelton et al. 2001; O'Neill et al. 2003,
311 2012). The SST-induced curl and divergence responses provide further constraints on spatial
312 scales of the SST-induced MABL response. These simple but powerful diagnostic metrics have
313 been broadly used to diagnose the simulated air-sea interaction over a range of scales in
314 numerical models (Bellucci et al. 2021), leading to refinements in the SST resolution (Chelton
315 2005) and the PBL parameterizations in NWP models (Song et al. 2017). However, the coupling
316 coefficients include contributions from broad scales represented in the high-pass filtered input
317 fields. Hence, other than the gross separation of small scales from large scales, it is difficult to

318 extract useful information about scale dependence from such calculations. Alternative statistical
319 and analytical approaches exist, including cross-spectral analysis (e.g., Small et al. 2005b;
320 O'Neill et al. 2012; Laurindo et al. 2019; Samelson et al. 2020), cross-covariance and correlation
321 functions between SST (and its tendency), wind and turbulent heat fluxes (e.g., Frankignoul and
322 Hasselmann 1977; Wu et al. 2006; Bishop et al. 2017; Small et al. 2019), and an analytical
323 model for MABL heat and momentum budgets (Schneider and Qiu 2015; Schneider 2020). The
324 analytical model for MABL is explored in detail in Section 2c.

325

326 2) Mesoscale current effects

327 Regions of strong SST gradients are also regions of substantial variability in ocean surface
328 current. The current feedback (CFB) mechanism directly modifies wind stress through the
329 relative motion of surface winds and currents, which in turn alters the low-level wind shear and
330 wind. That is, a negative current anomaly induces a positive stress anomaly acting on the
331 atmosphere, which causes a negative wind anomaly (Renault et al. 2016a). At the mesoscale,
332 CFB primarily impacts the surface wind stress curl but not its divergence due to the quasi-
333 geostrophic nature of ocean currents (Chelton et al. 2004). The wind stress and wind responses to
334 CFB can also be diagnosed using empirical relationships based on satellite and numerical
335 simulations. Renault et al. (2016a; 2019a) defined two coupling coefficients related to CFB: s_w is
336 the regression slope between mesoscale surface currents and 10 m wind and s_τ is the linear
337 regression coefficient linking mesoscale surface current and surface stress. The coefficient s_τ can
338 be interpreted as a measure of the damping efficiency of CFB to ocean eddy energy, as discussed
339 in greater detail in Section 4.

340

341 The SST and current-induced stress responses are challenging to separate since mesoscale SST
342 and current variations co-vary strongly near ocean fronts and eddies. Nonetheless, estimates of
343 the contributions of the current-induced wind stress response via the linear coupling coefficients
344 indicate that the current-induced stress anomalies exceed the SST-induced response over strong
345 WBCs and within isolated ocean eddies (e.g., Gaube et al. 2015; Renault et al. 2019a). The
346 current-induced stress response exists in scatterometer and direct air-sea flux observations and
347 coupled ocean-atmosphere simulations, but it is not directly apparent in atmosphere-only
348 simulations and reanalyses, such as the ERA5 wind stress anomalies used in Figure 2. Including

349 both current and SST-induced stress anomalies strongly impacts the mesoscale wind stress curl
350 field (e.g., Renault et al. 2019a).

351

352 *c. Analytic framework for SST-induced boundary layer response*

353 The MABL response to ocean mesoscale current must incorporate coupling between the MABL
354 thermodynamics and dynamics to adequately represent the influence of SST and surface current
355 on the surface wind stress and sensible and latent heat fluxes. An analytical framework for SST
356 impacts was recently proposed, which incorporates MABL heat and momentum budgets that
357 capture the first-order response of the MABL to SST forcing (Schneider and Qiu 2015;
358 Schneider 2020) and includes a representation of the processes shown in the literature to be of
359 primary importance. This framework considers an MABL capped by an inversion (Battisti et al.
360 1999). Within this layer, air temperature is assumed to be well mixed and vertically constant, and
361 subject to horizontal advection and air-sea heat exchanges. The system is driven by winds with
362 horizontal scales far larger than the ocean mesoscale that satisfy a drag law at the sea surface and
363 experience zero vertical momentum flux at the inversion. The large-scale winds \vec{U} form a
364 modified Ekman spiral (Holton 1965a,b), which is considered horizontally homogeneous on
365 scales commensurate with the ocean mesoscale.

366

367 SST T enters the heat budget of the layer via the air-sea heat exchanges due to the air-sea
368 temperature difference with a rate γ . The MABL air temperature Θ results, to first order, from a
369 quasi-steady balance of surface sensible heat fluxes with advection by large-scale winds (e.g.,
370 Small et al. 2005a),

371

$$372 \quad \vec{U} \cdot \nabla \Theta = \gamma(T - \Theta). \quad (1)$$

373

374 The air temperatures Θ adjust to SST T over a length-scale of U/γ , forming a wake of elevated
375 values of the air-sea temperature differences in the lee of spatial SST variations. Thermal
376 adjustment rates of the boundary layer γ correspond to adjustment times of a few hours to half a
377 day (Schubert et al. 1979), yielding length-scales of the response of $O(100 \text{ km})$. The momentum
378 equations govern the wind response to the ocean mesoscale SST-induced acceleration \vec{F}
379 (Schneider and Qiu 2015) such that,

380
$$\underbrace{\vec{U} \cdot \nabla \vec{u}}_I + \underbrace{\frac{w^*}{H} \partial_s \vec{U}}_{II} + \underbrace{f \hat{e}_3 \times \vec{u}}_{III} - \underbrace{\frac{1}{H^2} \partial_s A \partial_s \vec{u}}_{IV} + \underbrace{g' \nabla h}_V = \vec{F}. \quad (2)$$

381
 382 On the left-hand side of Eq. (2), the term I represents the horizontal advection by large-scale
 383 winds \vec{U} of SST-induced winds \vec{u} . The term II is the vertical advection by w^* of the large-scale
 384 shear, where a sigma coordinate is used in the vertical so that $s=0$ is the sea surface and $s=1$ is
 385 the mean inversion height, H . The term III indicates the Coriolis acceleration with Coriolis
 386 frequency f , where \hat{e}_3 denotes the unit vector in the vertical. The term IV is the divergence of
 387 vertical fluxes of horizontal momentum due to large-scale mixing with eddy coefficient A . The
 388 term V is the hydrostatic pressure gradient forces, including the so-called back pressure effect
 389 (e.g., Hashizume et al. 2002), due to ocean mesoscale-induced changes of inversion height (h).
 390 Together with the continuity equation and boundary conditions of a drag law at the sea surface,
 391 and a material inversion with no flux of momentum, these equations provide a complete
 392 analytical solution for the wind response to ocean mesoscale SSTs.

393
 394 The changes in θ due to ocean mesoscale SSTs impact acceleration \vec{F} to the horizontal
 395 momentum equation

396
 397
$$\vec{F} = \underbrace{\frac{gH}{\theta_0} (1-s) \nabla \theta}_{VI} + \underbrace{\frac{1}{H^2} \partial_s (\dot{A} \partial_s \vec{U})}_{VII} \quad (3)$$

398 through the modulation of the hydrostatic pressure gradients (the term VI) and the sensitivity of
 399 the vertical mixing to the fluxes at the air-sea interface (the term VII). Here, θ_0 is a reference
 400 temperature, g is the earth's gravitational acceleration, and \dot{A} is the sensitivity of vertical mixing
 401 coefficient A to SST.

402
 403 The pressure effect (the term VI), originally formulated by [Lindzen and Nigam \(1987\)](#),
 404 designates the acceleration of surface winds to the baroclinic pressure gradient imparted by air
 405 temperature gradients, which drive secondary wind circulations and up/down-drafts (e.g., Wai
 406 and Stage 1989; Wenegrat and Arthur 2018; Sullivan et al. 2020; Figure 1d). [Lindzen and Nigam](#)
 407 (1987) neglected advection and assumed that air temperature decays linearly from the SST to
 408 zero at a height of 3000 m. In contrast, we include advection in the momentum budget in Eq. (2)

409 and assume that the SST imprint is vertically constant, consistent with a reduced gravity
410 formulation (Battisti et al. 1999).

411
412 The vertical mixing effect (the term VII) is a linearization of the 'nonlinear' term envisioned by
413 Wallace et al. (1989) and Hayes et al. (1989) that captures the modulations of the vertical mixing
414 acting on the large-scale wind profile. The dynamics, amplitude, and vertical structure of
415 \dot{A} determine the character of mixing sensitivity. Mixing can intensify and change its vertical
416 scale. The dependence of vertical mixing on the non-equilibrium air-sea temperature difference
417 is but one possibility. Alternatively, SST induces convective adjustment of the lapse rate and
418 permanently deepens the atmospheric boundary layer over warmer waters (Samelson et al.
419 2006). These diagnostic formulations for \dot{A} are endpoints of the non-equilibrium evolution of
420 vertical mixing simulated by Large Eddy Simulations (LES, e.g., de Szoeke and Bretherton
421 2004; Skyllingstad et al. 2007; Sullivan et al. 2020), which allow for changes in the vertical
422 mixing that lag modulations of boundary layer stability (Wenegrat and Arthur 2018). As such,
423 the coupling between surface winds and SST is sensitive to the MABL turbulence closure
424 schemes (e.g., Song et al. 2009, 2017; Perlin et al. 2014; Samelson et al. 2020). The MABL
425 turbulence subsequently affects the SST by altering mixing and entrainment in the ocean surface
426 boundary layer, indicating co-dependence of the turbulent boundary layer schemes in the
427 atmosphere and oceans (Fox-Kemper et al. 2022).

428
429 Advection by large-scale winds allows for disequilibrium in air-sea temperature and shifts
430 responses of winds or stress as a function of the SST spatial scales and the large-scale wind
431 direction and speed (e.g., Small et al. 2005a, 2008). Spectral transfer functions, or their
432 corresponding physical-space impulse response functions, capture these non-local relationships
433 and generalize the widely used coupling coefficients to include spatial lags. Estimates from
434 satellite observed winds and SST of spectral transfer functions suggest scale-dependent, lagged
435 dynamics as a function of the Rossby number determined by large-scale winds, the wavenumbers
436 of ocean mesoscale SST, and the Coriolis frequency f , or thermal or frictional adjustment rates γ
437 or A/H^2 (Schneider 2020; Masunaga and Schneider 2022). For small Rossby numbers, the
438 pressure effect dominates, while large Rossby numbers favor the vertical mixing effect, and
439 order one Rossby numbers combine both with rotational effects, consistent with modeling

440 studies of boundary layer responses to prototype SST fronts (Spall 2007a; Kilpatrick et al. 2014,
441 2016) and ocean eddy fields (Foussard et al. 2019a) in the presence of large-scale winds.

