

1 **Improving wave-based air-sea momentum flux**
2 **parameterization in mixed seas**

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5 **Key Points:**

- 6 • Surface stress at moderate to high winds is dominated by short wind waves.
7 • COARE3.5 wave based formulation can underestimate surface stress by more
8 than 10 % in mixed sea conditions under moderate to high wind.
9 • Using the mean wave period or including the directional alignment between
10 wind and wave in COARE3.5 alleviates this issue.

Abstract

In winter, the Northwest Tropical Atlantic Ocean can be characterized by various wave age-based interactions among ocean current, surface wind and surface waves, which are critical for accurately describing surface wind stress. In this work, coupled wave-ocean-atmosphere model simulations are conducted using two different wave roughness parameterizations within COARE3.5, including one that relies solely on wind speed and another that uses wave age and wave slope as inputs. Comparisons with the directly measured momentum fluxes during the ATOMIC/EUREC⁴A experiments in winter 2020 show that, for sea states dominated by short wind waves under moderate to strong winds, the wave-based formulation increases the surface roughness length in average by 25% compared to the wind-speed-based approach. For sea states dominated by remotely generated swells under moderate to strong wind intensity, the wave-based formulation predicts significantly lower roughness length and surface stress ($\approx 15\%$), resulting in increased near-surface wind speed above the constant flux layer ($\approx 5\%$). Further investigation of the mixed sea states in the model and data indicates that the impact of swell on wind stress is over-emphasized in the COARE3.5 wave-based formulation, especially under moderate wind regimes. Various approaches are explored to alleviate this deficiency by either introducing directional alignment between wind and waves or using the mean wave period instead of the wave period corresponding to the spectral peak to compute the wave age. The findings of this study are likely to be site-dependent, and mostly concern specific regimes of wind and waves where the original parameterization was deficient.

Plain Language Summary

Accurately understanding and describing air-sea interactions is critical for weather forecast and regional climate. In this work, we use numerical experiments with and without taking into account the ocean waves to describe air-sea interactions. Most of the momentum exchange between the ocean and the atmosphere is done through locally wind-generated waves, however remotely generated waves, such as swells, can also interfere in these air-sea interactions. Comparisons with observations made during the ATOMIC/EUREC⁴A field campaigns in winter 2020 show in particular that our numerical experiment overestimated the impact of the swell on the atmosphere. Various approaches are explored here to alleviate this deficiency, one of those being the introduction of the effect of the alignment between wind and waves.

1 Introduction

Over the ocean, most of the momentum, heat, and mass exchanges with the atmosphere are supported by short wind-waves on spatial scales of $O(0.1-10\text{m})$. These wind-waves enhance the surface drag and roughness at the air-sea interface, thereby increasing the wind stress. The wind stress is coupled with the planetary boundary layer (PBL) processes in the atmosphere, modifying the kinematic and thermodynamic profiles in this lowest part of the atmosphere (Janssen, 1989; Moon et al., 2004). In addition to locally generated wind-waves, the sea state is also influenced by the remotely generated swell, especially in the lower latitudes, whose propagation direction is often uncorrelated with local winds. The fast-propagating swell wave that is strongly misaligned with or outruns the local wind can be a conduit for upward momentum and energy transfer from waves to the wind, forming a wave-driven low-level jet (e.g., Harris, 1966; Sullivan et al., 2008; Hanley & Belcher, 2008) and dissipating the swell waves (M. Donelan, 1999; Kahma et al., 2016; Liu et al., 2017).

60 In numerical models, the wind stress over the oceans is parameterized using
61 bulk flux algorithms, such as the Coupled Ocean-Atmosphere Response Experi-
62 ment (COARE, Fairall et al., 1996, 2003; Edson et al., 2013). If no coincident wave
63 fields are available, COARE parameterizes the wave roughness length (z_0) using
64 wind speeds only. In this study, this approach will be referred to as the wind-speed-
65 dependent formulation (WSDF). Since wind and wind-waves are in near-equilibrium
66 in many cases over the extratropical open oceans, the COARE’s WSDF tends to
67 accurately predict the surface roughness and thereby the surface stress (Edson et al.,
68 2013). However, under trade-wind regimes in the tropics such as our study region
69 in boreal winter, remotely-generated swell significantly shape the sea state, whose
70 effect on wind stress cannot be accurately characterized by local wind alone. To
71 improve estimates of the fluxes under these conditions, “wave-based” formulations
72 exist in many bulk flux algorithms that model z_0 as a function of wave age or wave
73 age/slope (e.g., Taylor & Yelland, 2001; Oost et al., 2002; Drennan et al., 2003;
74 Edson et al., 2013; Sauvage et al., 2020). As there are increasing interests and op-
75 portunities to incorporate the wave effects on surface fluxes in numerical models,
76 such wave-based formulations (WBF) in bulk formulas will likely be adopted more
77 in such models. Since the parameterized surface fluxes serve as lower boundary con-
78 ditions for turbulent exchanges within the atmospheric and oceanic boundary layers,
79 the simulation and forecast skills will be influenced by the physics and assump-
80 tions represented in the bulk formulas. Therefore, it is imperative to understand
81 the assumptions and deficiencies in current WBFs and offer possible revisions to the
82 formulations for air-sea fluxes with increased accuracy. The goal of this paper is to
83 enhance a regime-based understanding of wave-wind interactions via detailed valida-
84 tion of the parameterized air-sea flux from high-resolution coupled model simulations
85 against directly measured air-sea fluxes.

86 This study focuses on air-sea momentum flux during the ATOMIC/EUREC⁴A
87 field campaign. The ATOMIC (Atlantic Tradewind Ocean-Atmosphere Mesoscale
88 Interaction Campaign) is the U.S. complement to the European field campaign,
89 EUREC⁴A (Elucidating the Role of Cloud-Circulation Coupling in Climate,
90 Stevens et al., 2021), both of which took place in the Northwest Tropical Atlantic
91 Ocean in January-February 2020 (Figure 1). The primary objective of this study is
92 to determine how well the current WBF in an advanced bulk flux algorithm such as
93 COARE3.5 reproduces the observed wind stress in the mixed sea conditions com-
94 pared to the WSDF. By exploiting the fully-coupled ocean-atmosphere-wave model
95 simulations and extensive analyses of the in situ observational datasets, we will at-
96 tempt to explain the causes for discrepancies between simulated and measured wind
97 stresses. Our results indicate that the current COARE3.5 WBF underestimates z_0
98 and wind stress, particularly over the mixed sea state. We will show that this is due
99 to either a missing physics of the wave-wind interaction or using an inappropriate
100 wave input parameter to describe the mixed sea condition.

101 The paper is organized as follows. Section 2.1 describes the technical details
102 of the latest z_0 formulation in COARE3.5. Sections 2.2 and 2.3 discuss the fully
103 coupled ocean-atmosphere-wave modeling system used in the investigation, followed
104 by the details on the experimental design and observational datasets in Section 2.4
105 and Section 2.5, respectively. The wave impact on z_0 , wind stress, and low-level
106 winds are discussed in a case study investigation in Section 3. Section 4 provides an
107 in-depth comparison of the parameterized momentum flux against the direct mea-
108 surements, identifying the areas and regimes for further improvement. In section
109 5, possible approaches are proposed and tested to alleviate the biases. Section 6
110 provides a summary and discussion.

2 Air-sea flux parameterization and coupled model

This section provides a brief overview of the wave-mediated momentum flux implemented in the Coupled Ocean-Atmosphere Response Experiment parameterization (COARE3.5, Fairall et al., 1996, 2003; Edson et al., 2013). Hereafter, we will focus on the COARE3.5 version, although a slightly updated version, COARE3.6, has been made publicly available. However, the findings of this study would stay unchanged when using COARE3.6 (not shown).

2.1 Roughness length and momentum flux in COARE3.5

The along wind stress in the COARE framework is defined as:

$$\tau = \rho C_D(z, z_0, \psi_m) U_r(z) S_r(z) = \rho u_*^2, \quad (1)$$

where ρ_a is the air density, $U_r(z)$ is the magnitude of the along-wind component of the wind vector, $S_r(z)$ is the scalar wind speed, where the subscript r denotes relative to the ocean surface; and u_* the friction velocity. C_D is the drag coefficient defined as:

$$C_D(z, z_0, \psi_m) = \left[\frac{\kappa}{\ln(z/z_0) - \psi_m(\zeta)} \right]^2, \quad (2)$$

where κ is the von Kármán constant, $\psi_m(\zeta)$ is an empirical function of atmospheric stability, ζ is the z/L ratio with L the Obukhov length and z the height above the surface (Fairall et al., 1996). The surface roughness length z_0 is parameterized in COARE3.5 as the sum of two terms:

$$z_0 = z_0^{smooth} + z_0^{rough}, \quad (3)$$

where z_0^{smooth} and z_0^{rough} represent the smooth and rough flow components of z_0 , respectively (Edson et al., 2013). The smooth flow component is parameterized as

$$z_0^{smooth} = \gamma \frac{\nu}{u_*}, \quad (4)$$

where γ is the roughness Reynolds number for smooth flow, set to be constant at 0.11 based on laboratory experiments, and ν is the kinematic viscosity. For smooth flow, the wind stress is mainly supported by viscous stress where $z_0 \approx z_0^{smooth}$.

The rough part of the roughness length, z_0^{rough} , is meant to parameterize the wind-driven gravity waves that support most of the stress above approximately 5 m s^{-1} when the sea becomes aerodynamically rough. This component of the roughness is formulated currently in several ways in COARE3.5. The simplest and the most broadly used way is to parameterize it as a function of wind speed only. The so-called wind speed dependent formulation without explicit wave and sea states inputs estimates z_0^{rough} using the Charnock's relation (Charnock, 1955):

$$z_0^{rough} = \frac{\alpha_{ch} u_*^2}{g}, \quad (5)$$

where g is the acceleration of gravity and α_{CH} is the Charnock coefficient that is dependent only on wind speed. COARE3.5 formulates α_{CH} as

$$\alpha_{ch} = mU_{r10N} + b, \quad (6)$$

144 where U_{r10N} is the 10-m wind speed relative to the sea surface under neu-
 145 tral conditions (Edson et al., 2013, Appendix) and coefficients $m = 0.0017$ and $b =$
 146 -0.005 (?). Hereafter, U_{r10N} is defined such as:

$$U_{r10N} = \frac{u_*}{\kappa} \ln(10/z_0), \quad (7)$$

147 The coefficients m , and b in Eq. 6, have been determined to fit the average
 148 data used in COARE3.5 over wind speeds between 5 and 18 ms^{-1} . If wind speed
 149 is below 5 ms^{-1} , the surface roughness is mainly determined by z_{smooth} in Eq. 4.
 150 For wind speeds greater than 18 ms^{-1} , COARE3.5 fixes the value of the Charnock
 151 coefficient to its value at 18 ms^{-1} . Note, however, that although α_{CH} is fixed above
 152 18 ms^{-1} , z_0^{rough} , C_D and τ all continue to increase with the wind speed, just at a
 153 lower rate.

154 An alternative way to define z_0^{rough} in COARE3.5 is to use the so-called wave-
 155 based formulation (WBF), which requires contemporary information about the wave
 156 field and its state of development, such as significant wave height (H_s) and phase
 157 speed of the waves at the peak of the spectrum (c_p). Two WBFs are currently avail-
 158 able in COARE3.5, one that uses the wave age only and another that uses both the
 159 wave age and wave steepness. In the second form, which is explored in this study in
 160 great detail, z_0^{rough} is expressed as

$$z_0^{rough} = H_s D \left(\frac{u_*}{c_p} \right)^B, \quad (8)$$

161 where u_*/c_p is the inverse wave age based on the friction velocity, and D and
 162 B are numerical constants given by $D = 0.09$ and $B = 2$ in Edson et al. (2013).
 163 Hereafter, we will use a definition of wave age based on the ratio of the phase speed
 164 of the waves at the spectral peak over the surface wind speed at 10 m defined as

$$\chi = \frac{c_p}{U_{10}}. \quad (9)$$

165 The wave age is used to describe the state of development of the wave field.
 166 For example, a wave age close to 1.2 represents a fully developed sea when the sur-
 167 face waves and stress are largely in equilibrium (e.g., Phillips, 1985), in which the
 168 rate that wind does work on the surface waves is balanced by the dissipation rate of
 169 breaking waves (microbreakers and whitecaps) and nonlinear wave-wave interactions
 170 (e.g., Csanady & Gibson, 2001). Wave ages under 1 are associated with developing
 171 seas and young waves, while wave ages well above 1.2 describe decaying seas and
 172 swell. It should be noted that in the current COARE3.5, c_p is defined using the peak
 173 period of the waves, T_p , in deep water such that:

$$c_p = g \frac{T_p}{2\pi}. \quad (10)$$

174 In Section 3, we will examine the sensitivity of the estimated momentum flux
 175 based on the current COARE3.5 algorithm. Guided by comparison to the observa-
 176 tions in Section 4, we will then explore the impacts of revised COARE3.5 WBF in
 177 Section 5.

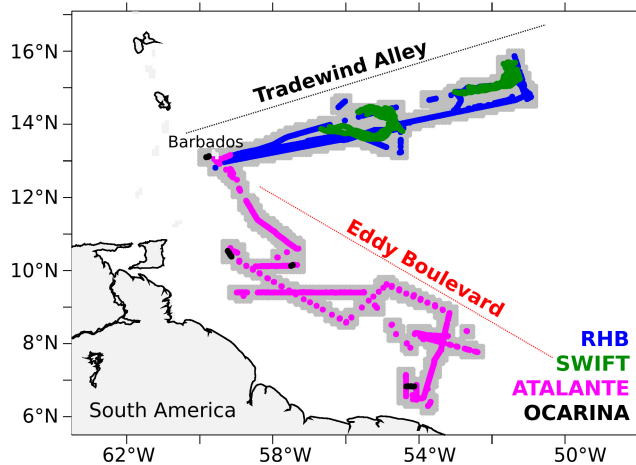


Figure 1. Tracks of the different platforms measuring surface stress. The gray area denotes where the model outputs are sampled along the tracks of observations. RHB provided data from January 9 to February 13, 2020. SWIFT drifters were deployed from 14 January to 22 January 2020 and from 30 January to 11 February 2020. R/V ATALANTE provided data from January 19 to February 19, 2020 and Ocarina was deployed periodically from January 25 to February 17, 2020.

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2.2 SCOAR regional coupled model system

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We use the Scripps Coupled Ocean-Atmosphere Regional (SCOAR) model (Seo et al., 2007, 2021), which couples the Weather Research and Forecast (WRF, Skamarock et al., 2008) Model to the Regional Ocean Modeling System (ROMS, Shchepetkin & McWilliams, 2005) via the COARE3.5 bulk flux algorithm (Fairall et al., 1996, 2003; Edson et al., 2013). In the absence of wave coupling, ROMS is driven by the surface heat flux (Q_{NET}), momentum flux (τ), and freshwater flux (Q_{FW}) computed from the wind speed-only formulation in COARE3.5 implemented in WRF. In turn, ROMS inputs SST and surface current vectors (U_s) to the COARE3.5 to compute the surface fluxes (Figure 2).