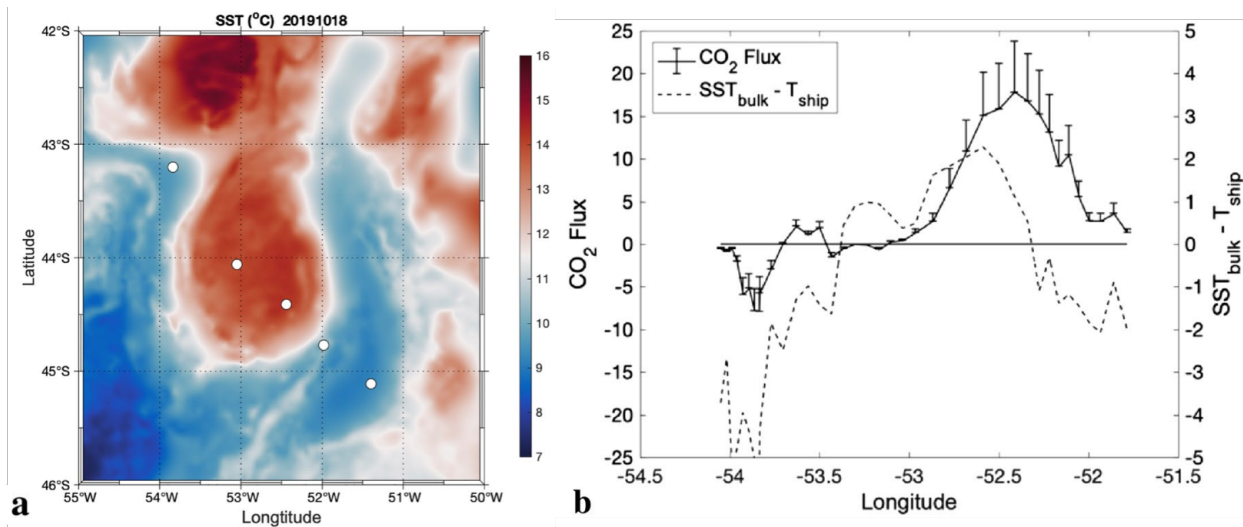
442
443 The analytical model described above considers a dry MABL without incorporating MABL
444 moisture or latent heat fluxes. The contribution of moisture to buoyancy fluxes, latent
445 heating/cooling, and overall MABL structure has not been investigated in as much detail within
446 the context of the mesoscale MABL response. However, it is anticipated to have a non-negligible
447 impact on the MABL dynamical response to mesoscale SSTA (Skylingstad and Edson 2009).
448 For instance, during CLIMODE, the buoyancy heat flux was approximately 20% larger than the
449 sensible heat flux due to moisture, and the average magnitude of the latent heat flux was ~ 2.5
450 times greater than the sensible heat flux (Marshall et al. 2009). In the tropics, the ratio of latent to
451 sensible heat flux is even larger (e.g., de Szoeke et al. 2015), so the moisture contribution is often
452 an order of magnitude greater than the sensible heat contribution. The impact of moist
453 convection during a cold air outbreak over the Gulf Stream was investigated with an LES
454 (Skylingstad and Edson 2009), showing that the latent and sensible heat fluxes are enhanced
455 over a simulated SST front resulting in stronger turbulent mixing and precipitation compared to a
456 constant SST simulation. The simulation across the SST front shows that relatively low humidity
457 values near the surface are maintained by the continual expansion of the boundary layer in the
458 entrainment layer, which mixes dry air from aloft into the MABL. This maintains the large air-
459 sea specific humidity and temperature differences necessary for strong latent and sensible heat
460 fluxes in the surface layer. Additional simulations and measurements are required to investigate
461 the role of moisture in response to mesoscale SST. For example, the analytical model could
462 provide insight by using the virtual temperature at both the sea surface and aloft.

463

464 *d. Modulation of air-sea fluxes of tracers*

465 Air-sea gas fluxes of tracers depend on the air-sea disequilibrium and processes driving
466 exchange, such as winds and breaking waves. From the ocean perspective, the disequilibrium can
467 be understood as the difference of the concentrations of a gas in the seawater, C , relative to the
468 concentration the gas would have at equilibrium with the atmosphere, C_{eq} , which, in turn, is
469 determined by the solubility of the gas in seawater. The air-sea flux F_x of a gas x then is
470 estimated as $F_x = k (C - C_{eq})$, where k is the gas transfer velocity (e.g., Woolf 1993; McGillis

471 et al. 2001; Wanninkhof et al. 2009; Dong et al. 2021). Impacts of ocean mesoscale features on
 472 the net F may be introduced via k or C_{eq} , each of which varies nonlinearly with wind speed and
 473 depends on sea state. The mesoscale may also affect C by impacting biological sources and sinks
 474 of tracers (Section 4d). Indeed, studies find local modulations of air-sea CO_2 fluxes due to the
 475 effects of mesoscale eddies on solubility, productivity, or winds (Jones et al. 2015; Song et al.
 476 2015, 2016; Olivier et al. 2021). One such study in the Southwest Atlantic Ocean detected clear
 477 spatial covariations of CO_2 flux with the MABL stability over a warm-core eddy (Figure 3; Pezzi
 478 et al. 2021). Yet, on the basin-to-global scales, positive and negative mesoscale anomalies of
 479 CO_2 fluxes appear to essentially cancel (Wanninkhof et al. 2011; Song et al. 2015). Clear
 480 separation and quantification of the individual and rectified effects of mesoscale phenomena on
 481 k , C , and C_{eq} from observations and models remain challenging, given the difficulty of capturing
 482 transient mesoscale variations in the ocean and atmosphere, including the concentration of
 483 tracers such as carbon.
 484



485 **a**
 486 **Figure 3:** (a) Observed SST ($^{\circ}\text{C}$) in the Southwestern Atlantic Ocean on 18 October 2019. The white circles
 487 denote the Po/V *Almirante Maximiano* trajectory. (b) In situ CO_2 fluxes ($\mu\text{mol m}^{-2}\text{s}^{-1}$) measured by Eddy
 488 Covariance method (solid) and atmospheric stability parameter, $\text{SST}_{\text{bulk}} - T_{\text{ship}}$ ($^{\circ}\text{C}$) (dotted), where SST_{bulk} and
 489 T_{ship} denote the sea surface and near-surface air temperatures, respectively. The error bars denote the standard
 490 error representing a 95% confidence interval. Figures adapted from Pezzi et al. (2021). The figure needs
 491 permission to reproduce.
 492

493

494 **3. Free-tropospheric, extratropical atmospheric circulation responses**

495 This section investigates atmospheric response beyond the MABL (Section 2) by focusing on
496 local and non-local circulation responses in the extratropics to SSTA patterns observed in the
497 WBC regions, including the semi-permanent SST fronts and transient mesoscale eddies. Some
498 aspects of deep convective response in the tropical atmosphere have also been attributed to
499 MABL adjustments to the mesoscale SST fields (Li and Carbone 2012; Skillingstad et al. 2019;
500 de Szoeke and Maloney 2020), although much of the studies on deep atmospheric responses
501 published to date is based on the extratropics. We start with a summary of previous studies on
502 the role of extratropical SSTA in quasi-equilibrium atmospheric circulation and storm tracks. We
503 then revisit the debates about the observed near-surface wind convergence and precipitation in
504 WBC regions diagnosed as a response to either SST variations or extratropical storms. Finally,
505 we will consider whether these processes may be important to future climate, focusing on the
506 difference between projections at high and low resolution in the oceans. The feedback processes
507 examined in this section are schematically illustrated in Figure 1a-c.

508

509 *a. Time-mean general circulation responses*

510 The question of how the extratropical atmosphere responds to variability in ocean fronts and/or
511 extratropical SSTA has been addressed over many decades. Early studies considered the linear
512 response (Hoskins and Karoly 1981; Frankignoul 1985), which predicted a shallow heating
513 response characterized by a downstream trough with a baroclinic structure. This was argued
514 against by Palmer and Sun (1985), who found a downstream ridge, with an advection of
515 temperature anomalies by mean flow acting against anomalous advection of mean temperature
516 gradients. Later, Peng et al. (1997) showed that the transient eddy response was important in
517 forming an equivalent barotropic high. More recent observational analyses find a weak low-
518 pressure response east of warm SSTA near the Gulf Stream (Wills et al. 2016) and Kuroshio
519 (Frankignoul et al. 2011; Wills and Thompson 2018). Deser et al. (2007) demonstrated that the
520 initial linear, baroclinic response is quickly (within 2 weeks) replaced with the equilibrium
521 barotropic response with a much broader spatial extent and magnitude (Ferreira and Frankignoul
522 2005, 2008; Seo et al. 2014). The adjustment time is shorter near WBC regions (Smirnov et al.
523 2015). This literature is well summarized in existing review papers (Kushnir et al. 2002; Small et
524 al. 2008; Kwon et al. 2010; Czaja et al. 2019).

525 Recent studies also indicated a strong sensitivity to the spatial resolution of the atmospheric
526 dynamics governing the large-scale circulation response. For example, Smirnov et al. (2015)
527 show that a low-resolution (1°) model induces a weak response resulting from shallow
528 anomalous heating balanced by equatorward cold air advection, consistent with the results from
529 steady linear dynamics. This contrasts with the higher resolution ($1/4^\circ$) model showing that the
530 anomalous diabatic heating is balanced by a deep vertical motion mediated by the transient
531 eddies (Hand et al. 2014; Wills et al. 2016; Lee et al. 2018). The anomalous diabatic heating and
532 the induced vertical motions maintain the climatological circulation pattern over the WBCs.

533

534 *b. Synoptic storms and storm track responses*

535 Storm tracks typically occur in the $30\text{-}50^\circ$ latitude band coincident with the climatological SST
536 fronts (Figure 4) and are associated with strong and frequent precipitation, particularly via
537 atmospheric fronts. Midlatitude storm tracks can be primarily defined in two ways (Chang et al.
538 2002; Hoskins and Hodges 2002): either using distributions of the tracks and intensity of
539 synoptic cyclones (the Lagrangian view) or as regions of strong variability or co-variability of
540 winds, geopotential height, temperature, and humidity in the lower to upper troposphere (the
541 Eulerian perspective). To better elucidate the forcing of near-surface weather by the oceans,
542 other studies also use the surface-based storm track, defined as the variance of near-surface
543 meridional winds (Booth et al. 2010, 2017; O'Neill et al. 2017; Small et al. 2019). The concept
544 of the surface storm track stems from earlier scatterometer measurements illustrating strong
545 imprints of the free-tropospheric storm tracks in the surface wind fields over the warm WBCs
546 (Sampe and Xie 2007; Bourassa et al. 2013). The reduced static stability and the enhanced
547 vertical mixing within the MABL (Figure 1d) synchronize the locations of the surface storm
548 track with the warm currents (Figure 4). The surface and free-tropospheric storm tracks are, thus,
549 dynamically coupled via deep moist convection (Czaja and Blunt 2011).

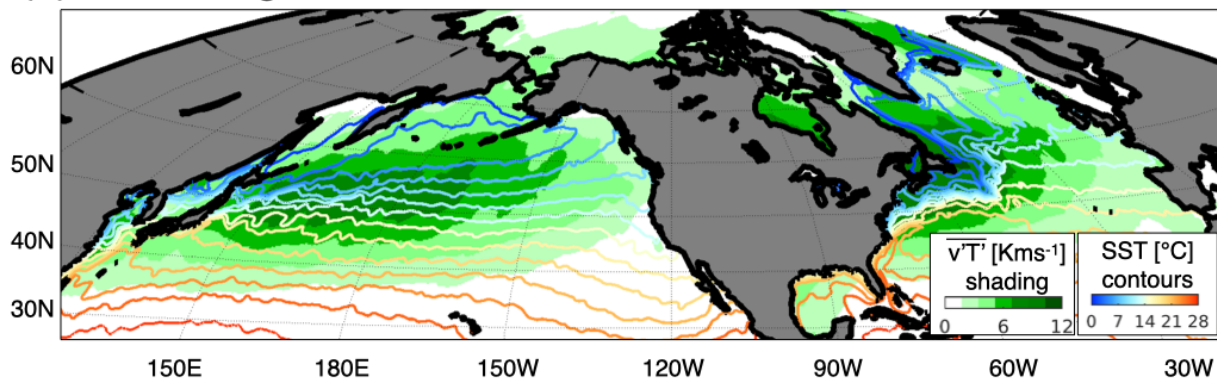
550

551 One possible mechanism of midlatitude oceanic influence on the storm track was suggested by
552 Hoskins and Valdes (1990), which found that enhanced diabatic heating by surface fluxes over
553 WBCs supports atmospheric baroclinicity, a vital element in setting the location of the storm
554 track (Hawcroft et al. 2012; Kaspi and Schneider 2013). Nakamura and Shimpo (2004) and
555 Nakamura et al. (2004) further argued that SST gradients directly influence low-level air

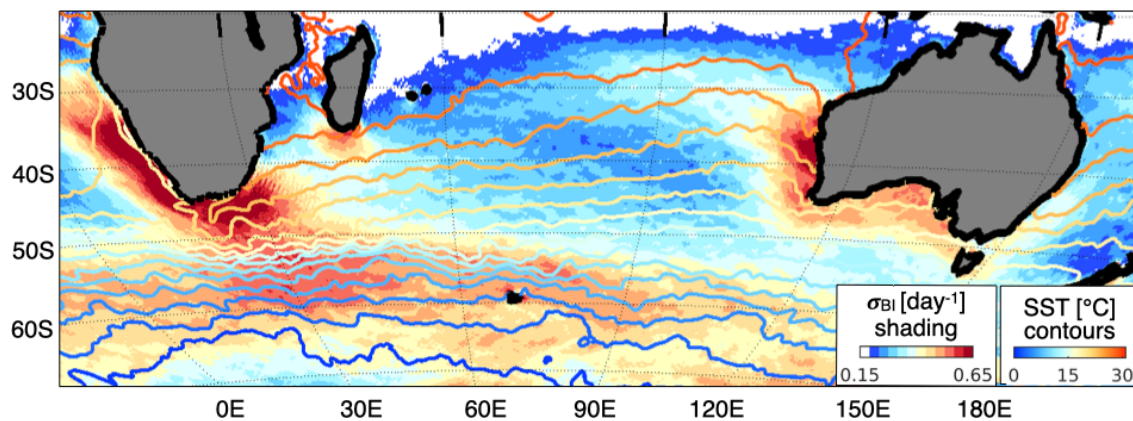
556 temperature gradients via cross-frontal gradients in sensible heat flux (Nakayama et al. 2021).
 557 The baroclinicity is measured as the atmospheric maximum Eady growth rate (Charney 1947;
 558 Eady 1949; Lindzen and Farrell 1980), such that stronger low-tropospheric baroclinicity is
 559 associated with weaker static stability and a stronger meridional air temperature gradient (see
 560 Figure 4 caption). Both conditions are observed over WBCs. Hence, the anchoring effect by
 561 cross-frontal differential heat supply from the ocean is consistent with the formation of a storm
 562 track over the WBC SST fronts (Nonaka et al. 2009; Hotta and Nakamura 2011), while diabatic
 563 heating over the warm portion of the WBC SST fronts to the warm and cold sectors of the
 564 cyclones supports the growth of transient baroclinic waves (Booth et al. 2012; Willison et al.
 565 2013; Hirata and Nonaka 2021; Figure 1b,c).

566

(a) Climatological storm track over Kuroshio and Gulf Stream



(b) Climatological “baroclinicity” over Agulhas and ACC



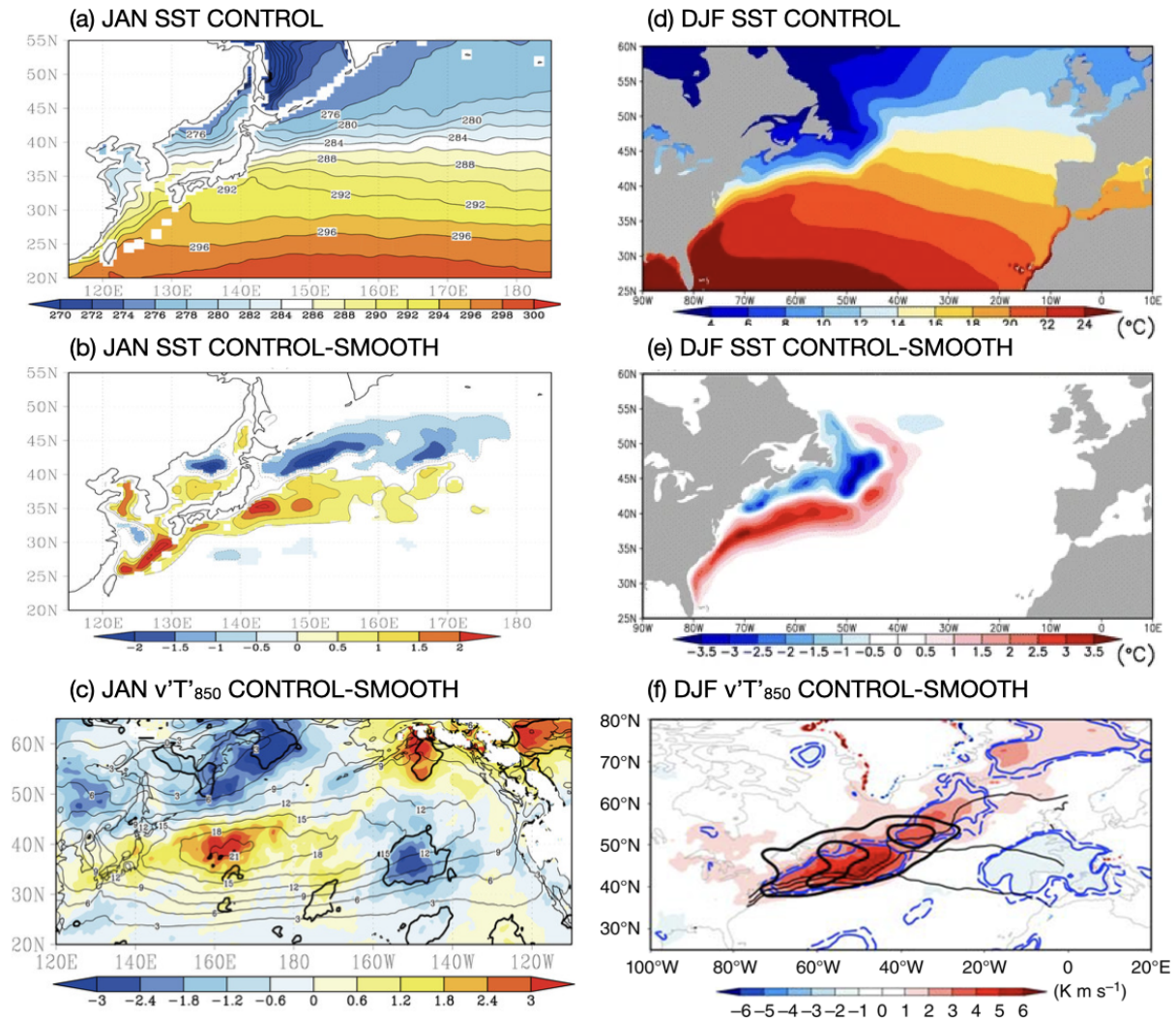
567

568 **Figure 4:** The climatological relationship of the extratropical storm tracks with the SST fields in (a) Kuroshio-
 569 Oyashio Extension and Gulf Stream in the Northern Hemisphere and (b) Agulhas Current and the Antarctic
 570 Circumpolar Current systems in the south Indian Ocean. The atmospheric storm track is estimated in (a) as the
 571 time-mean meridional heat transport by atmospheric transient eddies, $\overline{v'T'}$ at 850 hPa (low troposphere), where
 572 primes denote the 2-8-day bandpass filtered fields and the over-bar indicates the time-mean, and in (b) as the
 573 atmospheric maximum Eady growth rate, defined as the most unstable baroclinic mode whose growth rate is

574 scaled as the magnitude of the baroclinicity vector, $|\sigma_{BI}| = 0.31 \left(\frac{g}{N\theta} \right) \left| -\frac{\partial\theta}{\partial y}, \frac{\partial\theta}{\partial x} \right|$, at 850 hPa, where g is the
575 gravitational acceleration, N is the buoyancy frequency, and θ is the potential temperature. These storm track
576 quantities are derived from ERA5. The SST climatology is obtained from the NOAA daily Optimum
577 Interpolation dataset. The climatologies are calculated from 2010 to 2015.
578

579 A standard method to diagnose the SST forcing mechanism of the storm track is to run a pair of
580 AGCM simulations, one using observed SSTs (CONTROL), and another using a spatially-
581 smoothed SST field with weaker gradients (SMOOTH), which also alters absolute SST (Figure
582 5). Alternatively, AGCMs are forced by shifting the latitude of the SST fronts or filtering
583 mesoscale eddy SSTs (Seo et al. 2017). Such AGCM simulations indicate a strengthening of the
584 storm track near the Kuroshio-Oyashio Extension (KOE) (Kuwano-Yoshida and Minobe 2017)
585 and the Gulf Stream (O'Reilly et al. 2017) in CONTROL near the climatological maximum
586 cyclogenesis (Figure 5). Altered storm activity over the WBC regions influences the intensity of
587 the coastal storms, and, thereby, inland weather near the Kuroshio (Nakamura et al. 2012;
588 Hayasaki et al. 2013; Sugimoto et al. 2021), Gulf Stream (Infanti and Kirtman 2019; Hirata et al.
589 2019; Liu et al. 2020), and the Agulhas Current (Singleton and Reason 2006; Nkwinkwa
590 Njouodo et al. 2018).

591



592
 593 **Figure 5:** (Left) January observed SST, its difference (CONTROL-SMOOTH), and the difference
 594 (CONTROL-SMOOTH) in storm tracks over the North Pacific Ocean. The thin black contours show $\overline{v'T'}$
 595 from the CONTROL case. Thick contours denote the 95% confidence level. (Right) As in (left) but for over the
 596 North Atlantic. Black contours in (f) denote atmospheric Eady growth rate at 775 hPa. The dashed and solid
 597 blue contours indicate significant differences at the 10 and 5% levels, respectively. Figures adapted from
 598 Kuwano-Yoshida and Minobe (2017) and O'Reilly et al. (2016, 2017). The figure needs permission to
 599 reproduce.

600

601 Recent studies indicate that atmospheric mesoscale phenomena within the storm tracks, such as
 602 atmospheric fronts, directly interact with the WBC fronts. Parfitt and Czaja (2016) used
 603 reanalysis data over the Gulf Stream, and Parfitt et al. (2016) used AGCM simulations over the
 604 KOE to argue that the cross-frontal sensible heat flux gradients across the SST fronts exert
 605 "thermal damping or strengthening" of atmospheric fronts depending on the space-time
 606 alignment between the SST gradients and atmospheric fronts with shared cross-frontal length-
 607 scales (Figure 1b-c). The most significant diabatic heating by surface fluxes is concentrated on

608 the narrow space-time scales at which the cold sectors of the atmospheric front coincide with the
609 warm sector of the SST fronts (Figure 1c), significantly enhancing precipitation associated with
610 the atmospheric fronts and often facilitating explosive cyclogenesis (Hirata and Nonaka 2021
611 and references therein).

612

613 In contrast, other studies emphasize the limited role of SST fronts on extreme cyclones. AGCM
614 experiments by Tsopouridis et al. (2021) indicated that the direct impacts of sharp SST fronts on
615 individual cyclones over the Gulf Stream and KOE are weak, although SST fronts induce
616 significant indirect responses in large-scale environments in which such storms form. Using an
617 analytic model, Reeder et al. (2021) showed that diabatic frontogenesis over the WBCs
618 intensifies atmosphere fronts only when strong and rapidly propagating synoptic systems are not
619 already in the environment.

620

621 Much uncertainty remains in model simulations and observational analysis regarding the relative
622 importance of SST gradients causing cross-atmospheric frontal sensible heat flux gradients vs.
623 absolute SST affecting the large-scale condensational heating over warm currents. Another
624 critical issue is that since the SST contributions to the precipitation from the warm and cold
625 sectors of extratropical cyclones differ in terms of magnitude and spatial distribution (i.e.,
626 broader for the warm sectors and more "anchored" to the SST fronts for the cold sectors, e.g.,
627 Vanni re et al. 2017), the cold sector contribution might have been dominating the sensitivity of
628 relatively high-resolution (~50 km) AGCM simulations to SST smoothing. It remains an open
629 question whether even higher resolution AGCMs might amplify a sensitivity from the dynamics
630 of the warm sectors, including atmospheric mesoscale instabilities developing on the warm
631 conveyor belt (Czaja and Blunt 2011; Sheldon et al. 2017).