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2.3 Wave coupling in SCOAR

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This study implemented the coupling of the third-generation spectral wave model WaveWatch-III (WW3 Tolman et al., 2002; The WAVEWATCH III Development Group, 2016) into the SCOAR. Currently, two different ways are implemented to allow coupling waves to the atmosphere. The first option described in Figure 2 is based on the total friction velocity output from WW3 and used to estimate the wind stress and the resulting surface roughness length for computing turbulent heat fluxes. This option won't be used in this study. The second and third options described in Figure 2 are the focus of this manuscript and respectively take advantage of the COARE's WBF from (Edson et al., 2013), and the finding of this study. In this configuration, the centerpiece of the model coupling is the COARE3.5 implemented in the surface layer scheme in WRF to compute the air-sea fluxes. In this study, we use the Mellor-Yamada-Nakanishi-Niino (MYNN) surface layer scheme (Nakanishi & Niino, 2009; Jiménez et al., 2012), which over the ocean grid points computes the surface fluxes using the COARE3.5 WBF. WW3 is forced by the surface wind (U_{10}) from WRF and ocean current (U_s) from ROMS. WW3 then returns the significant wave height (H_s) and the phase speed of the dominant waves (c_p)

205 determined based on T_p (Eq. 10) to the MYNN surface layer scheme. In lieu of c_p ,
 206 WW3 can alternatively send the mean phase speed (c_m) and peak wave direction
 207 (Section 5). Spatially varying Charnock coefficients (α_{CH}) are then updated to pa-
 208 rameterize the surface roughness length (z_0) as a function of dominant wave age
 209 (χ) and wave steepness (Eq. 8). For this to work in WRF, the MYNN surface layer
 210 scheme has been modified to allow ingestion of wave age and significant wave height
 211 (H_s) from WW3. The MYNN PBL scheme (Nakanishi & Niino, 2004, 2006) is cou-
 212 pled to this modified surface layer scheme, allowing for the adjusted z_0 , wind stress
 213 (τ), and latent (Q_{LH}) and sensible (Q_{SH}) heat fluxes to influence the kinematic
 214 and thermodynamics processes in the PBL. The surface layer scheme has also been
 215 modified to take the ocean surface currents (U_s) from ROMS to compute the rela-
 216 tive wind and thus represent wind-current interaction. This so-called relative wind
 217 effect is represented in all simulations analyzed here. Wave to ocean coupling is also
 218 made available and ROMS can be forced by wave fields such as H_s and wave energy
 219 (FOC) fields. Wave-supported stress (τ^w) and wave dissipation (τ^{ds}) terms can also
 220 be send to ROMS to compute the ocean-side stress (τ^{oc}). For the purpose of this
 221 study, wave to ocean coupling is not included and thus on Figure 2 it is assumed
 222 that $\tau^{oc} = \tau^a$, where τ^a is the air-side stress.

2.4 Experiments

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 224 In WRF, the deep cumulus convection is represented through the Multi-scale
 225 Kain-Fritsch scheme (Zheng et al., 2016), the cloud micro-physics by the WRF
 226 single-moment 6-class scheme (Hong & Lim, 2006). The Goddard radiation scheme
 227 (Chou & Suarez, 1999) is used for shortwave and longwave radiation. The land
 228 surface process is treated with the Noah land surface model (F. Chen & Dudhia,
 229 2001). In ROMS, the KPP (K profile parameterization) scheme (Large et al., 1994)
 230 determines vertical eddy viscosity and diffusivity. The vertical grid in ROMS is
 231 stretched to enhance the resolutions near the surface and the bottom, using the so-
 232 called stretching parameters of $\theta_s = 7.0$, $\theta_b = 2.0$, and $h_{cline} = 300$ m. In WW3, the
 233 set of parameterizations from Ardhuin et al. (2010) is used, including swell dissipa-
 234 tion scheme (Ardhuin et al., 2009). Nonlinear wave-wave interactions are computed
 235 using the discrete interaction approximation (Hasselmann et al., 1985). Reflection
 236 by shorelines are enabled through Ardhuin and Roland (2012) scheme. The depth-
 237 induced breaking is based on Battjes and Janssen (1978), and the bottom friction
 238 formulation follows Ardhuin et al. (2003).

239 The model domain covers the Northwest Tropical Atlantic Ocean (Figure 3).
 240 The horizontal resolutions in WRF, ROMS, and WW3 are identical 10 km, with
 241 matching grids and land-sea masks. This horizontal resolution allows us to have
 242 reasonable description of the mixed sea state influenced by the remotely-generated
 243 swell and trade winds in the open oceans, which is the focus of this work. However,
 244 much finer-scale wind-wave and wave-current interactions, as studied in (Ardhuin et
 245 al., 2017; Bôas et al., 2020; Iyer et al., 2022), are not likely captured at this resolu-
 246 tion, especially in the regions of strong currents and eddy variability. ROMS (WRF)
 247 is run with a stretched vertical grid with a total of 30 (33) vertical levels, with ap-
 248 proximately 10 layers in the upper 150 m (below 1300 m). The model coupling is
 249 activated every 3 hours to account for the diurnal cycle.

250 A set of coupled model simulations presented in Section 4 is run for 6 months
 251 (November 1, 2019 to May 1, 2020), covering the ATOMIC/EUREC⁴A period, with
 252 a specific aim to compare with the measurements. In these simulations, the WRF
 253 model is initialized and driven by 3-hourly ERA5 global reanalysis at 0.25° resolu-
 254 tion (Hersbach et al., 2018a, 2018b), ROMS by the daily MERCATOR International
 255 global reanalysis at $1/12^\circ$ resolution (Lellouche et al., 2018), and WW3 by seven
 256 spectral points obtained from the global $1/2^\circ$ resolution WW3 simulations (Rasche

257 & Ardhuin, 2013). The initial conditions for ROMS and WW3 were obtained from
 258 the respective ROMS-only and WW3-only spin-up simulations forced by ERA5 at-
 259 mospheric forcing (starting from January 1, 2019). In ROMS, the tidal forcing is
 260 obtained using the Oregon State University Tidal Prediction Software (Egbert &
 261 Erofeeva, 2002) and applied as a 2-D open boundary condition by prescribing the
 262 tidal period, elevation amplitude, current phase angle, current inclination angle, the
 263 minimum and maximum tidal current, and ellipse semi-minor axes for 13 major tidal
 264 constituents. Daily climatology estimates of the Amazon and River and Orinoco
 265 River discharges are obtained from the Observatory Service SO-HyBAM database
 266 (<https://hybam.obs-mip.fr/>), which are prescribed as point sources close to the river
 267 mouths in our grid.

268 The second set of simulations presented in Section 3 is identical to that of the
 269 6-month-long simulations, except that WRF, ROMS, and WW3 are initialized from
 270 respectively 3-hourly ERA5 global reanalysis for the atmosphere and ROMS-only
 271 and WW3-only spin-up simulations for the ocean and waves as described above
 272 and run on a particular day (January 8, 2020) as a case study investigation. The
 273 motivation for the short simulations with the identical initial condition is to isolate
 274 the immediate impacts on z_0 and τ before the coupled feedback begins to alter the
 275 state variables. One could use the identical input state variables to estimate the
 276 air-sea fluxes offline using different COARE formulations. This yields similar results
 277 (not shown), indicating that the difference we show in Section 3 is not due to the
 278 difference in state variables, but due to the formulation difference. One notable ad-
 279 vantage to use the fully coupled model simulation is that it allows for evaluating the
 280 wind response beyond the surface layer (e.g., Figure 6c), and potentially large-scale
 281 feedback effects via the coupling.

Table 1. Summary of the different SCOAR experiments.

Experiments	z_0 parameterization	Relative wind	Wave period	misaligned wave
WSDF	wind speed [Eq. 5]	yes	/	/
WBF	wave age + wave steepness [Eq. 8]	yes	T_p	no
WBF $_{\theta}$	wave age + wave steepness [Eq. 11]	yes	T_p	yes
WBF $_{T_m}$	wave age + wave steepness [Eq. 12]	yes	T_m	no

282 Table 1 summarizes 4 experiments conducted in this study, where the only
 283 difference is in the way z_0 is parameterized in COARE3.5. In the first run (dubbed
 284 WSDF), the wind speed only formulation is used (hence, only WRF-ROMS cou-
 285 pling), while in the second run (WBF), the default wave-based formulation is used
 286 (WRF-ROMS-WW3). These two runs are examined in detail in Sections 3-4. Two
 287 additional runs, discussed in Section 5, are conducted with a modified wave-based
 288 formulation. WBF $_{\theta}$ takes into account the directional misalignment between wind
 289 and wave, while WBF $_{T_m}$ modifies the definition of wave age based on mean wave
 290 period rather than the peak wave period.

291 All simulations used in this study produce output every 3h. Since this output
 292 interval is much coarser than the typical sampling intervals used in the observations
 293 (Section 2e), there is inevitable inconsistency in sampling frequency and the number
 294 of samples between the model and data. We attempt to increase the model sample
 295 size and capture more spatio-temporal variability by sampling a slightly broader
 296 region of the model domain encompassing the particular observational tracks (gray
 297 areas in Figure 1a). By doing this we assume that the spatial variability sampled in

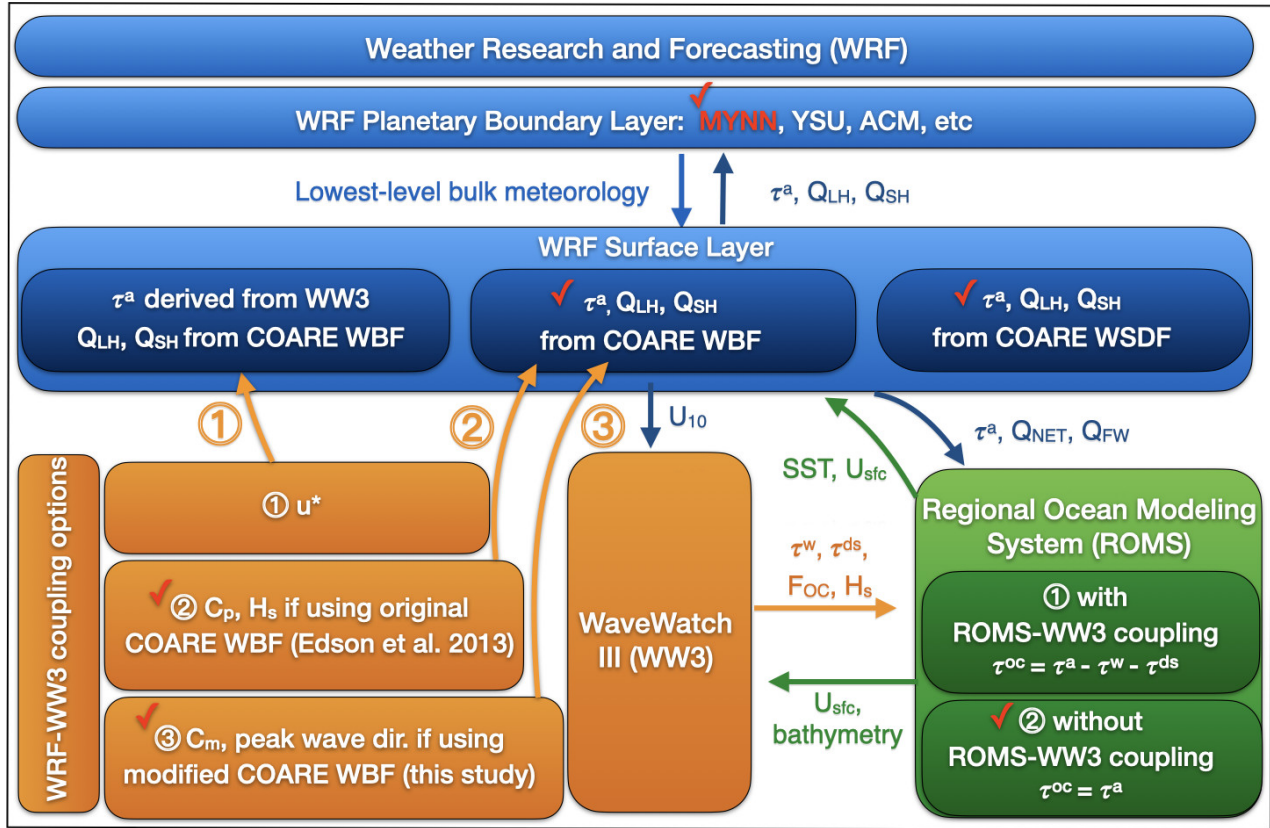


Figure 2. SCOAR WRF-ROMS-WW3 coupling flowchart. See the text for the variable names that are exchanged across the model components. Red ticks denote of the specific schemes and coupling methodology used in this study.

298 the model would resemble the temporal variability observed, considering that the
 299 spatial extent of our model sampling is still relatively close to the different platform
 300 tracks.

301 2.5 ATOMIC/EUREC⁴A observations

302 This study will exploit direct and indirect measurements of momentum fluxes
 303 and relevant wave fields (i.e., significant wave height and wave period) from various
 304 platforms deployed during the ATOMIC/EUREC⁴A experiment, summarized in
 305 Table 2. Figure 1 shows the tracks of the different observational platforms, includ-
 306 ing the NOAA R/V Ronald H. Brown (RHB, Quinn et al., 2021; Thompson et al.,
 307 2021), R/V ATALANTE (Bourras, Geyskens, et al., 2020), SWIFT drifters (Surface
 308 Wave Instrument Float with Tracking, Thomson, 2012; Thomson et al., 2019, 2021),
 309 and OCARINA (Ocean Coupled to Atmosphere, Research at the Interface with a
 310 Novel Autonomous platform, (Bourras, Branger, et al., 2020)) surface naval drone.
 311 The RHB provides direct momentum flux measurements every 10 minutes, using the
 312 eddy covariance method, in the so-called “Tradewind Alley” region from January 9
 313 to February 13, 2020. The SWIFT drifters were deployed from the RHB, from which
 314 the hourly stress can be estimated using the equilibrium frequency range in the wave
 315 spectrum. More specifically, the directional wave spectra and bulk wave paramet-
 316 ers were estimated from inertial motion observations. Then, the friction velocity at
 317 equilibrium u_* is calculated from the wave spectra, assuming a constant equilibrium

318 frequency range over which the source and sink of wave energy is balanced (Iyer
 319 et al., 2022). They were deployed from 14 January to 22 January 2020 and from
 320 30 January to 11 February 2020. The R/V ATALANTE measured the wind stress
 321 mostly in the “Eddy Boulevard” region based on the inertial dissipation method
 322 during the period of January 19 to February 19, 2020. OCARINA was deployed
 323 periodically from the R/V ATALANTE from January 25 to February 17, 2020, pro-
 324 viding direct wind stress measurements every minute through the eddy covariance
 325 method.

Table 2. Summary of the different ATOMIC/EUREC⁴A observations used in this study.