632

633 *c. Near-surface wind convergence and vertical motion over the WBCs*

634 A crucial part of the storm track response to SST is precipitation, which tends to cluster around
635 the WBCs and is associated with high near-surface wind convergence (NSWC) and substantial
636 vertical ascent. The climatological NSWC coincides with the ocean fronts and the Laplacians of
637 SST and SLP, which indicates that the boundary layer process depicted by linear Ekman
638 dynamics is germane to the observed NSWC and precipitation responses (Feliks et al. 2004;

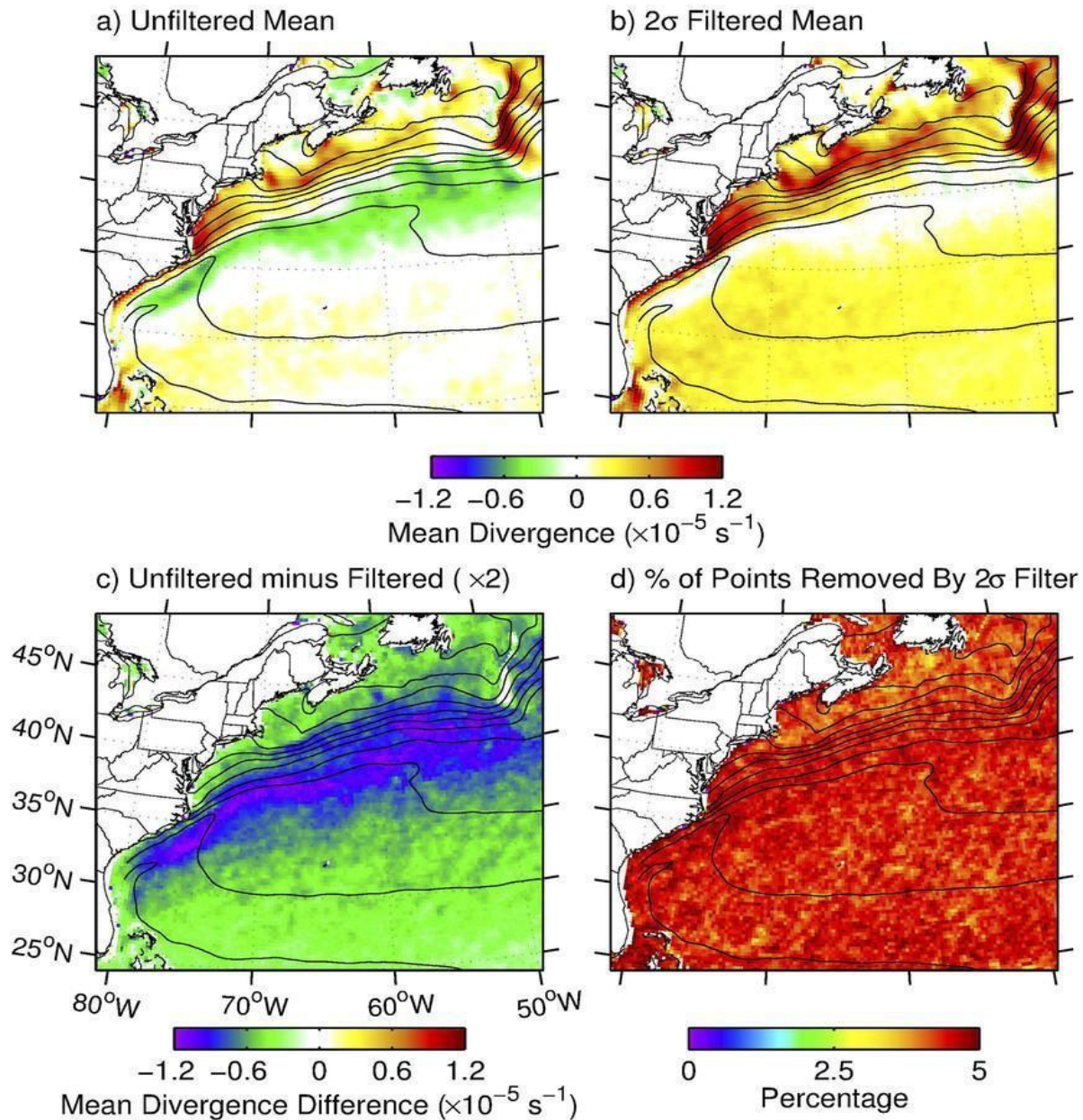
639 Minobe et al. 2008, 2010). However, the unambiguous attribution of NSWC to the steady
640 Ekman-balanced mass adjustment mechanism remains difficult due to the coexistence of
641 extratropical storm tracks with the WBC currents, which also induce minima in the time-mean
642 SLP Laplacian over the SST fronts (O'Neill et al. 2017).

643
644 O'Neill et al. (2015) show from QuikSCAT observations and a regional atmospheric model that
645 linear boundary layer dynamics cannot explain the daily time-scale occurrence of NSWC since,
646 on rain-free days, surface divergence dominates even though the SST Laplacian would indicate
647 convergence (Figure 6). Using an extreme value filter, O'Neill et al. (2017) further show that
648 NSWC and vertical motion over the Gulf Stream are highly skewed and consist of infrequent yet
649 extreme surface convergence events and more frequent but weak, divergent events, such that the
650 median surface flow field is weakly divergent or nearly non-convergent (Figure 6). Parfitt and
651 Czaja (2016) and Parfitt and Seo (2018) argue that much of the precipitation and NSWC are
652 associated with atmospheric fronts, given that only a weak near-surface divergence remains
653 when the contribution from atmospheric fronts is removed (Rousseau et al. 2021). In contrast,
654 Masunaga et al. (2020a,b) showed that storms and fronts of moderate intensity are significant
655 contributors to the time-mean convergence observed over the Gulf Stream and KOE.

656
657 Current research emphasizes identifying how and why atmospheric fronts align with and linger
658 over ocean fronts in all major WBCs and whether there is an additional underlying, steady,
659 small-scale boundary layer effect. There might exist a distinct temporal dependence of the
660 NSWC over WBC SSTs, where atmospheric fronts govern its day-to-day variability, while the
661 pressure adjustment and vertical mixing mechanisms provide lower frequency modulations (e.g.,
662 Brachet et al. 2012; Small et al. 2022).

663

10-yr Mean All-Weather QuikSCAT Divergence



664

665 **Figure 6:** Maps of the 10-yr-mean QuikSCAT all-weather divergence (a) consisting of all points; (b) after
666 application of the 2σ temporal extreme-value filter; (c) difference between (a) and (b); and (d) the percentage
667 of divergence points removed by the 2σ extreme-value filter. The contours in each panel are of the 10-yr-mean
668 Reynolds SST with a contour interval of 2°C . From O'Neill et al. (2017). The figure needs permission to
669 reproduce.

670

671 *d. Non-local downstream atmospheric circulation responses*

672 The upstream storm track variability leading to downstream development of the storm track is an
673 essential characteristic of midlatitude baroclinic waves (Chang 1993). The altered synoptic-scale

674 disturbances over the baroclinically unstable western basins (Section 3b) radiate energy
675 downstream, influencing the growth of a subsequent baroclinic wave toward the eastern basins
676 (e.g., Chang and Orlanski 1993). The downstream atmospheric circulation also results from the
677 synoptic eddy-mean flow interactions, where low-frequency atmospheric circulation is coupled
678 with the transient eddy activity modified over the WBCs (e.g., Haines and Marshall 1987;
679 Nakamura and Wallace 1990). Here, downstream (or remote, or non-local) refer to the region
680 immediately east of the SST forcing and the tail-end of the storm track abutting the west coasts
681 of the continents, as illustrated in Figure 1a.

682

683 Many AGCM studies demonstrate a non-local, downstream response in the storm track to WBC
684 SST forcing. Using the observational datasets, Wills et al. (2016) and Joyce et al. (2018)
685 identified significant transient atmospheric circulation responses (storm track and atmospheric
686 blocking) downstream that lag the SSTA in the Gulf Stream extension by several weeks to
687 months. The modeling studies by O'Reilly et al. (2016, 2017) showed that a strengthened storm
688 track over the Gulf Stream leads to the northward shifted atmospheric eddy-driven jet and the
689 increased European blocking frequency far downstream. Along a similar line, Lee et al. (2018)
690 suggested that SST biases near the Gulf Stream trigger extended biases in the simulation of deep
691 convection and downstream circulation via Rossby wave response.

692

693 In the North Pacific, O'Reilly and Czaja (2015) found that baroclinic eddies grow faster when the
694 KE front is in its stable regime (stronger SST gradients). The local shift in baroclinic wave
695 activity leads to the early barotropitization of the baroclinic eddies downstream, resulting in
696 weaker poleward eddy heat flux and increased occurrence of blocking in the eastern Pacific. An
697 AGCM study by Kuwano-Yoshida and Minobe (2017) also suggested the enhanced storm track
698 by the KOE SST fronts leads to a northward shifted storm track in the eastern Pacific. Ma et al.
699 (2015, 2017) showed from AGCM simulations that the transient SSTA associated with the KOE
700 mesoscale eddies leads to a northward shifted storm track and reduced precipitation in parts of
701 western North America (Foussard et al. 2019b; Liu et al. 2021; Siqueira et al. 2021).

702

703 In the Southern Ocean, Reason (2001) showed that amplified cyclone activity over the warm
704 Agulhas Current yielded an enhanced storm track in the southeast Indian Ocean. Recent

705 aquaplanet AGCM experiments have also demonstrated the critical role of the oceanic fronts in
706 shaping the structure of the baroclinic annular mode variability (e.g., Sampe et al. 2013; Ogawa
707 et al. 2016; Nakayama et al. 2021), leading modes of variability of the extratropics (e.g.,
708 Thompson and Wallace 2000). Evidence exists that the oceanic frontal zones also impact the
709 troposphere-stratosphere interactions (e.g., Hurwitz et al. 2012; Ogawa et al. 2015; Omrani et al.
710 2014), potentially affecting the entire hemispheric climate patterns.

711

712 *e. Climate change*

713 Climate change simulations for the 21st Century have emphasized the critical role of ocean
714 circulation leading to natural modes of variability such as ENSO and PDO (Seager et al. 2001),
715 the projected weakening of the Atlantic Meridional Overturning Circulation (AMOC; Weaver et
716 al. 2012), and the delayed warming of the Southern Ocean (Marshall et al. 2014). These changes
717 are relevant to the observed and projected intensification and poleward shift of the Kuroshio and
718 Agulhas, weakening of the Gulf Stream, and changes in the frontal systems of the Antarctic
719 Circumpolar Current (ACC) (e.g., Wu et al. 2012; Yang et al. 2016; Sen Gupta et al. 2021).

720

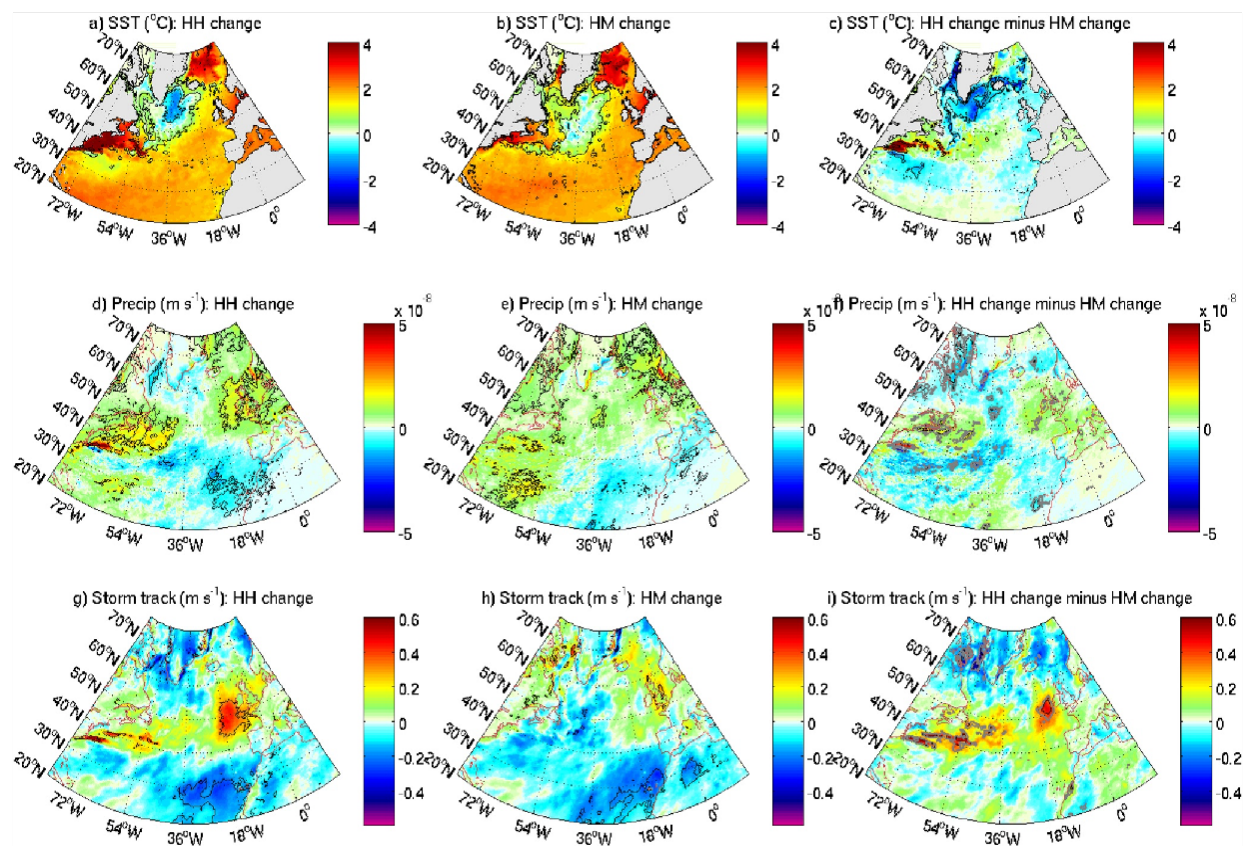
721 The IPCC report (IPCC, 2022) indicates that during the 21st Century, the North Pacific storm
722 track will most likely shift poleward, the North Atlantic storm track is unlikely to have a simple
723 poleward shift, and the Southern Hemisphere storm track will likely shift poleward.

724 Understanding these regional differences in projected changes in midlatitude storm tracks and
725 precipitation and their association with the predicted WBC changes have been the primary goals
726 of high-resolution CGCM studies, especially those that contrast the CGCMs with the eddy-rich
727 ocean (typically 0.1° resolution) to those with the eddy-parameterized ocean (0.5-1°). These
728 studies with increased ocean model resolution to mitigate the known biases in representing the
729 WBC dynamics and separation show distinct responses in SSTs and storm tracks in the WBC
730 regions to anthropogenic climate change.

731

732 In these eddy-rich simulations, the KOE front shifted equatorward, contrary to projections by the
733 eddy-parameterized IPCC-class CGCMs, which likely reflects the large natural variability in the
734 North Pacific (Taguchi et al. 2007; Seager and Simpson 2016). In the North Atlantic, the Gulf
735 Stream separation tends to be too far north in lower resolution models, an issue common to other

736 WBCs, but is improved in eddy-rich models. This makes it possible for the separation to move
 737 northwards as a response to AMOC weakening in eddy-rich models (Gervais et al. 2018;
 738 Moreno-Chamarro et al. 2021; Grist et al. 2021), leading to a significant projected ocean
 739 warming near the US eastern coastline (Figure 7; Karmalkar and Horton 2021). In the Southern
 740 Ocean, CMIP5-based climate change simulations indicate delayed warming, often attributed to
 741 stratospheric ozone depletion (McLandress et al. 2011; Polvani et al. 2011). However, the recent
 742 satellite observations and eddy-rich CGCMs simulations indicate a ubiquitous cooling trend
 743 (1961–2005) poleward of the ACC due to the effects of resolved ocean eddies (Bilgen and
 744 Kirtman 2020). Analysis of eddy-rich ocean simulations also indicates warmer and stronger
 745 Southern Hemisphere WBCs, suggesting that resolved ocean eddies play a critical role in long-
 746 term SST changes.
 747



748
 749 **Figure 7:** (a–c) 2031–2050 minus 1951–1970 differences simulated by the HadGEM3-GC3.1, with 25 km
 750 atmospheric resolution coupled to 1/4° ocean (eddy-permitting, HM) and 1/12° ocean (eddy-rich, HH): SST
 751 (°C) (a) HH and (b) HM, precipitation (ms^{-1}) (d) HH and (e) HM; surface storm track (ms^{-1}) (g) HH and (h)
 752 HM. Panels (c, f, i) show the difference between the HH future change and the HM change. The black lines
 753 denote the 95% significance. Gray lines in (c,f,i) denote the 90% significance. From Grist et al. (2021). The
 754 figure needs permission to reproduce.

755 The reorganization of the oceanic frontal zone and its associated eddy field modulates the
756 atmospheric low-level baroclinicity and the strength and location of the diabatic heating source
757 for the atmosphere. It is clear from this and other studies (Woollings et al. 2012; Winton et al.
758 2013; Keil et al. 2020) that such features would not occur without ocean circulation changes.
759 However, the exact pattern of large-scale SST change is highly dependent on the ocean model
760 and its resolution (Saba et al. 2016; Menary et al. 2018; Alexander et al. 2020), which also
761 affects the projected WBC responses to climate change (Jackson et al. 2020). Climate projections
762 with eddy-rich oceans have typically been performed with a small number of realizations and for
763 short durations due to high computational costs (e.g., Haarsma et al. 2016). Currently, high-
764 resolution coupled climate modeling projects are underway with much longer integration and
765 multi-ensembles (e.g., Chang et al. 2020; Wengel et al. 2021). These efforts will enable a robust
766 assessment of the forced responses in WBC and ocean circulation from natural variability in
767 response to projected changes in the large-scale climate.

768
769

770 **4. Feedback of atmospheric responses onto the ocean**

771 The new insights gained from the studies discussed in Section 3 have also led to improved
772 process understanding and notable revisions of theories of ocean circulation. This section
773 discusses current knowledge of ocean feedback mechanisms, including feedback impacts on
774 ocean biogeochemical cycles, and theories of ocean circulation and model parameterizations to
775 account for eddy-atmosphere interaction. The processes covered in this section correspond
776 mainly to Figures 1e,f.

777

778 *a. Feedback on ocean circulation*

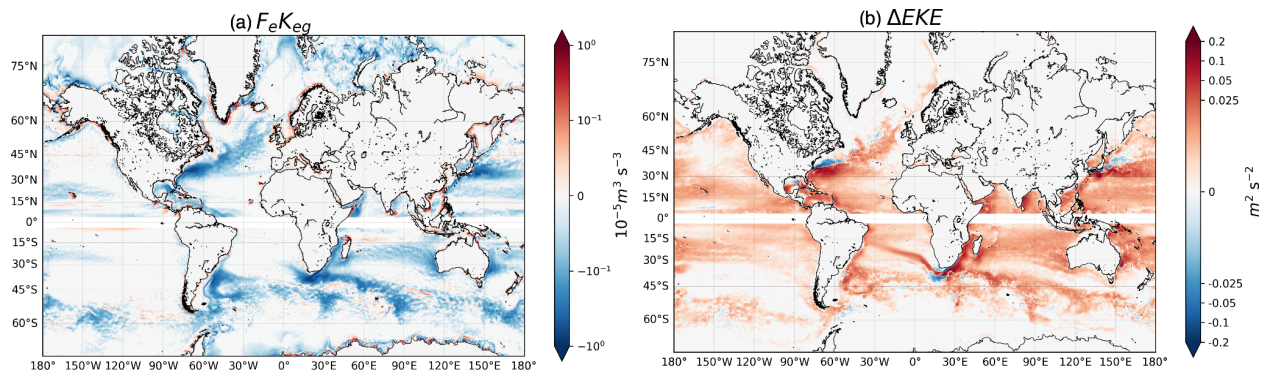
779 For simplicity, we consider two categories of oceanic mesoscale effects on air-sea fluxes: SST
780 impacts (thermal) described in Section 2b1 and surface current impacts (mechanical) in Section
781 2b2. The thermal feedback (TFB) results from kinematic and thermodynamic responses in the
782 MABL to mesoscale SSTs, modifying the wind stress and heat fluxes. The current feedback
783 (CFB) represents the frictional processes by which the surface ocean current alters the wind
784 stress, near-surface wind, and turbulent heat fluxes. This subsection focuses on the respective
785 feedback impacts of the air-sea fluxes on ocean circulation.

786 1) Thermal feedback (TFB) effect
787 Observed near-surface wind stress responses to mesoscale processes by Chelton et al. (2004)
788 were interpreted based mainly on the TFB effect. Vecchi et al. (2004) and Chelton et al. (2007)
789 hypothesized that the wind stress curl responses to SST fronts exert a vital feedback mechanism
790 driving the evolution of SST fronts via resulting anomalous Ekman pumping. Spall (2007b)
791 considered the impacts of SST-induced Ekman pumping on baroclinic instability in the ocean in
792 the modified linear theory by Eady (1949), showing that the SST-induced Ekman pumping
793 adjusts the growth rate and wavelength of the most unstable waves, especially the low-latitude
794 flows with strong stratification. Hogg et al. (2009) extended SST-induced Ekman pumping to an
795 idealized double-gyre circulation in mid-latitudes, showing that it destabilizes the eastward jet
796 with the enhanced cross-gyre potential vorticity fluxes, stabilizing the double gyre circulation by
797 30-40%.
798
799 Mesoscale SSTAs are damped by induced turbulent heat fluxes (THF), resulting in a negative
800 SST-THF correlation at oceanic mesoscales. Over the KOE, Ma et al. (2016) examined this
801 mesoscale SSTA damping in the context of the eddy potential energy (EPE) budget and the
802 Lorenz energy cycle. Compared to the eddy-filtered coupled model simulation (using a 1000 km-
803 by-1000 km boxcar filter), the eddy-unfiltered simulations showed a significant increase (>70%)
804 in diabatic EPE dissipation, leading to a decrease in eddy kinetic energy (EKE) by 20-40%, most
805 strongly at wavelengths shorter than 100 km (Figure 1d). Other studies find that TFB has a weak
806 impact on EKE (Seo et al. 2016; Seo 2017). It is possible that a large filter cutoff, as used in Ma
807 et al. (2016), overestimates EKE damping and may also smooth large-scale meridional SST
808 gradients, altering the large-scale wind curl and the mean circulation. Bishop et al. (2020)
809 evaluated the EPE damping over the global oceans using eddy-resolving climate model
810 simulations to find that the diabatic EPE damping was systematically stronger over warm-core
811 eddies (Figure 1c,e). Other studies point out that the efficacy of the negative SST-THF
812 correlation in the maintenance of the mesoscale SSTA and their gradients depends on the
813 distribution of the mixed-layer depth, which modulates the effective heat capacity, vertical eddy
814 heat transport, and hence the sensitivity of the SST to the heat flux anomaly (e.g., Tozuka et al.
815 2017; 2018; Jing et al. 2020).
816

817 2) Current feedback (CFB) effect

818 Although weaker than surface winds, surface currents modify surface stress directly by altering
819 wind speed (Bye 1986). By modulating the stress, the CFB exerts a "bottom-up" effect on the
820 wind, where a positive current anomaly causes a positive wind anomaly via a negative stress
821 anomaly (Renault et al. 2016a; 2019a). The CFB effect has initially focused on impact on wind
822 stress. Using satellite and in situ data, Kelly et al. (2001) showed that CFB reduces the median
823 wind stress from 20% to 50% near the equator, and Chelton et al. (2004) observed a clear imprint
824 of the Gulf Stream flow on the surface stress and the curl.

825



826

827 **Figure 8:** (a) Geostrophic eddy wind work ($10^{-5} \text{ m}^3 \text{ s}^{-3}$) estimated from the EC-Earth global coupled simulation
828 (15 km atmosphere coupling $1/12^\circ$ ocean) with current feedback (CFB). The negative values indicate a
829 momentum transfer from geostrophic mesoscale currents to the atmosphere. This sink of energy is the primary
830 driver of the damping of EKE illustrated in (b), as the difference of EKE ($\text{m}^2 \text{ s}^{-2}$) between the simulations
831 without CFB and with CFB. The positive values indicate the relative increase in EKE in the *absence* of CFB
832 due to the transfer of the momentum to the atmosphere. The geostrophic wind work and EKE are both
833 estimated over 30 years. Details about the coupled model and experiments can be found in Renault et al.
834 (2019c).

835

836 Several studies have highlighted the role of CFB as a "top drag" (Dewar and Flierl 1987), acting
837 on the oceanic circulation over a wide range of space-time scales. At the large-scale where the
838 currents tend to flow downwind (Figure 1e), CFB reduces the mean energy input from the
839 atmosphere to the ocean and slows down the mean circulation (Pacanowski et al. 1997). By
840 weakening net energy input to the ocean, CFB triggers a host of changes in eddy-mean flow
841 interactions and the inverse cascade of energy, weakening baroclinic and barotropic instabilities
842 and mesoscale activity (Renault et al. 2017b, 2019a; Figure 8). When the wind and current are in
843 the opposite sense, the CFB serves as a conduit of energy from the ocean to the atmosphere,
844 which can be seen from satellite data as negative mean and eddy wind work (Figure 8a; Scott
845 and Xu 2009; Renault et al. 2016a,b, 2017a). Numerous studies have demonstrated a strong EKE

846 damping effect of ~30% (See references in Jullien et al. 2020; Figure 8b). CFB also induces
847 additional Ekman pumping that weakens an eddy (Gaube et al. 2015) and influences the upper-
848 ocean stratification and SST (Seo et al. 2019; Song et al. 2020).