Platforms	R/V Ronald H. Brown	SWIFT	R/V ATALANTE	OCARINA
Observations	wind stress wave periods significant wave height	wind stress wave periods significant wave height	wind stress	wind stress
Methods used in estimating wind stress	eddy covariance	estimated through wave equilibrium subrange	inertial dissipation	eddy covariance
Periods	January 9 to February 13, 2020	14 January to 22 January 2020	January 19 to February 19, 2020	January 25 to February 17, 2020 (periodically)

326 RHB provided data from January 9 to February 13, 2020. SWIFT drifters were
 327 deployed from 14 January to 22 January 2020 and from 30 January to 11 February
 328 2020. R/V ATALANTE provided data from January 19 to February 19, 2020 and
 Ocarina was deployed periodically from January 25 to February 17, 2020

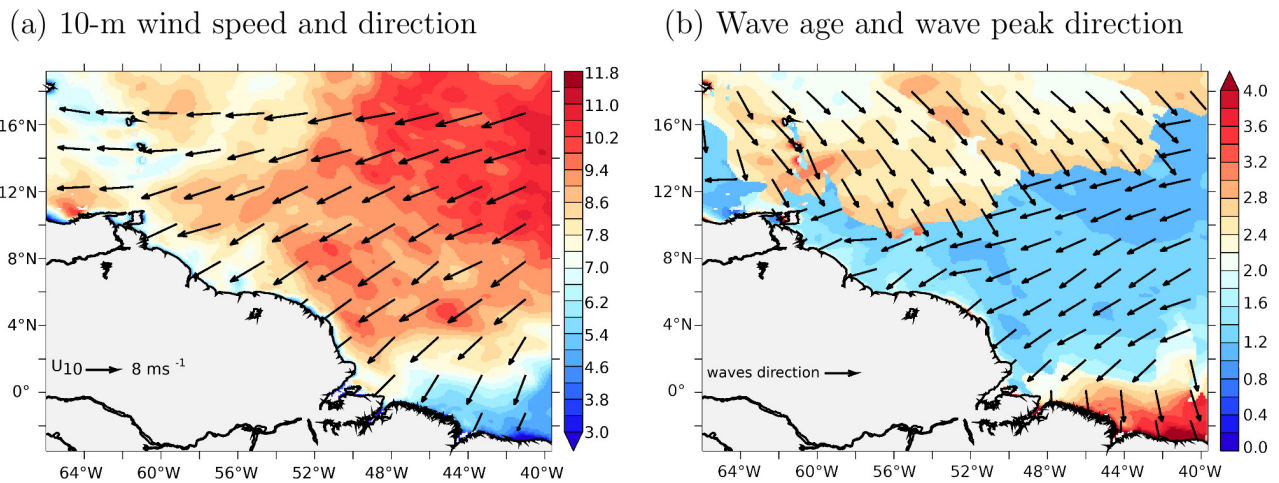


Figure 3. Snapshots of (a) 10-m wind speeds (shading, ms^{-1}) and direction (arrows) and (b) peak wave age (shading) and wave peak direction (arrows) on January 8, 2020 at 0600 UTC.

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3 Impacts of wave and sea state: a case study

To demonstrate the immediate effect of including waves on z_0 and τ in the COARE3.5 using a coupled model, we will first compare the simulation results close to the initial condition. By doing so, the input state variables into the bulk formula remain largely identical, and any differences in simulated z_0 and τ can be attributed to the difference in the formulations. From this set of experiments, we will compare the results 3 hours after the initial condition.

The sea state and wind fields on January 8, 2020 at 0600 UTC, shown in Figure 3a, illustrate the archetypal synoptic condition observed in this region during the boreal winter. Much of the domain was under the influence of northeasterly trade winds with wind speeds of 7-13 ms^{-1} , while the northern and southeastern parts of the domain experienced much weaker ($<7 ms^{-1}$) easterly and northerly winds, respectively. Figure 3b shows the corresponding wave age and peak wave direction. In the Tradewind Alley region, surface waves were predominantly downwind with relatively small wave age, indicating the developing seas with young waves. Away from the trade winds, especially in the northern part of the domain, the wave vectors are generally misaligned with the local wind vectors, and the wave age is high, indicative of the swell-dominated sea state.

To illustrate sea state distribution differently, Figure 4a shows the probability density function (PDF) of wave age for the same period. Two distinct peaks of wave age stand out clearly. The first peak resides on wave age between 0.8 and 1.7, corresponding to developing (young) waves to fully developed (mature) seas. The secondary peak is found over a wide range of wave age greater than 1.7, reaching up to 4-5, the latter representing swell. Indeed, the fact that there is a gap at 1.7 strongly suggests that the older waves are swell, as opposed to the continuum of longer/older wind waves. Thus, in this case, we choose to use 1.7 as a threshold for fully developed seas and not the usual value of 1.2 which is what you might expect for wind waves dominated region. As a matter of fact, this swell-dominated sea state is frequently observed in the ATOMIC region in the boreal winter (e.g., Semedo et al., 2011; Jiang & Chen, 2013). Indeed, if considering the entire month of January 2020 in our simulations, we find that wave ages greater than 2 occur more than 60% of the time in this domain.

Figure 4b compares the z_0 against wind speed from the WSDF (black) and WBF (color) runs for this period. z_0 from WBF is color-coded to denote the corresponding wave age. The bottom panel shows stacked PDFs of 10-m wind speeds from WBF, with the red (gray) parts representing the proportion of wind associated with wave age over (under) 1.7. The WSDF in COARE3.5 assumes young seas under moderate to high winds, and hence the parameterized z_0 (black) obeys the well-known quadratic dependence on wind speed. The surface roughness z_0 from WSDF shows less scatter because it is based solely on wind speed.

In contrast, WBF captures the two wave age-dependent regimes of z_0 that appear distinct from WSDF. The first is the cluster of z_0 , which increases more rapidly with wind speed than WSDF z_0 and occurs over 4-12 ms^{-1} . The wave age of this cluster (shading) is typically less than 1.7, corresponding to the first wave age peak in Figure 4a of small-scale young waves. Thus, the developing and equilibrium waves under these wind speeds and wave age conditions increase z_0 in WBF compared to WSDF.

The second cluster indicates significantly decreased z_0 in WBF with wind speed up to 12 ms^{-1} . This cluster can be further split into two different wind speed groups, under and above 8 ms^{-1} , color-coded by the PDF of winds (Figure 4b). Below 8 ms^{-1} (red, weak winds), the wave age mainly constitutes the tail of the

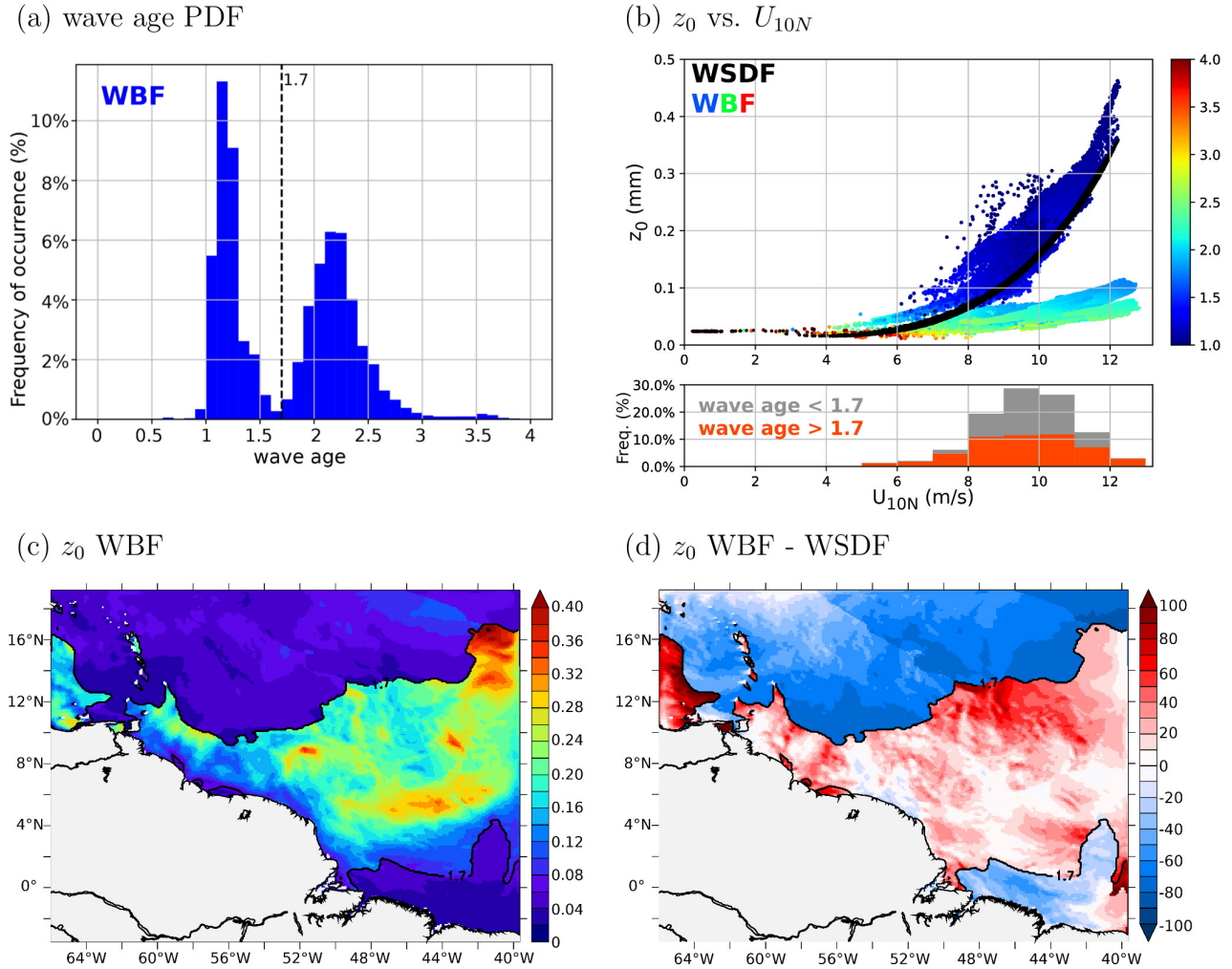


Figure 4. (a) PDF of wave age from the entire model domain on January 8, 2020 at 0600 UTC. The dotted vertical line denotes the wave age of 1.7, below (above) which the sea state is characterized as developing, equilibrium and slightly old waves (mature waves and swell). The upper panel of (b) is a scatter plot of z_0 (mm) vs. U_{10N} (ms^{-1}). z_0 from WSDF is shown in black, while z_0 from WBF is color-coded to denote the corresponding wave age. The stacked PDFs of U_{10N} in the lower panel of (b) are constructed when wave age is above 1.7 (red) and below 1.7 (gray). (c) A map of z_0 from WBF, superposed with a contour of wave age = 1.7. (d) A map of percentage difference of z_0 between WBF and WSDF

381 PDF distribution shown in Figure 4a with an average wave age of 2.7. It is where re-
 382 motely generated swell appears to dominate the sea state. However, the wind speeds
 383 under $8 ms^{-1}$ account for less than 10% of the total wind speed data, and thereby it
 384 has a relatively small impact on the space/time-averaged z_0 . Indeed, when averaged
 385 for wind speed below $8 ms^{-1}$, the percentage difference in z_0 between WSDF and
 386 WBF, defined as $(WBF-WSDF/WSDF)*100$, is only -1.7%.

387 During this day, most of the wind speed is above $8 ms^{-1}$. In addition to the
 388 proportion of low wave age expected under this moderately high wind speed, we also
 389 find an increased occurrence of large wave age, accounting for 44% of the data (Fig-
 390 ure 4b). The co-existence of high wind and swell indicates a mixed sea condition. In

391 this case, when averaged over wind speed above 8 ms^{-1} , the swell impact appears
 392 much more significant, with z_0 in WBF being 15.7% lower than that in WSDF. The
 393 working hypothesis is that the use of the phase speed at the spectral peak causes the
 394 WBF to assume that the swell is supporting most of the stress even under moderate
 395 winds. This strong impact of swell on z_0 at such moderately strong winds is ques-
 396 tionable, in the sense that the majority of air-sea momentum exchanges should still
 397 be supported by short-scale coupled wind waves despite the co-existence with the
 398 long-wave swell.

399 The spatial distribution of z_0 from WBF is shown in Figure 4c. The z_0 dif-
 400 ference between WBF and WSDF is shown in Figure 4d. As in Figure 4a,b, two
 401 distinct regimes of z_0 are readily apparent on the map, delineated sharply by the
 402 contour of wave age 1.7 (black). The horizontal discontinuities in the wave and z_0
 403 fields (Figure 4c,d) appear only with the use of the peak period, while the use of
 404 average wave period produces much smoother fields (not shown). The location of the
 405 front is only because this is a snapshot of the sea state on 8 January at 0600 UTC.
 406 Snapshots 3h before/after would show the swell front displaced to another location
 407 as the swell is moving/dissipating. In the first regime of increased z_0 in WBF under
 408 moderate to strong trade winds, the WBF predicts an increased z_0 by on average
 409 25% compared to WSDF. This increased z_0 is expected as the WBF z_0 formulation
 410 (Eq. 8) takes into account the effect of wave slope on the aerodynamic roughness
 411 of the sea surface. That is, Figure 5a,b show that wave slope under young waves is
 412 higher, where the choppy sea surface increases z_0 . Figure 5c,d shows the angle (θ)
 413 between the wind direction and peak wave direction. If $\theta = 0^\circ$, wind and waves are
 414 perfectly aligned, whereas $\theta = 180^\circ$ means wind and waves are opposed. Collocated
 415 with the regime of increased z_0 , the peak wave direction is largely downwind, since
 416 θ is generally less than 50° . This corroborates that these waves are young waves
 417 driven by local winds. In the present study only the peak wave direction is used to
 418 defined alignment/misalignment with the local wind. However, at times, the wave
 419 field can yield significant directional spreading, this aspect is discussed later on in
 420 Section 5.2.

421 Figure 4d also shows the second regime of decreased z_0 with the inclusion of
 422 waves, especially in the northern part of the domain. In this region, the remotely
 423 generated swell propagates into the domain through the northern boundary and
 424 forms a sea state with the aerodynamically smooth sea surfaces (Figure 5a,b) and
 425 with waves whose direction is strongly misaligned ($\theta = 60-160^\circ$) with the local wind
 426 (Figure 5c,d). In particular, the reduced z_0 over swell persists under wind speed of
 427 up to 12 ms^{-1} (Figure 3a), despite the expectation that under such a high wind, the
 428 wind-waves would still strongly increase the aerodynamic roughness and stress.

429 Figure 6a,b compare the parameterized wind stress in WBF and WSDF. One
 430 can see from these plots a consistent difference in wind stress due to the inclusion of
 431 waves. Wind stress decreases sharply in wind speeds of $8-12 \text{ ms}^{-1}$ over the northerly
 432 swell, where wave age >1.7 . At the highest wind speed during the event, the per-
 433 centage difference in wind stress magnitude exceeds 10%. Conversely, wind stress
 434 is increased in WBF by $\approx 4\%$ over fully developed seas (wave age <1.7) and high
 435 winds, consistent with the increase in z_0 there (Figure 4c). By comparing to the
 436 direct momentum flux observations, we will determine in Section 4 if such reduced
 437 z_0 and τ over swell conditions at moderate to high wind speeds are consistent with
 438 the observations. As COARE3.5 does not consider the misaligned waves with winds,
 439 these conditions may constitute a source of uncertainty in the parameterized z_0 and
 440 τ via COARE3.5 WBF. As for the large wave age in the southeastern corner of the
 441 domain, it is concurrent with weaker winds (Figure 3a), and hence the assumptions
 442 about the swell under weaker wind seem valid in this region. This leads to a small
 443 difference in z_0 between WBF and WSDF.

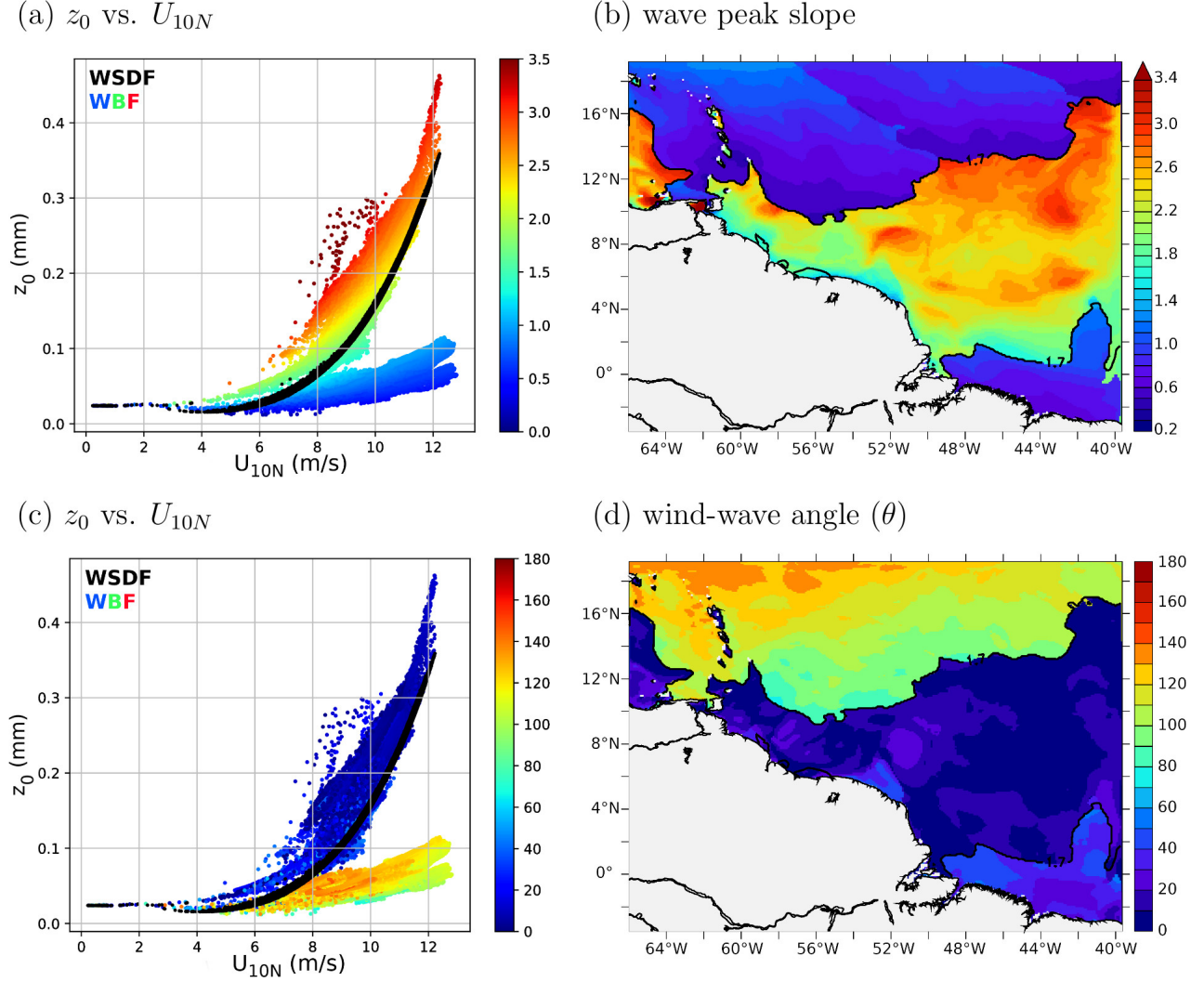


Figure 5. (a) Scatter plot of z_0 (mm) vs. U_{10N} (ms^{-1}) from WSDF in black and WBF color-coded to denote the corresponding wave peak slope (10^{-2}) defined as H_s/L_p where L_p is the peak wavelength. (b) A map of wave slope peak (10^{-2}), superposed with a contour of wave age = 1.7 on January 8, 2020 at 0600 UTC. (c,d) As in (a-b) except that colored scatters and shading denote the angle between the wind and wave directions ($^\circ$).

444 The altered stress directly influences the low-level winds via the surface drag.
 445 Here, we estimate the response in low-level winds at the lowest WRF model layer,
 446 at about 27 m above the sea surface. Figure 6c shows that the low-level wind is
 447 increased over the aerodynamically smooth sea surface due to swell by $>0.5 ms^{-1}$,
 448 accounting for 5-20% of the wind speed in WBF. In contrast, where young waves
 449 dominate in WBF, the wind stress is increased by 5% and the wind speed is de-
 450 creased.

451 One relevant physical process that represents the air-sea momentum transfer
 452 affecting the winds and surface currents, is the wind work (P),

$$P = \frac{1}{\rho_o} (\overline{u_s \tau_x} + \overline{v_s \tau_y}), \quad (11)$$

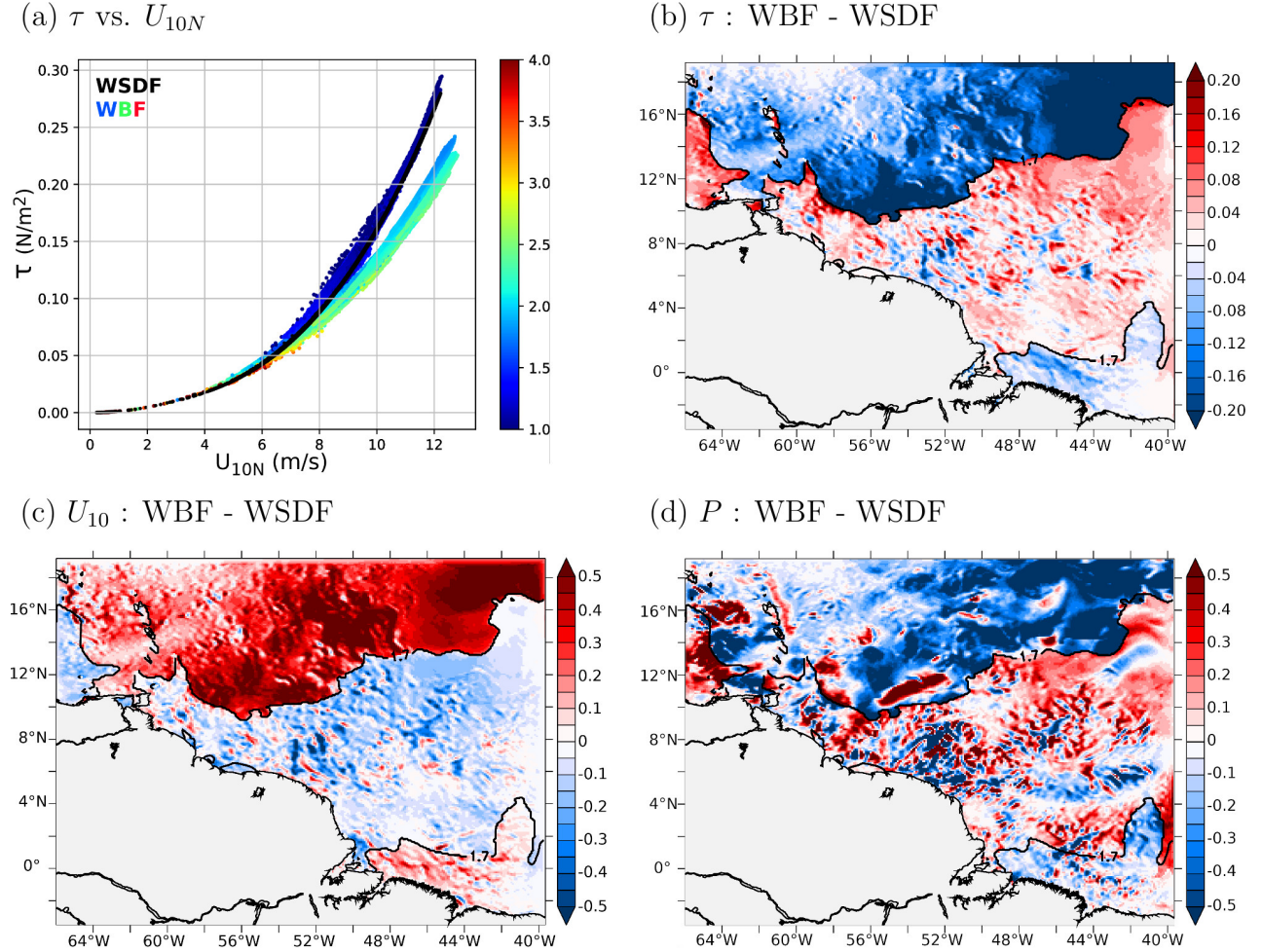


Figure 6. (a) Scatter plot of τ (Nm^{-2}) vs. U_{10N} (ms^{-1}) from WSDF in black and WBF color-coded to denote the corresponding wave age. (b,c,d) Difference maps between WBF and WSDF of (b) τ (10^{-1}Nm^{-2}), (c) U_{10} (ms^{-1}), and (d) wind work (P , $10^{-5} \text{m}^3 \text{s}^{-3}$) on January 8, 2020 at 0600 UTC, superposed with a contour of wave age = 1.7.

453 where (u_s, v_s) are the surface current vectors, (τ_x, τ_y) are the wind stress vec-
 454 tors, and the overbar denotes the time-average. When P is positive, the mechanical
 455 work is done by the wind stress on the ocean surface currents, increasing the ocean
 456 kinetic energy (e.g., Wunsch, 1998). When negative, it represents the diversion of
 457 the ocean energy by the current to the wind, accelerating the low-level winds at
 458 the expense of weakened surface currents (e.g., Renault et al., 2016, 2017; Seo et
 459 al., 2019, 2021). Figure 6d shows the difference in P between WBF and WSDF for
 460 this snapshot. The region of reduced τ and increased low-level wind in the swell-
 461 dominated region is congruent with the region of the robust decrease in P , while the
 462 opposite is true in the Tradewind Alley region. The difference in P mainly reflects
 463 the changes in wind stress due to waves (Figure 6b).

4 Modeled and observed momentum fluxes during ATOMIC

464
 465 Determining whether or not the parameterized z_0 and τ with WBF represents
 466 an improvement over WSDF requires a detailed comparison to direct covariance

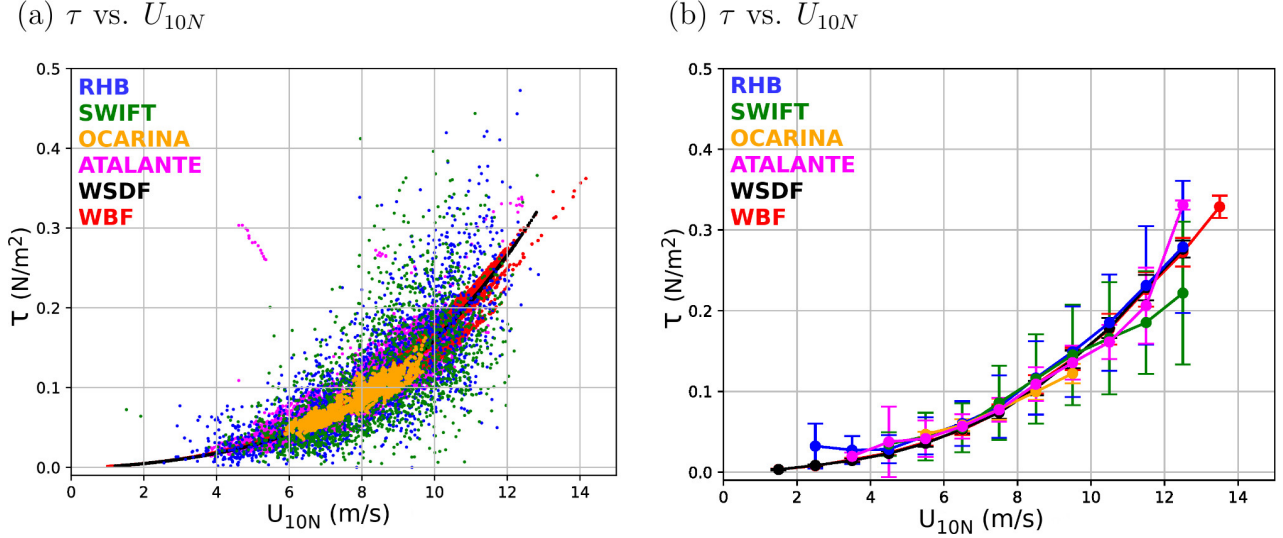


Figure 7. (a) Scatter plot comparing the two parameterized τ (Nm^{-2}) using COARE3.5 WSDF (black) and WBF (red) against the various types of measurements of τ (see Section 2e for a description of the various methodologies). (b) As in (a) except that measurements are bin-averaged with a wind speed bin-size of $U_{10N} = 1 \text{ ms}^{-1}$. The error bars represent ± 1 standard deviation. Only bins with more than 5 points are plotted.

467 stress measurements. In this section, we will compare the model simulation with
 468 the observations during the EUREC⁴A/ATOMIC experiments to evaluate the ac-
 469 curacy of the wave-based parameterized τ and identify the regimes where further
 470 improvements might be needed.

471 Figure 7a compares the two modeled stresses to the observations. All observa-
 472 tions and the two model simulations display the quadratic relationship of wind stress
 473 with wind speed. RHB and SWIFT, sampling the stress mainly in the Tradewind
 474 Alley region, produce greater scatter compared to ATALANTE and OCARINA,
 475 which were deployed further south in the Eddy Boulevard region (1a). The signifi-
 476 cant departure from this curve in the Tradewind Alley region may reflect the greater
 477 uncertainties in determining τ from these measurements. Between the model simu-
 478 lations, WBF produces a larger spread than WSDF, yet their averages at given wind
 479 speed are similar (Figure 7b). Overall, parameterized stresses by WSDF and WBF
 480 both agree well with the observations to within the observational errors during the
 481 campaign.

482 Figure 8a compares the histograms of the wave age from the WBF run to those
 483 from the SWIFT drifters and the RHB. It should be noted that in both the model
 484 and measurements, the wave age is estimated using the peak period (T_p). The ob-
 485 servations and model simulation show the bi-modal distribution of wave age as was
 486 seen from the snapshot case in Section 3 (Figure 4a), with the first peak near wave
 487 age 1.7 and the secondary, much broader, peak between 2.5-3. The SWIFT obser-
 488 vations (in red) capture a higher occurrence of young waves than the RHB obser-
 489 vations or the WBF simulation. WBF also features a fatter tail of the distribution
 490 toward larger wave ages, indicating that the model overemphasizes the occurrences
 491 of swell and decaying waves compared to these observed estimates.