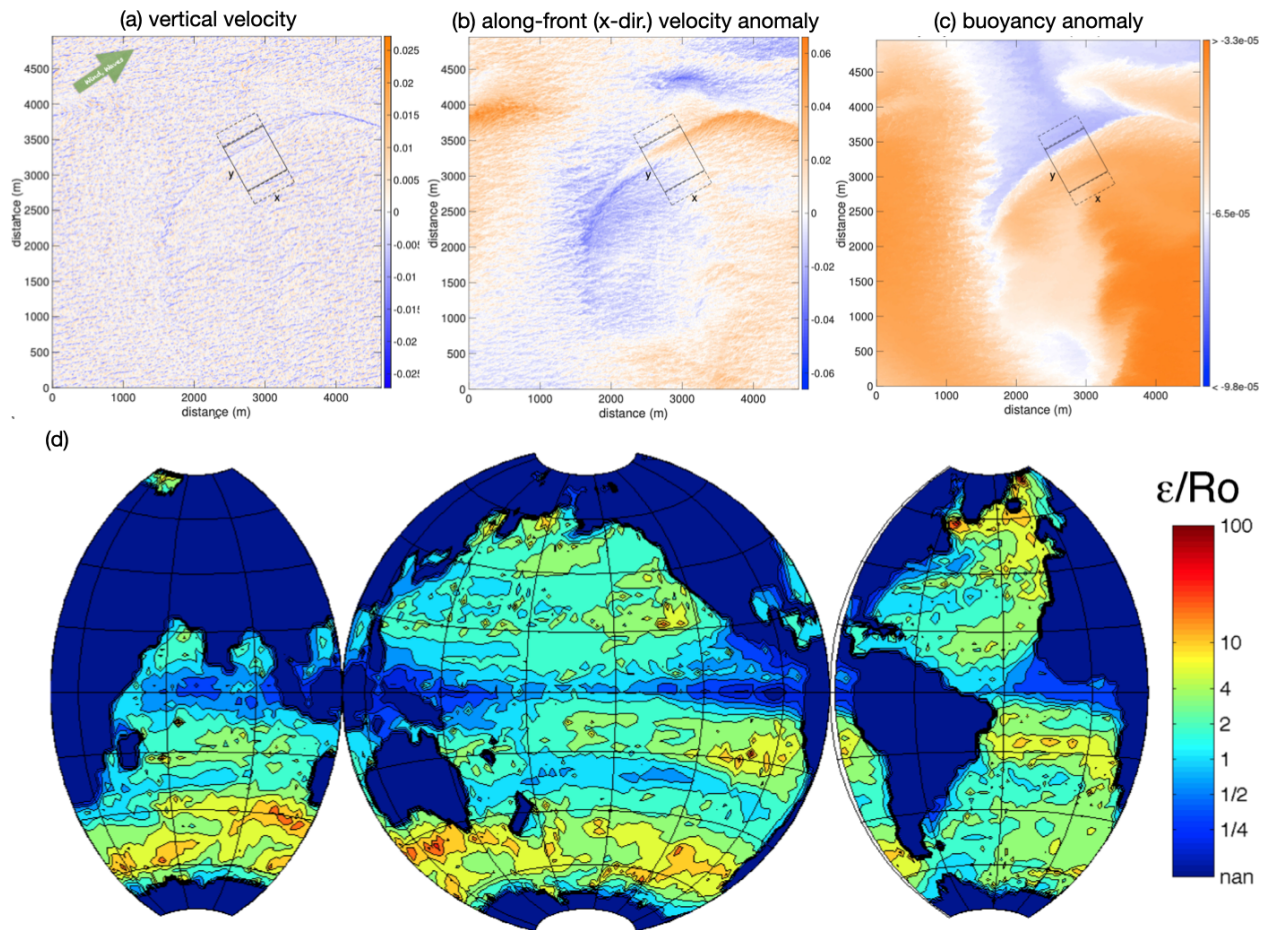
849
850 Recent studies also have emphasized the CFB impact on near-surface winds (Renault et al.
851 2016a, 2017a, 2019a). Over the shelf oceans where the current speed at tidal frequencies well
852 exceeds the wind speed, tidal currents induce tidal winds, with an amplitude of about one-third
853 of the underlying tidal currents (Renault and Marchesiello 2022). Since the wind curl is more
854 strongly impacted by current gradients (Shi and Bourassa 2019), the consideration of wind-
855 current coupling at tidal frequency might be necessary for the simulation and prediction of
856 surface winds and the MABL momentum EKE balances in the offshore environments.

857
858 There are several open questions. First, little is known about CFB at the submesoscale. For the
859 US West Coast, Renault et al. (2018) highlighted a submesoscale dual effect of CFB: it damps
860 submesoscale eddies but also catalyzes submesoscale current generation. By affecting mixing,
861 stratification, and eddy variability, Second, CFB modulates biogeochemical variability
862 (McGillicuddy et al. 2007), yet the detailed mechanisms behind the biogeochemical impacts are
863 not fully understood, although the impact depends highly on background stratification (e.g.,
864 Kwak et al. 2021). Finally, since CFB and TFB coexist where mesoscale currents are strong
865 (Song et al. 2006; Seo et al. 2007; Takatama and Schneider 2017; Renault et al. 2019b; Shi and
866 Bourassa 2019), CFB likely influences large-scale boundary-layer moisture, clouds,
867 precipitation, and atmospheric circulation via rectified effects. However, this downstream
868 influence is only beginning to be explored (e.g., Seo et al. 2021).

869
870 *b. Wave-current interactions near ocean fronts*

871 While sea state is a salient aspect of air-sea fluxes (Fairall et al. 1996; Cavaleri et al. 2012; Edson
872 et al. 2013), there are other aspects related to surface wave interactions with (sub)mesoscale
873 currents potentially important for small-scale air-sea interaction (Section 6c). For example, it has
874 long been known that sheared currents affect the propagation of surface wave rays (Villas Bôas
875 and Young 2020). In the open ocean, the spatial gradients in mesoscale surface currents
876 dominate the variability of significant wave height, leading to the refraction of waves near steep

877 vorticity gradients (Ardhuin et al. 2017; Villas Bôas et al. 2020). Similarly, the underpinnings of
 878 the Craik-Leibovich theory of Langmuir turbulence specify that rectification of wave-vorticity
 879 interactions in the upper ocean leads to Stokes forces, which can cause substantial wave effects
 880 on currents (Leibovich et al. 1983; Lane et al. 2007). The LES models that include vortex forces
 881 and regional models that include the wave refraction by currents (Romero et al. 2020) illustrate
 882 the frontal adjustment and frontogenesis triggered or enhanced by surface wave interactions
 883 (McWilliams and Fox-Kemper 2013; Suzuki et al. 2016; Sullivan and McWilliams 2019).
 884 Examples are provided in Figure 9 (upper panel), where a submesoscale density front in the
 885 downwind and down-Stokes direction interacts with Langmuir turbulence. Strong overturning
 886 circulation (downwelling) sharpens the front and strengthens the along-front jet. Classic balances
 887 are altered by waves to yield the wavy Ekman balance (McWilliams et al. 2012), the wavy
 888 geostrophic balance (McWilliams and Fox-Kemper 2013; Figure 9, lower panel), and the
 889 baroclinic and symmetric instabilities affected by waves (Haney et al. 2015).
 890



891

892 **Figure 9:** (Upper panel) Examples of a front interacting with Langmuir turbulence (box centered on this
893 feature), which is aligned in the downwind and down-Stokes direction. (a) Vertical velocity (ms^{-1}) at $z=-$
894 11.25m shows ubiquitous Langmuir cells, but also a long, coherent (downwelling) overturning circulation
895 along the front due to frontogenesis and accelerated by the Stokes shear force. (b) Along-front (x-direction)
896 velocity anomaly (with respect to the horizontal mean, ms^{-1}) at $z=-11.25$ m shows the frontal flow. (c)
897 Buoyancy anomaly (with respect to the horizontal mean, ms^{-2}) at $z=-11.25$ shows the front characterized by a
898 sharp transition in buoyancy (or temperature). Adapted from Suzuki et al. (2016). (d) Estimated ratio of ϵ
899 (strength of Stokes drift-induced vertical acceleration vs. buoyancy, an indicator of wave contributions added
900 to the traditional hydrostatic balance) to Rossby number (indicating geostrophic balance). This ratio implies
901 the deviation from the hydrostatic balance due to waves compared to the geostrophic balance due to advection.
902 This estimate is based on the de Boyer-Montegut et al. (2004) mixed layer depth climatology (h) and a global
903 simulation of WaveWatch3 and AVISO geostrophic velocity. Figures redrawn from McWilliams and Fox-
904 Kemper (2013).
905

906 *c. Physics of ocean mesoscale processes and air-sea interaction*

907 Traditionally, mesoscale and submesoscale eddy parameterizations have been deterministic and
908 focused only on effects on the mean and variance of tracers (Gent and McWilliams 1990; Fox-
909 Kemper et al. 2011), while neglecting rectified effects on air-sea coupling. However, in
910 simulations where some eddies are resolved, deterministic closures do not stimulate a resolved
911 eddying response or backscatter (e.g., Bachman et al. 2020). In response, there is a growing
912 desire to implement stochastic parameterizations of the eddy transport into non-eddy-resolving
913 models, for example, via uncertainty in location (Memin 2014), transport (Drivas et al. 2020),
914 closure (Nadiga 2008; Jansen and Held 2014; Zanna et al. 2017; Bachman et al. 2020), or
915 equation of state (Brankart 2013). These efforts should include stochastic parameterizations of
916 the eddy-driven air-sea coupling (Ma et al. 2016; Bishop et al. 2020; Jing et al. 2020). As
917 stratification and rotation parameters vary globally, building scale awareness into
918 parameterizations is also crucial (Hallberg 2013; Dong et al. 2020, 2021). Changing the relative
919 orientation of atmospheric winds and oceanic fronts leads to qualitatively different results (e.g.,
920 enhancement vs. suppression of submesoscales, Figure 1f), implying that directional subgrid
921 information will be necessary to consider (e.g., D'Asaro et al. 2011; Suzuki et al. 2016;
922 McWilliams 2016). Observed air-sea fluxes are highly variable, indicating a response to high
923 spatio-temporal variability (Yu 2019), scale dependence (Bishop et al. 2017, 2020), and sea state
924 dependence (Kudryavtsev et al. 2014), thus offering the potential for stochastic implementation.
925 While idealized studies have begun to develop a process-level understanding (Sullivan et al.
926 2020, 2021), no realistic model implementation of stochastic air-sea fluxes seems to have been
927 evaluated carefully.

928 *d. Impacts on primary productivity*

929 Mesoscale air-sea interaction can also influence biogeochemical environments and primary
930 productivity (e.g., McGillicuddy 2016). Satellite observations show that the wind stress
931 responses to mesoscale SST and currents introduce perturbation Ekman upwelling and
932 downwelling (e.g., Gaube et al. 2015), leading to dramatic mid-ocean mesoscale plankton
933 blooms, such as those observed in the nutrient-replete subtropics (e.g., McGillicuddy et al.
934 2007). Additionally, eddy-induced modifications of wind stress impact vertical mixing in the
935 upper oceans. Eddy effects on mixed-layer depths are asymmetric between anticyclones and
936 cyclones (e.g., Dufois et al. 2017; Hausmann et al. 2017). However, to what extent this
937 asymmetry stems from the mesoscale modulations of surface wind stress has yet to be
938 determined. Considering the prevalence and persistence of nonlinear mesoscale eddies in the
939 global oceans (Chelton et al. 2011a,b), the relevance of mesoscale eddy impacts on primary
940 productivity via eddy-wind interaction needs robust quantification.

941

942

943 **5. State of observational capabilities**

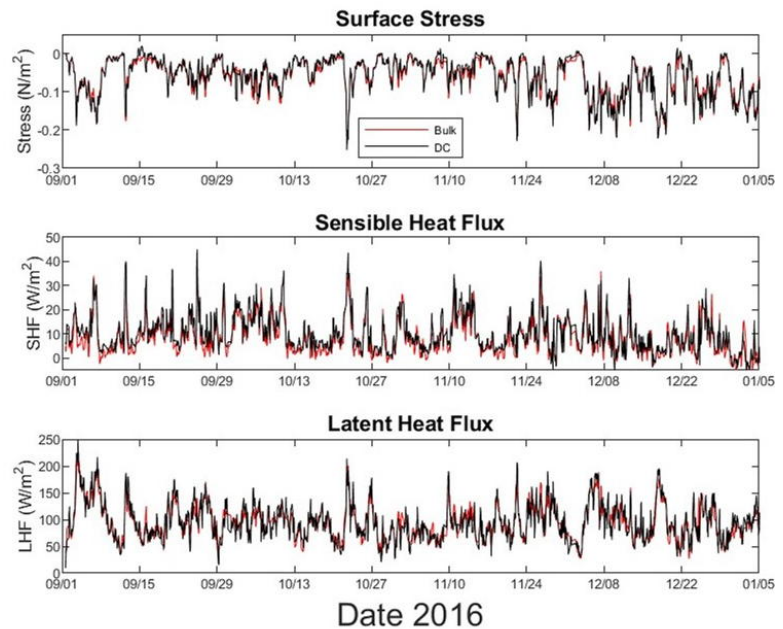
944 Observing mesoscale air-sea interaction processes is challenging since multiple oceanic and
945 atmospheric parameters must be measured with high accuracy and spatio-temporal resolution.
946 The past decade has seen the emergence of many novel in situ and remote sensing platforms that
947 increasingly better capture mesoscale and smaller processes with high accuracy and resolution
948 (e.g., Chapter 9 of Kessler et al. 2019). These novel observational technologies are expected to
949 provide opportunities for multi-platform, coordinated measurements for air-sea interaction
950 studies (e.g., Bony et al. 2017; Wang et al. 2018).

951

952 *a. In situ observations*

953 Oceanographic moorings can be equipped with meteorological instruments, including direct
954 covariance flux systems and bulk meteorological sensors, to provide directly measured and bulk-
955 estimated air-sea fluxes, respectively. An example system is shown in Figure 10 from the second
956 Salinity Processes in the Upper-ocean Regional Study (SPURS-2) experiment, which computed
957 and telemetered in near-real-time the motion-corrected surface wind stress and sensible and
958 latent heat fluxes from a surface mooring for the first time (Clayson et al. 2019). There is overall

959 a good qualitative agreement between the measured and estimated air-sea fluxes (Bigorre et al.
 960 2013). However, the bulk formula method underestimates the momentum flux and overestimates
 961 the buoyancy flux under high wind conditions. These biases are categorically related to
 962 deficiencies in formulations for the drag and heat transfer coefficients. Edson et al. (2013)
 963 revised the formulations for drag coefficient in COARE 3.5 to alleviate the low drag coefficient
 964 bias and proposed a new formula for heat transfer coefficients. Ayet and Chapron (2022)
 965 reviewed potential wave-atmospheric turbulence coupling mechanisms that allow for further
 966 refinements. Recently, buoy arrays have been deployed as part of the Ocean Observatories
 967 Initiative (OOI, Trowbridge et al. 2019) and operated for years on both coasts. These in situ data
 968 and the simultaneous measurements of surface meteorology and wave conditions are crucial to
 969 reducing the uncertainty in air-sea flux estimates in modern bulk formulas (Edson et al. 2013;
 970 Cronin et al. 2019; Villas-Boas et al. 2019).
 971



972
 973 **Figure 10:** (left) The SPURS-2 central mooring with instrumentation at the upper right includes a sonic
 974 anemometer, infrared hygrometer, and sensors to remove buoy motion. The sensor package can directly
 975 measure the surface stress, sensible heat, and latent heat fluxes (See Clayson et al. 2019 for more details on
 976 instrumentations). (right) Time series of these fluxes showing bulk estimates in red and direct covariance (DC)
 977 fluxes in black. A good qualitative agreement is seen between the bulk and DC estimates, with the most
 978 significant discrepancies visible in the sensible heat flux (Bigorre et al. 2013). The coincident measurements of
 979 direct flux and bulk meteorology from SPURS-2 and prior field campaigns (e.g., CBLAST, DYNAMO,
 980 CLIMODE, etc.) are being used for improving the bulk flux algorithm for turbulent heat flux transfer
 981 coefficients. Photo by James B. Edson (WHOI).

982 Autonomous surface vehicles (ASVs) are piloted wave- or wind-propelled surface platforms that
983 can be instrumented with ocean, atmospheric, and biogeochemical sensors. Widely-used ASVs
984 include Saildrones (Meinig et al. 2019) and Wave Gliders (Thomson and Girton 2017), which
985 have long-endurance (~6 months) and can sample in remote locations and be piloted across
986 fronts. Using numerous instruments can mitigate issues with cross-frontal sampling and thus
987 capture mesoscale and smaller variations in air-sea interaction (Quinn et al. 2021; Stevens et al.
988 2021).

989
990 Drifting platforms can be instrumented with various sensors that capture air-sea interaction. The
991 Global Drifter Program, a global network of surface drifters that typically measure currents, SST,
992 and barometric pressure, has contributed to understanding global mesoscale circulation
993 (Laurindo et al. 2017; Centurioni et al. 2019). Drifting spar buoys (Graber et al. 2000; Edson et
994 al. 2013) have been measuring surface fluxes in situ for decades. In recent years, sophisticated
995 low-profile Lagrangian platforms have been developed, such as SWIFTs (Thomson 2012), to
996 measure surface currents, waves, and near-surface ocean turbulence over various wave
997 conditions. Benefits of drifters include relatively low cost and Lagrangian sampling. However,
998 they tend to converge at fronts; therefore, multiple drifters are necessary to characterize cross-
999 frontal structure (D'Asaro et al. 2018).

1000
1001 Recent advancements in biologging technology may help facilitate autonomous measurements
1002 and real-time monitoring of essential ocean variables that may be important for air-
1003 sea interaction studies (Harcourt et al. 2019). As the biologging data can track mesoscale eddies
1004 and fronts in greater detail (Charrassin et al. 2008; Miyazawa et al. 2019) and can be assimilated
1005 into operational models (Yoda et al. 2014; Miyazawa et al. 2015), the application of animal-
1006 borne sensors has the potential to advance predictive capabilities of extratropical cyclones that
1007 strongly interact with the oceans (Section 3b).

1008
1009 Aircraft measurements are crucial for air-sea interaction studies. The platform's mobility is
1010 advantageous because of its ability to obtain in situ measurements of the horizontal and vertical
1011 variability in and above the MABL in a short time. With carefully designed flight patterns, it can
1012 also derive mesoscale forcing to the boundary layer using the velocity field measured at flight

1013 level (Lenschow et al. 1999; Stevens et al. 2003). In the past 20 years, air-deployable sensor
1014 packages such as GPS dropsondes, AXBT, AXCTD, and instrumented floats have further
1015 expanded the sampling capability to depict the entire column of the atmosphere and the upper
1016 ocean, particularly when low-level flights are not feasible (Doyle et al. 2017). In recent years,
1017 airborne measurements have been extended to 10 m above the sea surface using a controlled
1018 towed vehicle (Wang et al. 2018). This new capability is significant to air-sea interaction studies,
1019 particularly surface flux parameterization.

1020

1021 *b. Remote sensing*

1022 Emerging remote sensing platforms, including satellite, ground-based, or airborne
1023 measurements, present promising means to estimate air-sea fluxes at ocean mesoscale and
1024 smaller. Scatterometer and microwave measurements provide collocated global views of ocean
1025 vector winds and SST under all wind conditions at daily scales. However, considerable
1026 uncertainty exists under extreme conditions due to inconsistent in situ reference wind speeds
1027 from dropsondes and moored buoys to calibrate satellite winds (e.g., Polverari et al. 2021). This
1028 also implies uncertainties in modeling ocean drag and air-sea interaction. The virtual
1029 constellation of scatterometers (Stoffelen et al. 2019) provides good temporal coverage of the
1030 extremes, with now 7 scatterometers in space with revisits globally within 30 minutes or a few
1031 hours (Gade and Stoffelen 2019). Future satellite observations will need to resolve synoptic
1032 variability under strong wind and rain and increase the resolution of the vertical profiles within
1033 the MABL to better estimate the relationship between the surface flux and flux profiles.

1034

1035 For momentum fluxes, key variables are surface winds, currents, and waves. In coastal regions,
1036 high-frequency radar systems provide surface currents at $O(1)$ km resolution (Kim 2010; Paduan
1037 and Washburn 2013; Kirincich et al. 2019), which can be used to infer surface wave conditions
1038 and wind stress (e.g., Saviano et al. 2021). The airborne DopplerScatt system simultaneously
1039 captures surface wind stress, waves, and currents (Wineteer et al. 2020) and is central to the Sub-
1040 Mesoscale Ocean Dynamics Experiment (S-MODE; Farrar et al. 2020). Similar concepts for new
1041 satellite observations have been proposed (see Villas Bôas et al. 2019) and are currently in
1042 various development stages (e.g., Bourassa et al. 2016; López-Dekker et al. 2019;
1043 Gommenginger et al. 2019; Wineteer et al. 2020). Surface waves are crucial for accurate

1044 estimates of momentum flux; new satellite missions such as CFOSAT (Chinese-French
1045 Oceanography Satellite) simultaneously measuring waves and winds (Ardhuin et al. 2019) are
1046 expected to improve the accuracy of the wind-speed and wave-based formulations in the
1047 advanced bulk formula for air-sea flux. Satellite surface measurements of stress-equivalent winds
1048 more closely respond to stress than wind (e.g., de Kloe et al. 2017). Given the persistent large-
1049 scale and mesoscale errors in NWP reanalyses (Belmonte and Stoffelen 2019; Trindade et al.
1050 2020), these new satellite observations collocated with in situ measurements of surface stress
1051 will be valuable for understanding stress-related air-sea coupling and improving ocean modeling
1052 and marine forecasting (Bourassa et al. 2019).

1053
1054 In contrast to momentum flux, a critical gap remains in the current satellite remote sensing
1055 capability to provide accurate global estimates of turbulent heat and moisture fluxes. Current
1056 satellite remote sensing systems rely on bulk parameterizations to estimate net heat and gas
1057 fluxes (Cronin et al. 2019). Mesoscale air-sea interaction studies will benefit significantly from a
1058 satellite mission that measures co-located, small-scale state variables, including near-surface
1059 atmospheric temperature and humidity, SST, and wind speed, that allow accurate estimates of the
1060 turbulent heat fluxes (e.g., Gentemann et al. 2020). This will also help validate the numerical
1061 models to lower the uncertainty in air-sea heat flux and improve related predictions.

1062
1063

1064 **6. Discussion and synthesis**

1065 Since the first global-scale surveys of the mesoscale air-sea interactions by Chelton et al. (2004)
1066 and Xie (2004), our theoretical understanding and observational and modeling capabilities in the
1067 past two decades have advanced significantly, leading to a substantial body of literature related
1068 to ocean mesoscale air-sea interaction. Our current scientific understanding indicates that
1069 mesoscale eddies perturb the MABL via surface flux anomalies, leading to dynamic and
1070 thermodynamic adjustments (Section 2; Figure 1d). The MABL response is communicated to the
1071 free troposphere, especially over WBCs (Figure 1b,c), influencing downstream development of
1072 weather and short-term climate events (Section 3; Figure 1a,b). The MABL response feeds back
1073 to the ocean circulation, modifying WBC dynamics, air-sea gas exchanges, and nutrient
1074 distribution (Section 4; Figure 1e,f). This new knowledge has transformed our classical

1075 understanding of physical processes, leading to notable revisions of oceanic and atmospheric
1076 circulation theories that incorporate the coupled effects of ocean mesoscale processes, wave, and
1077 biogeochemical processes (Section 4). Our observing capability has advanced rapidly to
1078 characterize mesoscale air-sea interaction (Section 5). However, numerous challenges and open
1079 questions remain. The remainder of the chapter will focus on physical and biological aspects of
1080 modeling, observational, and diagnostic approaches that require further research in the coming
1081 years.