492 Given the wave age distributions, we then divide the distribution into 3 dif-
 493 ferent “Regimes” to better understand the wave age-dependent z_0 -wind speed and

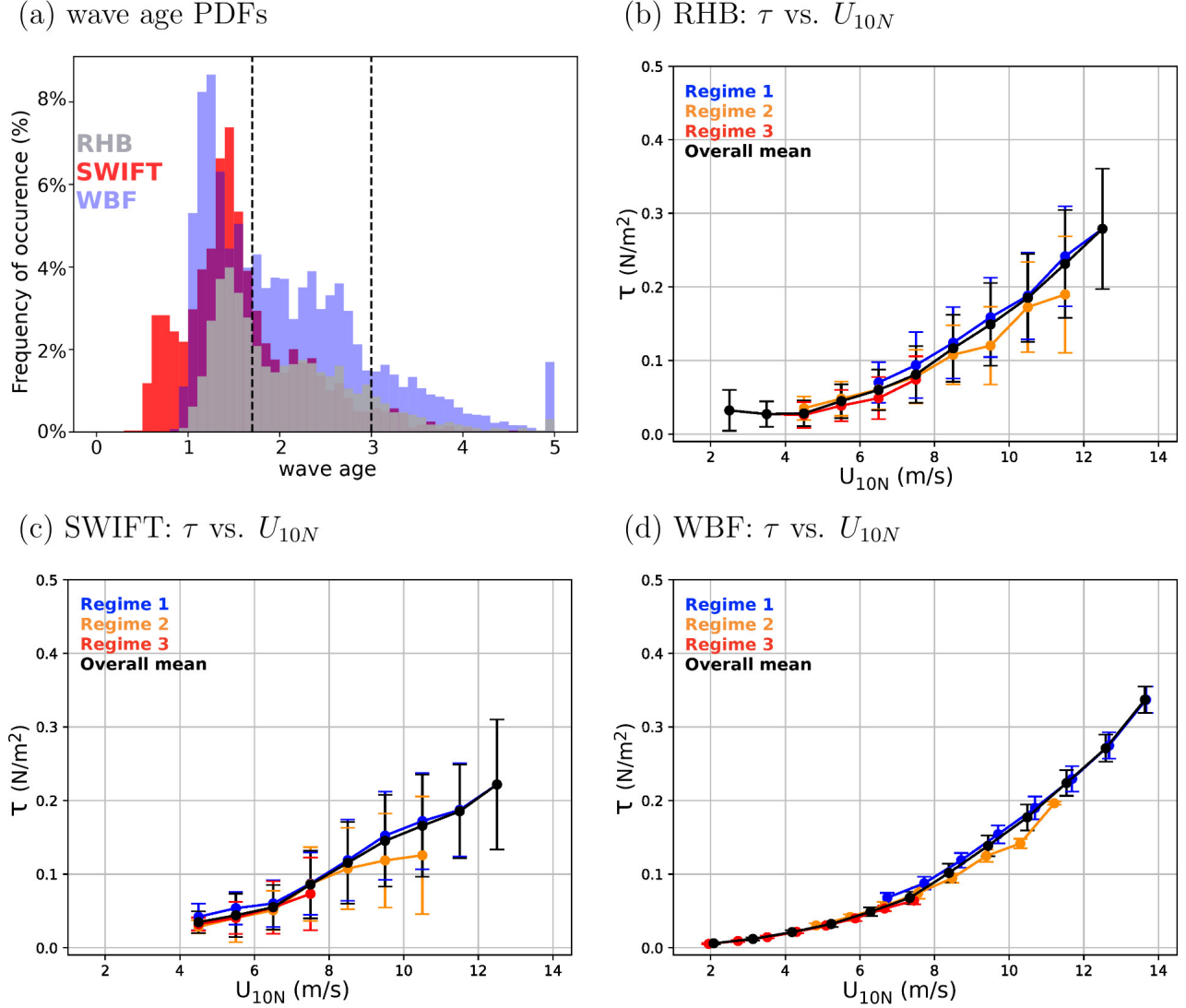


Figure 8. (a) Peak wave age distribution estimated from SWIFT (red), RHB (gray), and WBF (blue). Here, wave age is capped at 5. Three wave age regimes are defined: Regime 1 (blue) when wave age < 1.7 denotes the young sea to fully developed sea, Regime 2 (orange) when wave age is between 1.7 and 3 indicates the mature to old sea, and Regime 3 (red) when wave age > 3 represents the old sea and non-locally generated swell. (b-c) Binned scatter plots of τ (Nm^{-2}) vs. U_{10N} (ms^{-1}), color-coded to show the three different wave age Regimes, with the bin-average of 1 ms^{-1} . The error bars represent ± 1 standard deviation. Only bins with more than 5 points are plotted. The mean of all wave ages is shown in black. (d) As in (b) and (c) except from the WBF run. Here WBF is sampled along-track of the RHB and SWIFT.

494 τ -wind speed relationships. Regime 1 refers to young to fully developed seas, defined
 495 as when wave age < 1.7 , while Regime 2 indicates the mature to old sea, including
 496 mixed sea state, which is diagnosed as wave ages between 1.7 and 3. Finally, the old
 497 sea and non-locally generated swell characterizes Regime 3 estimated as when wave
 498 age > 3 . When using the peak period, and to stay consistent throughout the paper,
 499 thresholds are kept the same. However, these thresholds are not necessarily universal
 500 but can vary in different times or regions under consideration.

501 The colored lines in Figures 8b and c show the bin-averaged surface stress
 502 from the RHB and the SWIFT from the 3 Regimes. The black lines denote the
 503 bin-averaged surface stress across all wave age regimes. Despite the significant error
 504 bars, which represent ± 1 standard deviation, one can observe the consistent rela-
 505 tionship between the measured stress and the wind speed across different wave age.
 506 For example, the measured stress over Regime 1 (blue) is higher than the overall
 507 average (black) as the short-wind waves support the bulk of momentum exchanges.
 508 In contrast, the stress over Regime 2 (orange) and Regime 3 (red) is lower than the
 509 overall average, as the sea state is characterized by mixed and older seas. This sea
 510 state dependence of wind stress is also somewhat evident in the WBF simulation
 511 (Figure 8d) despite the smaller error bars likely due to smaller number of samples in
 512 the model, as discussed in Section 2d.

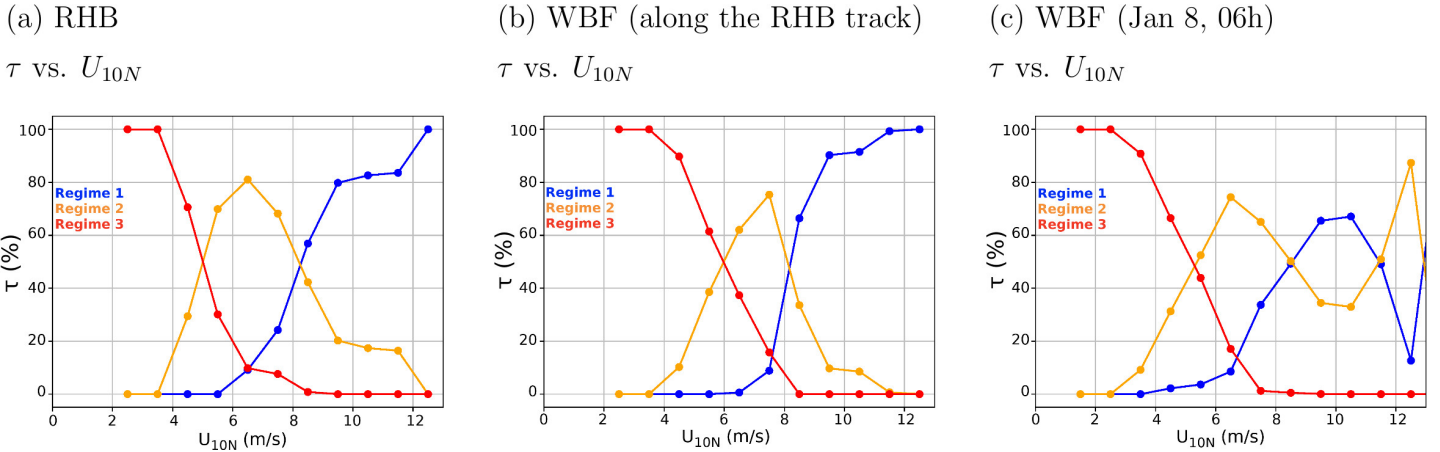
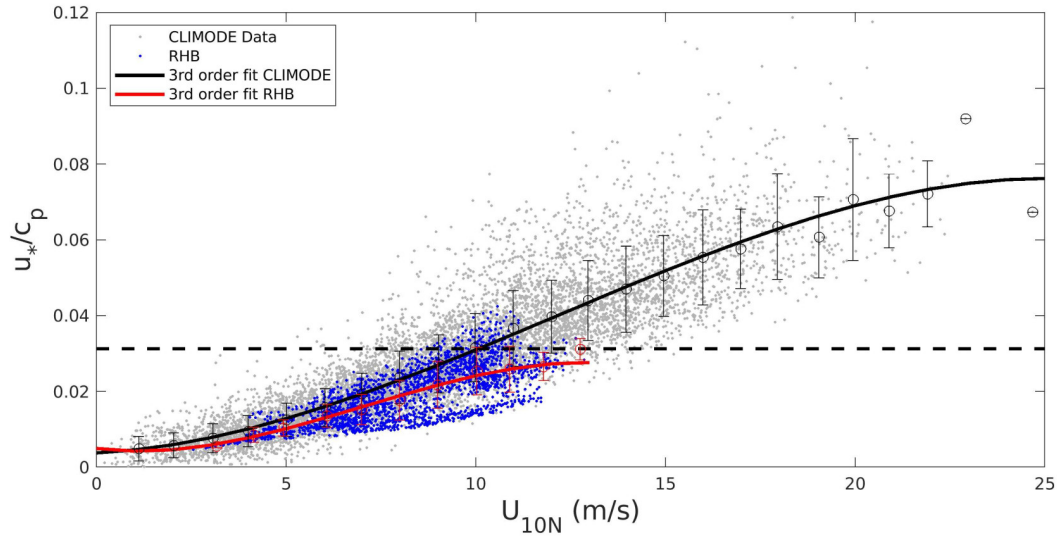


Figure 9. Percentage contribution of τ (%) by the three different wave age Regime at a given wind speed (bin averaged every 1 m s^{-1}) from (a) RHB, (b) WBF sampled along the RHB track between January 9 and February 13, 2020 and (c) WBF sampled over the whole model domain on January 8, 2020 at 0600 UTC. The different colors denote the different wave age categories described in Figure 8.

513 To further quantify this relationship, Figure 9a shows the percentage of stress
 514 supported by the different wave-age Regimes from the RHB observations, binned
 515 over 1 m s^{-1} intervals. Under 4 m s^{-1} wind speeds, the surface stress is mainly
 516 supported by Regime 3 (red), whereas above 8 m s^{-1} , Regime 1 (blue) dominates
 517 the contribution to the stress. Regime 2, which represents mixed sea conditions
 518 ($4\text{--}8 \text{ m s}^{-1}$) and contributes to less than 20% of the stress above 10 m s^{-1} . Figure 9b
 519 shows the same diagnostics, but for the WBF run sampled along the track of RHB.
 520 It shows that the WBF overall exhibits a similar fractional contribution to stress.
 521

522 When the model is compared to the observations at this particular track, WBF
 523 appears to accurately characterize the observed stress relationship with wave age
 524 (See also Figure 8). However, if sampled over a broader region of the same mixed
 525 sea conditions from the model, a different result is obtained. Figure 9c shows the
 526 same results as Figure 9b, except that the entire model domain is sampled under
 527 the same synoptic condition examined in Section 3. It shows that the parameterized
 528 stress under $8\text{--}12 \text{ m s}^{-1}$ wind speeds supported by Regime 2 (orange) is comparable
 529 to the stress supported by Regime 1 (blue) as also seen in Figure 6. In reality, short
 530 wind waves under such wind speeds should still support the increased stress despite
 531

(a) u_*/c_p vs. U_{10N}



(b) u_*/c_p vs. U_{10N}

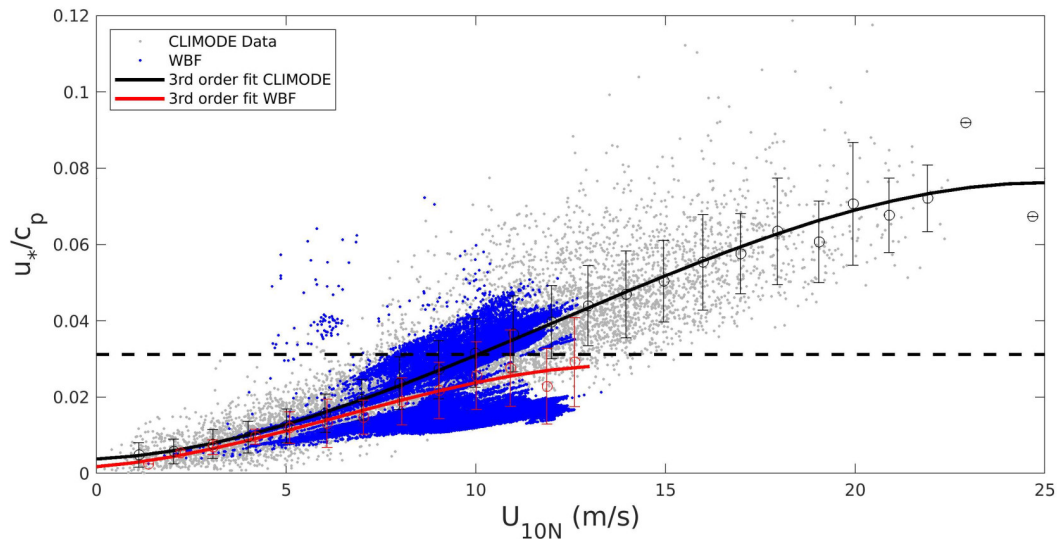


Figure 10. (a) Scatter plot of inverse peak wave age (u_*/c_p) vs. U_{10N} (ms^{-1}) for CLIMODE data (gray) and RHB data (a, blue). Bin-averages with the 1 standard deviation error bars are overlaid, at $1 ms^{-1}$ interval, along with the 3rd order fit (line) for CLIMODE (black) and RHB (red). The horizontal dashed line is $u_*/c_p = 0.03$, denoting the threshold for fully developed seas (equivalent to $c_p/U_{10N} = 1.2$). (b) As in (a) but RHB data is replaced with WBF, for the whole domain

on January 8, 2020 at 0600 UTC.

531 the higher wave age, we believe this is a form of deficiency in COARE3.5 WBF in
 532 representing the wind stress over mixed swell-dominated seas.

533 In fact, the COARE3.5's WBF was developed and tuned primarily by using
 534 the wave data collected from the extratropics, where sea state tends to be dominated
 535 by growing and fully-developed waves under high winds (see Figure 2 in Edson et

536 al., 2013). Figure 10 compares the sea state used to tune COARE3.5, taken during
 537 the CLIMODE campaign (CLIVAR Mode Water Dynamic Experiment, Marshall et
 538 al., 2009), with the sea state observed by RHB during January-February 2020 and
 539 modeled in WBF on January 8, 2020 at 0600 UTC in the ATOMIC region. It shows
 540 the relationship between the inverse wave age and U_{10N} . Here, a low inverse wave
 541 age is indicative of decaying seas and swells. An inverse wave age of 0.03 (dashed
 542 line) is roughly equivalent to an equilibrium wave age of 1.2. As expected, the sea
 543 state captured in the ATOMIC region is very different and much older than the
 544 one used in COARE3.5. Therefore, the wind stress under moderate winds and swell
 545 dominated conditions observed here, and possibly in other tropical oceans, may not
 546 be currently well parameterized in the COARE3.5 WBF. The specific deficiency
 547 identified from this analysis is that, for mixed seas (Regime 2) where high wave age
 548 and moderately strong wind co-occur, the current COARE3.5 WBF overemphasizes
 549 the swell impact on wind stress, leading to the low-stress bias despite the moderately
 550 strong winds.

551 5 The revised wave-based formulation in COARE3.5

552 In the following, we present two experimental revisions to the z_0 formulation in
 553 the current COARE3.5 WBF for swell conditions coincident with moderate to high
 554 winds, the condition that is frequently observed in the northern ATOMIC region
 555 in the boreal winter. One method is to replace the peak wave period (T_p) with the
 556 mean wave period (T_m) in the definition of the phase speed and thus wave age, and
 557 another is to incorporate the effect of misaligned waves with local wind on aerody-
 558 namic roughness in the z_0 parameterization. In essence, these two observationally-
 559 guided approaches desensitize the impact of swell on z_0 and τ estimates at moderate
 560 winds and alleviate the low biases in the current COARE3.5 WBF. For this, we now
 561 return to the case study on January 8, 2020 as in Section 3.

562 5.1 The mean wave period

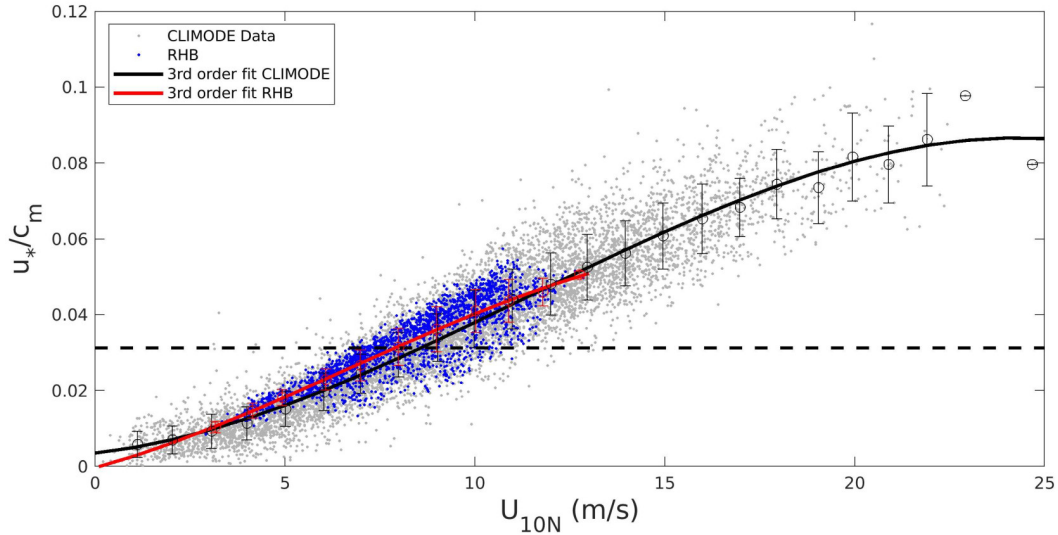
563 One possible approach to mitigate the overestimation of the swell impact on
 564 z_0 and τ under moderate to high winds is to use the wave’s mean period, T_m , to
 565 calculate the average phase speed, c_m , in the wave age definition. This change is mo-
 566 tivated by the finding that T_p does not accurately describe a mixed-sea state where
 567 swell and wind-sea co-exist, as shown in Figure 10. T_p can be also sensitive to the
 568 spectral shape of the wave energy and the chosen filter, while T_m can be reliably es-
 569 timated from observations and WW3 as either an energy-weighted average period or
 570 zero-crossing period. A similar argument has been made recently by (Colosi et al.,
 571 2021) as they chose to use a wave age dependent computed with the mean period to
 572 construct the seasonal probability of swell over global oceans.

573 We carried out an additional coupled simulation, dubbed WBF_ T_m , where T_p
 574 is replaced with T_m to get the mean phase speed of the waves c_m in Eq. 12:

$$z_{rough} = H_s D \left(\frac{u_*}{c_m} \right)^B, \quad (12)$$

575 where $D=0.39$ and $B=2.6$, which have been tuned using the COARE3.5 set
 576 of observations. We will estimate T_m based on the zero-crossing period, as it is the
 577 one used to describe T_m in the observation. Figure 11 shows the same diagnostics
 578 as in Figure 10 but this time using c_m to calculate the inverse wave age in both the
 579 observations, CLIMODE and RHB, and the WBF_ T_m run. The general trend of
 580 both sets of observations are now in good agreement (Fig. 11a). In WBF_ T_m , the
 581 use of c_m in eq. 12 alleviates the bias over the mixed sea (Regime 2) (Figure 10b vs.
 582 Fig. 11b) and shows a better agreement of the general trends from the observations.

(a) u_*/c_m vs. U_{10N}



(b) u_*/c_m vs. U_{10N}

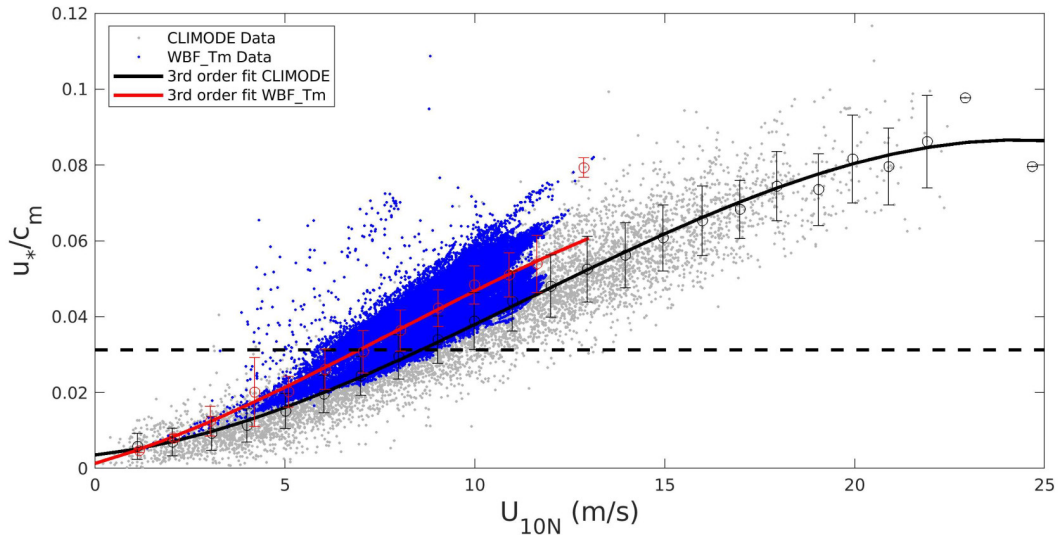


Figure 11. (a) As in Figure 10a, but with inverse mean wave age (u_*/c_m). The dashed line is $u_*/c_m = 0.03$, denoting the threshold for fully developed seas (equivalent to $c_m/U_{10N} = 1.2$). (b) As in Figure 10b except for showing the result from WBF_T_m

583 Further refinement of coefficients in eq. 12 will be addressed in more detail in the
 584 future release of the COARE4.0 algorithm.

585 Figure 12a shows the PDF of wave age for RHB (gray), SWIFT (red), and
 586 WBF_T_m (blue) computed using T_m . This figure should be compared to Figure 8a
 587 where RHB, SWIFT and WBF wave age PDFs were computed using T_p . Similar to
 588 Figure 8a, wave age is capped at 5 to show the tail of the distribution. In contrast
 589 to the bi-modal distribution of wave age with the pronounced secondary peak of
 590 wave age estimate with T_p , the use of T_m effectively removes this secondary peak
 591 in both the model and observations, yielding a markedly different distribution with
 592 an overall prevalence of younger sea state. We adjusted the different categories of

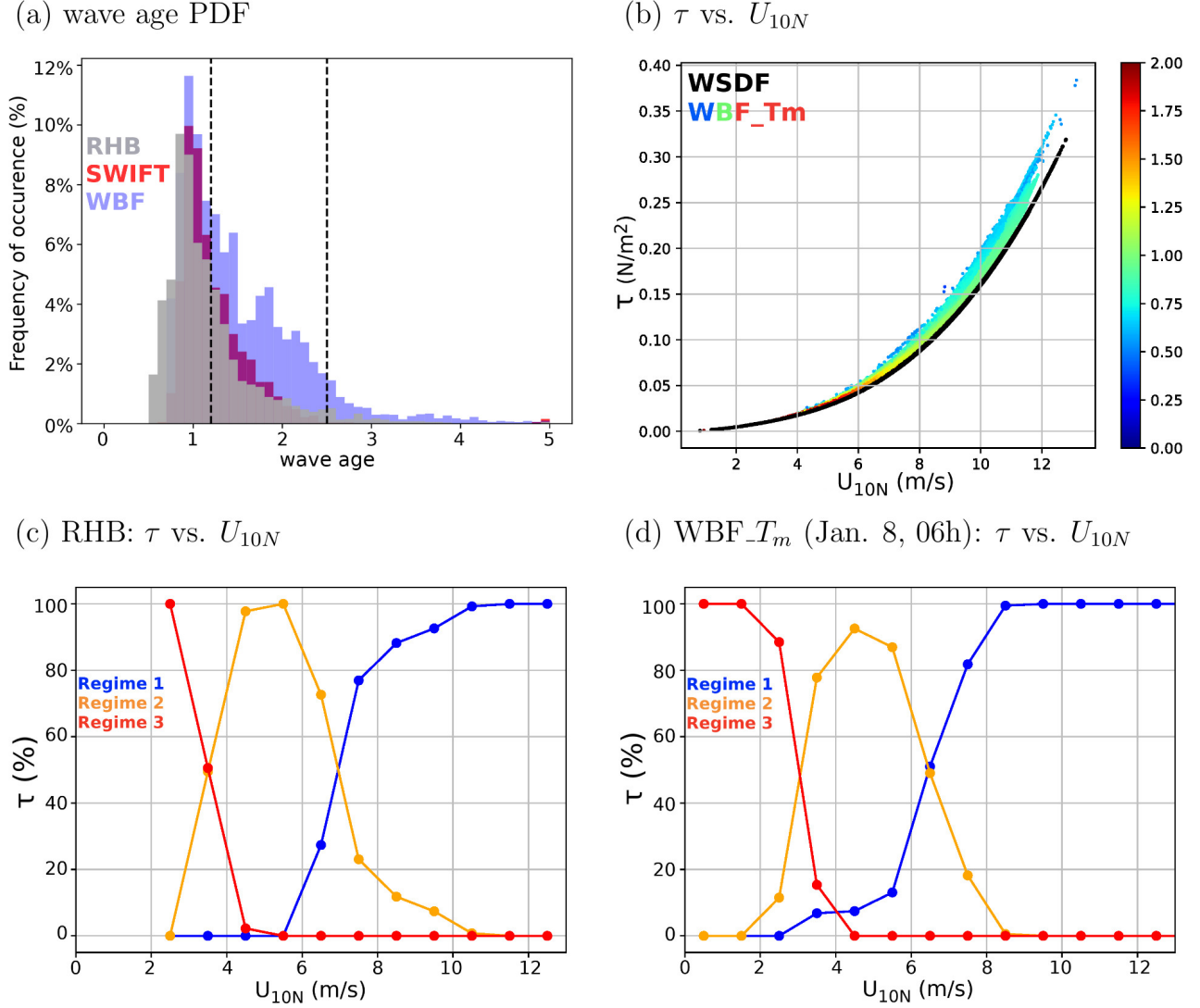


Figure 12. (a) Mean wave age distributions estimated from RHB (gray), SWIFT (red), and WBF_ T_m (blue). WBF_ T_m is sampled along-track of the RHB and SWIFT. (b) Scatter plot of τ (Nm⁻²) vs. U_{10N} (ms⁻¹) from WSDF in black and WBF_ T_m color-coded to denote the corresponding wave age on January 8, 2020 at 0600 UTC. (c,d) As in Figure 9a,c, except that the wave age is defined with T_m for (c) RHB and (d) WBF_ T_m .

593 wave age defined previously to fit the new wave age distribution based on T_m . Figure
 594 12b shows τ on January 8, 2020 at 0600 UTC from WBF_ T_m , with wave age
 595 color-coded. The cluster of low z_0 with high wave age seen in Figure 4b is elimi-
 596 nated in WBF_ T_m , because of the elevated z_0 and τ under moderate to high wind
 597 speeds. Finally, Figure 12c,d, to be compared to Figure 9a,c shows the percentage of
 598 τ supported by each category of wave age for RHB and for WBF_ T_m , respectively.
 599 With the use of T_m , WBF_ T_m agrees well with RHB concerning the fractional con-
 600 tribution from each sea state to the surface stress. Particularly over 7 ms⁻¹, most of
 601 the contribution to τ now comes from the wind sea (blue), whereas the contribution
 602 of mature seas and swell subsides rapidly with the increased wind speeds. This is
 603 a clear improvement from τ parameterized using T_p (Figure 9c) and is much more
 604 consistent with the observations (Figures 9a, 12c).

5.2 Including the (mis)aligned wind-wave directions

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As discussed in Section 2, the COARE3.5 assumes the wave stress as a scalar roughness parameter, and hence the direction of wave-stress vectors is aligned with the mean wind vectors. However, wave stress and mean wind vectors can be misaligned under various conditions, including under rapidly translating storms (e.g., S. S. Chen et al., 2013), near strong vorticity and divergence gradients and density fronts (e.g., Villas Bôas & Young, 2020), or over mixed seas where wind waves and swells co-exist under high winds. Such nonequilibrium wave motions can influence wave slope, roughness length, and wind stress (Janssen, 1991; Rieder et al., 1994; Zou et al., 2019; Patton et al., 2019; Porchetta et al., 2021; Deskos et al., 2021). Here, we attempt to incorporate the directionality of the wind and waves following Patton et al. (2019) and Porchetta et al. (2019), such that

$$z_{rough} = H_s D \cos(a\theta) \left(\frac{u_*}{c_p}\right)^{B \cos(b\theta)}. \quad (13)$$

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D and B are the coefficients taken from COARE3.5 (See Eq. 8), while the coefficients $a = 0.4$ and $b = 0.32$ are adopted from (Porchetta et al., 2019). In principle, all these coefficients require site-specific tuning. For example, (Porchetta et al., 2019) used the high wind conditions observed from the FINO platform in the North Sea and the Air-Sea Interaction Tower (ASIT) in the New England Shelf, which represents different wind speed and wave age conditions from the trade-wind and swell-dominated tropical oceans as in the ATOMIC domain. Additional tuning exploiting direct momentum flux measurements would be needed to develop a refined set of coefficients for the tropical oceans. This is beyond the scope of the study. Using this new formulation, we conducted an additional coupled experiment, dubbed WBF_θ , which is to be compared to the default wave-based formulation in COARE3.5, where $\theta = 0$.