1082

1083 *a. Attribution of near-surface wind convergence*

1084 While the WBC SST impact on the MABL dynamics is increasingly better understood, there are
1085 some critical remaining questions regarding the essential role of WBC SST forcing on the time-
1086 mean atmospheric state. The ongoing debates about the origin of the near-surface wind
1087 convergence (NSWC) and the maximum precipitation over WBCs are particularly relevant as
1088 they entail important implications pertinent to various aspects of the topics discussed in this
1089 article. That is, assessing whether the steady linear boundary layer dynamics account for the
1090 time-mean NSWC and vertical motion requires a detailed understanding of the modulation of
1091 boundary layer ageostrophic circulation by SST (Section 2; Figure 1d). On the other hand, the
1092 demonstrated impacts of storms and atmospheric fronts on the NSWC require a careful
1093 examination of extratropical cyclogenesis modulated by the diabatic forcing over the ocean
1094 fronts (Section 3; Figure 1b-c). Overall, any approach to quantifying the nature of the
1095 relationships between NSWC and SST will need to robustly separate the small magnitude
1096 convergence predicted by linear boundary layer theory from the large anomalous convergence
1097 induced by storm systems that are several orders of magnitude greater.

1098

1099 *b. Robust diagnostic framework*

1100 The debate about the role of SST fronts in the NSWC arises partly due to the lack of a robust
1101 process-based diagnostics and analytic framework to interpret the observed convergence
1102 patterns. The existing analytical model of Schneider and Qiu (2015) discussed in Section 2c
1103 offers a complete account of the role of boundary layer dynamics over the SST fronts, providing
1104 the two limiting cases of wind response to SST dependent on background wind speed. The model
1105 also suggests an extension of the diagnostic framework from the widely used coupling

1106 coefficients to lagged regression, impulse response, or corresponding spectral transfer functions.
1107 Yet, the model assumes a quasi-steady state and does not account for the stochastic and moist
1108 processes associated with the storm tracks and their synoptic-scale influence on NSW. A
1109 critical path forward is to incorporate the time-dependent and moist processes related to
1110 extratropical storms along SST frontal zones and the local SST-induced boundary layer response
1111 in a single analytical framework. Given the coexistence of the SST and current feedback effects
1112 along the frontal zones, any future development of diagnostic frameworks will also have to
1113 consider the mechanical coupling effects simultaneously along with the thermal effects (e.g.,
1114 Takatama and Schneider 2017; Seo 2017; Renault et al. 2019a).

1115

1116 *c. Large-scale impacts in climate models*

1117 Numerous studies have demonstrated WBC impacts on downstream atmospheric circulation
1118 (Figure 1e). Some studies argue that the sharpness of WBC fronts shifts the storm track and jet
1119 stream, influencing the blocking frequency in Europe and Northeastern Pacific (e.g., Kuwano-
1120 Yoshida and Minobe 2017; O'Reilly et al. 2015, 2016, 2017; Piazza et al. 2016). Other studies
1121 find that meridional shifts of WBC fronts alter the atmospheric transient eddy heat flux
1122 downstream (e.g., Frankignoul et al. 2011; Kwon and Joyce 2013; Seo et al. 2017; Joyce et al.
1123 2018). Warm-core eddies near the KOE act as significant oceanic sources of moisture and heat
1124 for large-scale circulation, altering downstream precipitation patterns (Ma et al. 2015, 2016; Liu
1125 et al. 2021). The importance of the seasonal background state in the atmosphere has also been
1126 recognized as it shapes the atmospheric response to SSTA (e.g., Taguchi et al. 2009; Huang et al.
1127 2020).

1128

1129 However, some aspects of the far-field circulation response and its statistical significance remain
1130 elusive (Kushnir et al. 2002; Kwon et al. 2010; Czaja et al. 2019). Deriving a robust conclusion
1131 on downstream influences is particularly challenging difficult because the studies adopt different
1132 methods to define WBC SST impacts, leading to distinct amplitudes/patterns of SST
1133 perturbations and atmospheric responses. This uncertainty is in addition to differences in model
1134 climatologies. To date, the relative impacts of sharpness of SST gradient, its meridional shift,
1135 and activity of warm or cold-core eddies remain unquantified (Parfitt and Seo 2018). The

1136 importance of the coordinated modeling and diagnostic approaches regarding this specific point
1137 is emphasized in Section 6d.

1138

1139 *d. Coordinated climate modeling and improved physical parameterizations*

1140 Significant progress can be made in understanding results and uncertainties in climate models of
1141 different complexity and resolutions via coordinated modeling experiments with resolutions at or
1142 beyond the ocean mesoscale and shared sets of diagnostics. The CMIP6 HighResMIP protocol
1143 (Haarsma et al. 2016) and PRIMAVERA project (Bellucci et al. 2021) well represent the
1144 community's interests in this direction. Analyses from a subset of these models reveal significant
1145 model resolution sensitivity (especially in the oceans) of the simulated air-sea interaction and
1146 climate regimes in the extratropics (e.g., Jullien et al. 2020; Moreton et al. 2021). Further
1147 advances in model resolution, for example, DYAMOND (Stevens et al. 2019) and the planned
1148 HighResMIP2, together with programs such as OASIS (Observing Air-Sea Interaction Strategy,
1149 <https://airseaobs.org>; Cronin et al. 2022) that aims to bring observations and models closer
1150 together, will build on these previous efforts and provide further insights into the fidelity of
1151 modeled mesoscale air-sea interactions. Furthermore, in the ocean and coupled models where the
1152 ocean eddies are not fully or only partially resolved, their rectified effects on the air-sea heat,
1153 momentum, and tracer fluxes are not currently parameterized. Various stochastic representations
1154 of eddy transports are being tested and implemented (Section 4c), which can potentially address
1155 this issue of low-frequency rectification effects by eddies on large-scale climate via air-sea
1156 interaction. (e.g., Siqueira and Kirtman 2016).

1157

1158 *e. Air-sea interaction mediated by ocean submesoscale and sea state*

1159 The ocean submesoscale processes with length-scales smaller than ~ 10 km are essential for the
1160 ocean energy cycle (Lorenz 1960), global heat balance (Su et al. 2018), and marine
1161 biogeochemistry and ecosystems (Omand et al. 2015; Lévy et al. 2018). While the dynamics of
1162 the submesoscale ocean instabilities are becoming better understood (e.g., Fox-Kemper et al.
1163 2008; D'Asaro et al. 2011), their direct impact on the MABL and heat and carbon uptake by the
1164 oceans (e.g., Johnson et al. 2016; Bachman et al. 2017; du Plessis et al. 2019) remain poorly
1165 understood. Thus far, only a few satellite-based studies provide direct observational evidence of
1166 relative wind stress response to submesoscale SST fronts (e.g., Beal et al. 1997; Xie et al. 2010;

1167 Gaube et al. 2019; Ayet et al. 2021), although prior in situ observational studies have long
1168 documented such interactions in localized regions (e.g., Sweet et al. 1981; Friehe et al. 1991;
1169 Mahrt et al. 2004). While results from high-resolution numerical simulations (e.g., LES) indicate
1170 submesoscale SST-driven MABL dynamics (Skylingstad et al. 2007; Lambaert et al. 2013;
1171 Wenegrat and Arthur 2018; Lac et al. 2018; Sullivan et al. 2020, 2021), they also recognize the
1172 importance of advection and convective organization in characterizing the nonlinear MABL
1173 dynamics that co-occur at the submesoscale. As for the oceanic impact, the ocean current
1174 feedback dominates the wind stress response at the submesoscale, influencing the kinetic energy
1175 cascade (Renault et al. 2018). Spatial variability in sea state and surface roughness is enhanced at
1176 the submesoscale, and hence wave-current interactions (e.g., Villas Bôas and Pizzo 2021) and
1177 wave-wind interactions (e.g., Deskos et al. 2021) are expected to be critical in determining wind
1178 stress, heat flux, and MABL variations (Ayet et al. 2021; Section 4b), yet such processes remain
1179 poorly observed, understood, and parameterized. Emerging in situ and satellite observations for
1180 near-surface processes (Section 5), combined with dedicated atmospheric and oceanic LES and
1181 high-resolution modeling studies, will help improve the physical understanding of air-sea
1182 interactions at the submesoscale.

1183

1184 *f. Air-sea gas flux exchange and ocean biogeochemistry processes*

1185 Estimates of air-sea gas exchange do not fully consider the effects of ocean mesoscale eddies and
1186 fronts. One issue is that the gas transfer velocity typically does not consider wind variations
1187 introduced by mesoscale air-sea interactions. The transfer velocity is also often based on wind
1188 speed (e.g., Wanninkhof 1992). Hence, it only implicitly accounts for the sea state variations.
1189 Studies with parameterizations that consider bubble-mediated gas exchanges due to breaking
1190 waves (e.g., Frew et al. 2007; Deike and Melville 2018) reveal their significant contribution to
1191 regionally-integrated CO₂ flux, especially under midlatitude storm tracks (e.g., Reichl and Deike
1192 2020). To accurately represent the sea state influence modulated by mesoscale processes in the
1193 transfer velocity-based flux parameterization (e.g., Fairall et al. 2011; Edson et al. 2011), it is
1194 imperative to increase direct measurements of CO₂ flux (e.g., McGillis et al. 2001) along with
1195 the coincident observations of wind, waves, solubility, and air-sea partial CO₂ pressure
1196 differences.

1197

1198 Further, mesoscale air-sea interaction feeds back to ocean primary productivity (Lévy 2008;
1199 McGillicuddy et al. 2016) and tracer concentrations, such as carbon. Since the physical
1200 properties of mesoscale eddies and their relationships with biogeochemical variables vary widely
1201 by region (e.g., Chelton et al. 2011; Gaube et al. 2013, 2014; Frenger et al. 2018), future work
1202 should aim to identify the specific aspects of this regional variability that are due to mesoscale
1203 air-sea interaction and subsequent impacts on upwelling and vertical mixing. Eddy-rich climate
1204 model simulations are one avenue to gain quantitative insight into the relevance of the complex
1205 coupling of ocean mesoscale features, biogeochemistry, and the atmosphere. Few such
1206 simulations exist due to their computational expense (e.g., Harrison et al. 2018), but we expect
1207 this to change in the coming years. Dedicated field experiments combined with eddy-resolving
1208 coupled physical-biogeochemical models are critical to determining what aspects of mesoscale
1209 air-sea interactions need to be considered and represented in non-eddy-resolving models.

1210

1211 *g. Final remarks*

1212 Prospects for significant advances in mesoscale air-sea interaction in the coming years are
1213 incredibly bright. Strong community efforts and enthusiasm exist for building sustained
1214 observational networks to characterize detailed physical and biogeochemical processes across the
1215 air-sea coupled boundary layers (e.g., OceanObs'19 White Papers; OASIS; US CLIVAR's air-
1216 sea interaction research initiatives). New satellite missions with advanced instrument technology
1217 and retrieval algorithms will continue to improve our capability to monitor state variables
1218 pertinent to air-sea interactions at fine scales and with increased accuracy. These new
1219 observations will lead to updated physical parameterizations that are becoming increasingly
1220 more scale-aware and that can be potentially built with stochastic schemes that account for
1221 rectified effects of eddy transports on air-sea flux and large scales. More field experiments are
1222 being coordinated via close integration with process-oriented and data assimilative modeling to
1223 help not only develop the sampling plans but also improve the parameterizations and skills in
1224 prediction models (e.g., Cronin et al. 2009; Cravatte et al. 2016; Kessler et al. 2019; Sprintall et
1225 al. 2020; Shroyer et al. 2021; Shinoda et al. 2021; Newman et al. 2022). The climate modeling
1226 community is developing and refining high-resolution Earth system model simulations with
1227 advanced physical parametrizations. International partnership and coordination are becoming
1228 increasingly solid, enabling the design of multi-model, multi-ensemble, high-resolution coupled

1229 modeling protocols and diagnostic frameworks. The identified common biases in mesoscale air-
1230 sea interaction in such climate models, in turn, guide the sampling strategy of observing systems
1231 and process studies. Ensemble data assimilation systems are rapidly advancing, yielding more
1232 accurate observationally constrained ocean, atmosphere, and biogeochemical state estimates
1233 critical for sub-seasonal to decadal predictions (e.g., Penny and Hamill 2017; Verdy and Mazloff
1234 2017). Overall, the successful coordination across observations, modeling, and theories has been
1235 critical, and these coordinated efforts will and should continue to enhance Earth system
1236 prediction skills across scales from weather forecasts to climate projection scales.

1237

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1264

1265 ***Data Availability Statement***

1266 Datasets used in the figures are based on ERA5 (Hersbach et al. 2020), NOAA OI SST
1267 (Reynolds et al. 2007), climate model simulations from the HighResMIP (Haarsma et al. 2016),
1268 or already published papers as cited in the figure captions.

1269

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