(a) τ vs. U_{10N}

(b) τ vs. U_{10N}

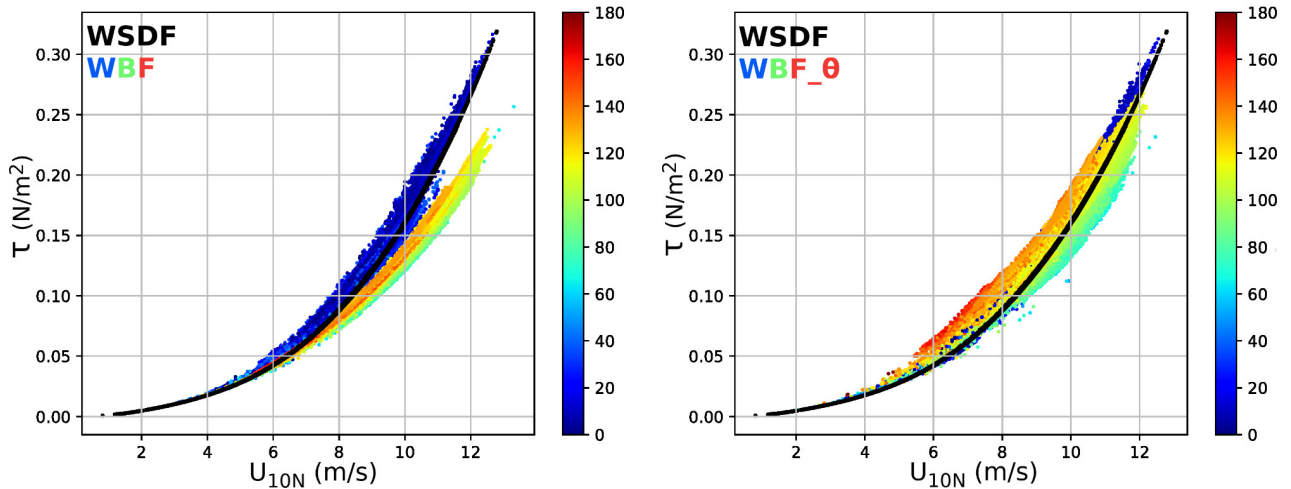


Figure 13. (a) Scatter plot of parameterized τ (Nm⁻²) vs. U_{10N} (ms⁻¹) from WSDF in black and WBF color-coded to denote the corresponding wind-wave angle (θ) on January 8, 2020 at 0600UTC. Note that in the z_0 formulation in WBF assumes $\theta = 0$. (b) As in (a) except from WBF_θ , where θ is treated as a non-zero quantity in the z_0 formulation.

629 Figure 13a compares the parameterized τ , color-coded by the angle (θ) be-
 630 tween the wind direction and peak wave direction in WBF. It shows that the lower
 631 τ from WBF compared to WSDF (and also observations) occurs when the swell
 632 waves are strongly misaligned with winds (e.g., $\theta > 60-90^\circ$). This indicates that the
 633 assumption of $\theta = 0$ in WBF can be attributed to the lower τ . When the directional
 634 misalignment is considered in the roughness length parameterization in COARE3.5
 635 (Figure 13b), τ over the misaligned waves has been effectively elevated as the waves
 636 opposing the wind increase the surface drag. This is shown to reduce the low τ bias
 637 significantly.

(a) Location: ($54^\circ\text{W} - 16^\circ\text{N}$) ; $\chi=2.1$

(b) Location: ($46^\circ\text{W} - 6^\circ\text{N}$) ; $\chi=1.1$

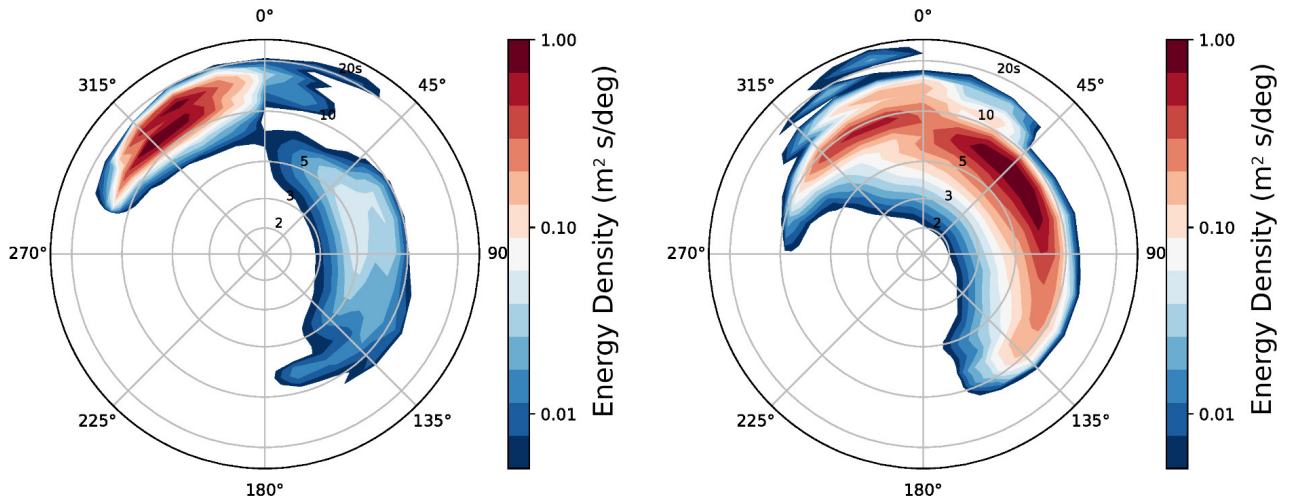


Figure 14. Normalized wave spectrum energy density ($\text{m}^2\text{sdeg}^{-1}$) plotted in period (s) space from (a) one point in the northern part of the domain under swell influence and (b) one point in the center part of the domain on January 08, 2020 at 0600UTC for WBF.

638 Here, the alignment between wind and waves has been defined only by using
 639 the wave peak direction. Figure 14 compares the normalized wave spectrum energy
 640 density ($\text{m}^2\text{sdeg}^{-1}$) shown in the period space between one grid point in the north-
 641 ern part of the domain under swell regime (Fig. 14a) and another grid point in the
 642 center part of the domain under wind waves regime for WBF. Both are sampled
 643 on January 08, 2020 at 0600UTC. On the northern grid point where the wave age
 644 was 2.1, Figure 14a shows the strong swell signal (with the periods of 10-20s)
 645 from the northwest direction. It does also show a large directional spreading,
 646 due to the concurrent shorter period wind waves (2-10s) originating from the
 647 northeast, east, and southeast direction. However, the energy density from the
 648 shorter-period waves is much weaker. In the center of the domain (Figure 14b),
 649 where the sea state is dominated by wind-waves and waves near equilibrium (the
 650 wave age here is 1.1), the directional spreading is also quite large, but with
 651 higher energy in the wind waves and weaker energy in the swell.

652 The sea state in this region appears to be mixed ubiquitously between wind
 653 waves and swell in winter, leading to a large wave directional spreading. How-
 654 ever, since the peak energy density is well separated between the swell (in the
 655 northern point, Fig. 14a) and the wind waves (in the southern point, Fig. 14b),
 656 we anticipate that the use of waves' direction variance in the bulk formula or
 657 the spectrally-averaged wave direction in the bulk formula, would yield qualitatively similar re-

sults. For this reason, in the present study, only the peak direction of the waves is used to account for the misaligned wave effect on z_0 in COARE. However, it is possible that by using the peak wave direction we would grossly underrepresent some unresolved processes contributing to the directional spread of waves, and its impact on z_0 .

6 Conclusion

This study investigated the role of surface waves in surface roughness length (z_0) and surface stress (τ) in the persistent and strong trade winds and swell-dominated Northwestern Tropical Atlantic Ocean during the boreal winter season. The main objective is to evaluate how accurately the air-sea momentum flux is represented in advanced bulk flux algorithms such as COARE3.5 when compared to the direct surface flux measurements. In this investigation, estimated z_0 and τ from four different SCOAR ocean-atmosphere-wave coupled model simulations are analyzed. The results show that the estimated z_0 and τ differences strongly depend on wind speeds and wave age regimes. Wind sea or fully-developed sea under high winds are characterized by the enhanced wave slope and choppy surface (Figure 5b), which effectively increases the surface drag, and τ . The increased surface drag decelerates the near-surface winds (Figure 6c).

However, in the mixed sea condition, where moderate to high wind speeds (10 to 12 ms^{-1}) co-occur with decaying swell, the WBF tends to underestimate z_0 compared to the WSDF and τ compared to the measurements. The weak stress then accelerates the near-surface wind speed by 5% over the region of negative change in wind work (Figure 6d). The sea state, in this high wave age region, is strongly misaligned with the local wind (Figure 5d), indicating the presence of remotely-generated swell. However, despite the swell-dominated sea state, the observations suggest that the wind seas in this mixed sea condition should continue to support the momentum flux due to moderate-to-high wind speeds, thereby increasing τ with wind speed (Figure 7).

The different approaches were explored in this study to alleviate the low-stress bias in the COARE3.5 WBF under the mixed sea regime. The first approach involves re-defining wave age using the mean period of the waves to more accurately represent the wave period in the mixed sea condition (Figure 4a). The second approach takes advantage of the fully coupled model by considering the directionality of waves with respect to winds (Eq. 12), the vital missing process in the current COARE3.5 WBF and many numerical modeling studies except for a limited number of Large Eddy Simulations (LES) and offshore wind energy studies (See Review by Patton et al., 2019). Our results show that both approaches produce equivalent results by effectively boosting z_0 and τ under the misaligned waves under moderate-to-high winds. Since both methods yield equivalent results, accounting for both (peak direction and wave mean period), without more dedicated tuning with the measurements, produces too strong correction for the low bias (not shown). Finally, it is important to note that these improvements are most likely to be site-dependent, as we are only using limited observations in one specific region. Moreover, the improvement of the parameterization is mostly over specific regimes of wind and waves where the original parameterization was deficient.

Our analysis reveals a notable deficiency in the ocean-wave and wave-atmosphere coupling components of the coupled model, which guides the direction of our future investigation. That is, the frequency of swell simulated by the coupled WW3 model is overestimated compared to the in situ observations (Figure 8a), more so with the use of peak wave period but nonetheless noticeable with the use of mean period. Since the wave model provide the parameters required by the WBF, some

709 of the issues described above are a result of inaccurate inputs as well as problems
 710 with the parameterization. The tendency toward the higher wave age indicates that
 711 the model under-represents critical dissipation mechanisms of the swell energy, and
 712 waves in general, which likely have contributed to the low-stress bias. There are at
 713 least two possible factors to consider.

714 First, the primary loss of swell energy is to the atmosphere in situations where
 715 the swell waves outrun the winds or propagate in the opposite direction to the local
 716 wind (e.g., M. Donelan, 1999; Raschle et al., 2008; Kahma et al., 2016; Liu et al.,
 717 2017). Tropical oceans, including our study region, have many low-wind regimes,
 718 where the wave-driven low-level wind jet (Harris, 1966) and turbulent mixing in
 719 the MABL (Kantha, 2006; Ardhuin & Jenkins, 2006; A. V. Babanin, 2006) consti-
 720 tute important sources for attenuation of the swell energy (Ardhuin et al., 2009;
 721 S. Chen et al., 2019). It is quite possible that the processes related to the upward
 722 flux of momentum and energy over swell are not adequately captured in our coupled
 723 wind-wave model. Previous studies find that the wave-driven wind jet is at heights
 724 of 5-10 m (Sullivan et al., 2008; Smedman et al., 2009). However, our experiments
 725 used the default vertical grid system in WRF, where the wind at the lowest height
 726 of the model is typically 30–50 m. The WRF PBL scheme expects this level to be
 727 within the constant-flux layer, where similarity theory is applied (Aligo et al., 2009;
 728 Shin et al., 2012). Yet, this level can be above the surface layer, especially in the
 729 low-wind and stable boundary layer conditions, as often observed in the northern
 730 part of the ATOMIC domain. If the turbulent mixing between the lowest model
 731 level and the swell at the sea surface is weak, the upward energy and momentum
 732 fluxes from the swell to the wind are likely to be under-represented. This might have
 733 been exacerbated by using a local PBL scheme (MYNN) in our model.

734 Moreover, parameterizations for the so-called negative wind input exist in
 735 standalone WW3 model through the use of the source term packages of wind input
 736 (M. A. Donelan et al., 2006; Ardhuin et al., 2010; A. Babanin, 2011; Rogers et al.,
 737 2012; Liu et al., 2017, 2019). With this, the standalone WW3 model forced with
 738 winds should better capture the loss of energy of swell waves. Yet, it is unclear how
 739 such parameterizations should be incorporated into the coupled model, as they do
 740 not represent the actual gain of momentum by the wind from the swell. Our future
 741 work will focus on adequately representing the near-surface wind responses to swell
 742 waves in the atmospheric model.

743 Secondly, the wave breaking and the induced near-surface mixing would in-
 744 fluence the wave energy growth and attenuation (e.g., Kudryavtsev et al., 2014).
 745 Also, Iyer et al. (2022), using the SWIFT drifters deployed during the ATOMIC
 746 campaign, showed that wave-current interactions can generate significant spatial and
 747 temporal variability in momentum fluxes in this region. However, here, since the
 748 current study does not include wave-ocean coupling, the question about the impacts
 749 of ocean-wave coupling on the skill of the simulated wave fields cannot be addressed.
 750 This is a subject of ongoing efforts.

751 **7 Open Research**

752 The observational datasets from the ATOMIC and EUREC⁴A experiments
 753 (Stevens et al., 2021) are available freely on [https://observations.ipsl.fr/
 754 aeris/eurec4a/\#](https://observations.ipsl.fr/aeris/eurec4a/\#/). ERA5 Atmospheric hourly reanalyses were made avail-
 755 able by the Copernicus Climate Change Service (Hersbach et al., 2018a, 2018b).
 756 Mercator Ocean International daily analyses (Lellouche et al., 2018) were
 757 made available by the Copernicus Marine Environment Monitoring Service on
 758 <https://doi.org/10.48670/moi-00016>. Global 3-hourly spectral wave analy-
 759 ses were made available by Ifremer (Raschle & Ardhuin, 2013) on a FTP server at

760 ftp://ftp.ifremer.fr/ifremer/ww3/HINDCAST/GLOBAL; WaveWatchIII model (The
 761 WAVEWATCH III Development Group, 2016) is available at [https://github.com/](https://github.com/NOAA-EMC/WW3)
 762 NOAA-EMC/WW3. WRF model (Skamarock et al., 2008) is available at [https://](https://github.com/wrf-model/WRF)
 763 github.com/wrf-model/WRF. ROMS model (Shchepetkin & McWilliams, 2005) is
 764 also freely available at <https://github.com/kshedstrom/roms>. The SCOAR (Seo
 765 et al., 2007) code is available at <https://github.com/hyodae-seo/SCOAR>. Finally,
 766 the original versions of COARE3.5 (Edson et al., 2013) bulk formula is available at
 767 <https://github.com/NOAA-PSL/COARE-algorithm>.

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778 References

- 779 Aligo, E. A., Gallus, W. A., & Segal, M. (2009). On the impact of WRF model ver-
 780 tical grid resolution on Midwest summer rainfall forecasts. *Wea. Forecasting*,
 781 *24*, 575–594.
- 782 Ardhuin, F., Chapron, B., & Collard, F. (2009). Observation of swell dissipation
 783 across oceans. *Geophys. Res. Lett.*, *36*(6). doi: 10.1029/2008GL037030
- 784 Ardhuin, F., Gille, S. T., Menemenlis, D., Rocha, C. B., Rasche, N., Chapron, B., ...
 785 Molemaker, J. (2017). Small-scale open ocean currents have large effects on
 786 wind wave heights. *J. Geophys. Res.*, *122*, 4500–4517. doi: 10.1002/2016JC012413
- 787 Ardhuin, F., & Jenkins, A. D. (2006). On the Interaction of Surface Waves and
 788 Upper Ocean Turbulence. *Journal of Physical Oceanography*, *36*(3), 551 - 557.
 789 doi: 10.1175/JPO2862.1
- 790 Ardhuin, F., O’Reilly, W. C., Herbers, T. H. C., & Jessen, P. F. (2003). Swell
 791 Transformation across the Continental Shelf. Part I: Attenuation and
 792 Directional Broadening. *J. Phys. Oceanogr.*, *33*(9), 1921–1939. doi:
 793 10.1175/1520-0485(2003)033<1921:STATCS>2.0.CO;2
- 794 Ardhuin, F., Rogers, E., Babanin, A. V., Filipot, J.-F., Magne, R., Roland, A., ...
 795 Collard, F. (2010). Semiempirical Dissipation Source Functions for Ocean
 796 Waves. Part I: Definition, Calibration, and Validation. *J. Phys. Oceanogr.*,
 797 *40*(9), 1917–1941. doi: 10.1175/2010JPO4324.1
- 798 Ardhuin, F., & Roland, A. (2012). Coastal wave reflection, directional spread,
 799 and seismoacoustic noise sources. *J. Geophys. Res.*, *117*(C11). doi:
 800 10.1029/2011JC007832
- 801 Babanin, A. (2011). *Breaking and Dissipation of Ocean Surface Waves*. Cambridge
 802 University Press. doi: 10.1017/CBO9780511736162
- 803 Babanin, A. V. (2006). On a wave-induced turbulence and a wave-mixed upper
 804 ocean layer. *Geophysical Research Letters*, *33*(20). doi: [https://doi.org/10](https://doi.org/10.1029/2006GL027308)
 805 [.1029/2006GL027308](https://doi.org/10.1029/2006GL027308)
- 806 Battjes, J., & Janssen, J. (1978, Jan.). Energy loss and set-up due to breaking of
 807 random waves. *Coastal Engineering Proceedings*, *1*, 32. doi: 10.9753/icce.v16
 808 .32
- 809 Bourras, D., Branger, H., Luneau, C., Reverdin, G., Speich, S., Geykens, N.,
 810 ... Cléménçon, A. (2020). EUREC4A-OA.OCARINA : OCARINA

- 811 Air-Sea Flux Data. *SEANOE*. (Accessed 23 December 2021) doi:
812 <https://doi.org/10.17882/77479>
- 813 Bourras, D., Geyskens, N., Reverdin, G., Cl  men  on, A., Barrois, H., Branger, H., &
814 Luneau, C. (2020). EUREC4A-OA experiment: Air-Sea Flux Mast Data. *SEA-*
815 *NOE*. (Accessed 23 December 2021) doi: <https://doi.org/10.17882/77341>
- 816 B  as, A. B. V., Cornuelle, B. D., Mazloff, M. R., Gille, S. T., & Arduin, F. (2020).
817 Wave-Current Interactions at Meso- and Submesoscales: Insights from Ideal-
818 ized Numerical Simulations. *Journal of Physical Oceanography*, *50*(12), 3483 -
819 3500. doi: 10.1175/JPO-D-20-0151.1
- 820 Charnock, H. (1955). Wind stress on a water surface. *Quart. J. Roy. Meteor. Soc.*,
821 *81*(350), 639-640. doi: 10.1002/qj.49708135027
- 822 Chen, F., & Dudhia, J. (2001). Coupling an Advanced Land Surface-Hydrology
823 Model with the Penn State-NCAR MM5 Modeling System. Part I: Model Im-
824 plementation and Sensitivity. *Monthly Weather Review*, *129*(4), 569 - 585. doi:
825 10.1175/1520-0493(2001)129(0569:CAALSH)2.0.CO;2
- 826 Chen, S., Qiao, F., Jiang, W., Guo, J., & Dai, D. (2019). Impact of Surface Waves
827 on Wind Stress under Low to Moderate Wind Conditions. *Journal of Physical*
828 *Oceanography*, *49*(8), 2017 - 2028. doi: 10.1175/JPO-D-18-0266.1
- 829 Chen, S. S., Zhao, W., Donelan, M. A., & Tolman, H. L. (2013). Directional
830 wind-wave coupling in fully coupled atmosphere-wave-ocean models: Re-
831 sults from cblast-hurricane. *Journal of the Atmospheric Sciences*, *70*(10), 3198
832 - 3215. doi: 10.1175/JAS-D-12-0157.1
- 833 Chou, M.-D., & Suarez, M. J. (1999). A Solar Radiation Parameterization for Atmo-
834 spheric Studies. *Technical Report Series on Global Modeling and Data Assim-*
835 *ilation*, *15*.
- 836 Colosi, L. V., Villas B  as, A. B., & Gille, S. T. (2021). The seasonal cy-
837 cle of significant wave height in the ocean: Local versus remote forcing.
838 *Journal of Geophysical Research: Oceans*, *126*(8), e2021JC017198. doi:
839 <https://doi.org/10.1029/2021JC017198>
- 840 Csanady, G. T., & Gibson, M. (2001). *Air-Sea Interaction: Laws and Mechanisms*.
841 Cambridge University Press. doi: 10.1017/CBO9781139164672
- 842 Deskos, G., Lee, J. C. Y., Draxl, C., & Sprague, M. A. (2021). "review of wind-wave
843 coupling models for large-eddy simulation of the marine atmospheric bound-
844 ary layer". *Journal of the Atmospheric Sciences*, *78*(10), 3025 - 3045. doi:
845 10.1175/JAS-D-21-0003.1
- 846 Donelan, M. (1999). Wind-induced growth and attenuation of laboratory waves.
847 *In: Institute of Mathematics and Its Applications Conference Series*, *69*, Ox-
848 ford;Clarendon, 183-194.
- 849 Donelan, M. A., Babanin, A. V., Young, I. R., & Banner, M. L. (2006). Wave-
850 Follower Field Measurements of the Wind-Input Spectral Function. Part II:
851 Parameterization of the Wind Input. *Journal of Physical Oceanography*, *36*(8),
852 1672 - 1689. doi: 10.1175/JPO2933.1
- 853 Drennan, W. M., Graber, H. C., Hauser, D., & Quentin, C. (2003). On the wave
854 age dependence of wind stress over pure wind seas. *Journal of Geophysical Re-*
855 *search: Oceans*, *108*(C3). doi: <https://doi.org/10.1029/2000JC000715>
- 856 Edson, J., Jampana, V., Weller, R., Bigorre, S., Plueddemann, A., Fairall, C.,
857 ... Hersbach, H. (2013). On the exchange of momentum over the open
858 ocean. *Journal of Physical Oceanography*, *43*(8), 1589 - 1610. doi:
859 10.1175/JPO-D-12-0173.1
- 860 Egbert, G. D., & Erofeeva, S. Y. (2002). Efficient inverse modeling of barotropic
861 ocean tides. *Journal of Atmospheric and Oceanic Technology*, *19*(2), 183 - 204.
862 doi: 10.1175/1520-0426(2002)019(0183:EIMOBO)2.0.CO;2
- 863 Fairall, C. W., Bradley, E. F., Hare, J. E., Grachev, A. A., & Edson, J. B. (2003).
864 Bulk parameterization of air-sea fluxes : Updates and verification for the
865 COARE algorithm. *Journal of Climate*, *16*, 571-591.

- 866 Fairall, C. W., Bradley, E. F., Rogers, D. P., Edson, J. B., & Young, G. S. (1996).
867 Bulk parameterization of air-sea fluxes for tropical ocean-global atmosphere
868 coupled-ocean atmosphere response experiment. *Journal of Geophysical Re-*
869 *search: Oceans*, *101*(C2), 3747-3764. doi: <https://doi.org/10.1029/95JC03205>
- 870 Hanley, K. E., & Belcher, S. E. (2008). Wave-Driven Wind Jets in the Marine At-
871 mospheric Boundary Layer. *Journal of the Atmospheric Sciences*, *65*(8), 2646 -
872 2660. doi: 10.1175/2007JAS2562.1
- 873 Harris, D. L. (1966). The Wave-Driven Wind. *Journal of Atmospheric Sciences*,
874 *23*(6), 688 - 693. doi: 10.1175/1520-0469(1966)023<0688:TWDW>2.0.CO;2
- 875 Hasselmann, S., Hasselmann, K., Allender, J. H., & Barnett, T. P. (1985). Compu-
876 tations and Parameterizations of the Nonlinear Energy Transfer in a Gravity-
877 Wave Spectrum. Part II: Parameterizations of the Nonlinear Energy Transfer
878 for Application in Wave Models. *J. Phys. Oceanogr.*, *15*(11), 1378-1391. doi:
879 10.1175/1520-0485(1985)015<1378:CAPOTN>2.0.CO;2
- 880 Hersbach, H., Bell, B., Berrisford, P., Biavati, G., Horányi, A., Muñoz Sabater, J.,
881 ... Thépaut, J.-N. (2018a). ERA5 hourly data on pressure levels from 1979 to
882 present. Copernicus Climate Change Service (C3S) Climate Data Store (CDS).
883 *[Dataset]*. doi: <https://doi.org/10.24381/cds.bd0915c6>
- 884 Hersbach, H., Bell, B., Berrisford, P., Biavati, G., Horányi, A., Muñoz Sabater, J.,
885 ... Thépaut, J.-N. (2018b). ERA5 hourly data on single levels from 1979 to
886 present. Copernicus Climate Change Service (C3S) Climate Data Store (CDS).
887 *[Dataset]*. doi: <https://doi.org/10.24381/cds.adbb2d47>
- 888 Hong, S.-Y., & Lim, J.-O. J. (2006). The WRF Single-Moment 6-Class Microphysics
889 Scheme (WSM6). *Journal of the Korean Meteorological Society*, *42*, 129-151.
- 890 Iyer, S., Thomson, J., Thompson, E., & Drushka, K. (2022). Variations in wave
891 slope and momentum flux from wave-current interactions in the tropical trade
892 winds. *Journal of Geophysical Research: Oceans*, *127*(3), e2021JC018003. doi:
893 <https://doi.org/10.1029/2021JC018003>
- 894 Janssen, P. A. E. M. (1989). Wave-induced stress and the drag of air flow over sea
895 waves. *Journal of Physical Oceanography*, *19*(6), 745 - 754. doi: 10.1175/1520
896 -0485(1989)019<0745:WISATD>2.0.CO;2
- 897 Janssen, P. A. E. M. (1991). Quasi-linear theory of wind-wave generation applied to
898 wave forecasting. *J. Phys. Oceanogr.*, *21*(11), 1631-1642.
- 899 Jiang, H., & Chen, G. (2013). A Global View on the Swell and Wind Sea Climate
900 by the Jason-1 Mission: A Revisit. *Journal of Atmospheric and Oceanic Tech-*
901 *nology*, *30*(8), 1833 - 1841. doi: 10.1175/JTECH-D-12-00180.1
- 902 Jiménez, P. A., Dudhia, J., González-Rouco, J. F., Navarro, J., Montávez, J. P.,
903 & García-Bustamante, E. (2012). A Revised Scheme for the WRF Sur-
904 face Layer Formulation. *Monthly Weather Review*, *140*(3), 898 - 918. doi:
905 10.1175/MWR-D-11-00056.1
- 906 Kahma, K. K., Donelan, M. A., Drennan, W. M., & Terray, E. A. (2016). Evi-
907 dence of Energy and Momentum Flux from Swell to Wind. *Journal of Physical*
908 *Oceanography*, *46*(7), 2143 - 2156. doi: 10.1175/JPO-D-15-0213.1
- 909 Kantha, L. (2006). A note on the decay rate of swell. *Ocean Modelling*, *11*(1), 167-
910 173. doi: <https://doi.org/10.1016/j.ocemod.2004.12.003>
- 911 Kudryavtsev, V., Chapron, B., & Makin, V. (2014). Impact of wind waves on the
912 air-sea fluxes: A coupled model. *Journal of Geophysical Research: Oceans*,
913 *119*(2), 1217-1236. doi: <https://doi.org/10.1002/2013JC009412>
- 914 Large, W. G., McWilliams, J. C., & Doney, S. C. (1994). Oceanic vertical mixing: A
915 review and a model with a nonlocal boundary layer parameterization. *Reviews*
916 *of Geophysics*, *32*(4), 363-403. doi: <https://doi.org/10.1029/94RG01872>
- 917 Lellouche, J.-M., Greiner, E., Le Galloudec, O., Garric, G., Regnier, C., Drevillon,
918 M., ... Le Traon, P.-Y. (2018). Recent updates to the Copernicus Marine
919 Service global ocean monitoring and forecasting real-time 1/12° high-resolution
920 system. *Ocean Science*, *14*(5), 1093-1126. doi: 10.5194/os-14-1093-2018

- 921 Liu, Q., Babanin, A., Fan, Y., Zieger, S., Guan, C., & Moon, I.-J. (2017). Nu-
 922 merical simulations of ocean surface waves under hurricane conditions: As-
 923 sessment of existing model performance. *Ocean Modelling*, *118*, 73-93. doi:
 924 <https://doi.org/10.1016/j.ocemod.2017.08.005>
- 925 Liu, Q., Rogers, W. E., Babanin, A. V., Young, I. R., Romero, L., Zieger, S., ...
 926 Guan, C. (2019). Observation-based source terms in the third-generation wave
 927 model wavewatch iii: Updates and verification. *Journal of Physical Oceanogra-*
 928 *phy*, *49*(2), 489 - 517. doi: 10.1175/JPO-D-18-0137.1
- 929 Marshall, J., Ferrari, R., Forget, G., Maze, G., Andersson, A., Bates, N., ...
 930 Thomas, L. (2009). The Climode Field Campaign: Observing the Cy-
 931 cle of Convection and Restratification over the Gulf Stream. *Bulletin of*
 932 *the American Meteorological Society*, *90*(9), 1337 - 1350. doi: 10.1175/
 933 2009BAMS2706.1
- 934 Moon, I.-J., Hara, T., Ginis, I., Belcher, S. E., & Tolman, H. L. (2004). Effect of
 935 surface waves on air-sea momentum exchange. part i: Effect of mature and
 936 growing seas. *Journal of the Atmospheric Sciences*, *61*(19), 2321 - 2333. doi:
 937 10.1175/1520-0469(2004)061(2321:EOSWOA)2.0.CO;2
- 938 Nakanishi, M., & Niino, H. (2004). An Improved Mellor-Yamada Level-3 Model
 939 with Condensation Physics: Its Design and Verification. *Boundary-Layer Mete-*
 940 *orology*, *112*, 1-31. doi: 10.1023/B:BOUN.0000020164.04146.98
- 941 Nakanishi, M., & Niino, H. (2006). An Improved Mellor-Yamada Level-3
 942 Model: Its Numerical Stability and Application to a Regional Predic-
 943 tion of Advection Fog. *Boundary-Layer Meteorology*, *119*, 397-407. doi:
 944 10.1007/s10546-005-9030-8
- 945 Nakanishi, M., & Niino, H. (2009). Development of an improved turbulence closure
 946 model for the atmospheric boundary layer. *Journal of the Meteorological Soci-*
 947 *ety of Japan. Ser. II*, *87*(5), 895-912. doi: 10.2151/jmsj.87.895
- 948 Oost, W., Komen, G., Jacobs, C., & Van Oort, C. (2002). New evidence for a
 949 relation between wind stress and wave age from measurements during AS-
 950 GAMAGE. *Boundary-Layer Meteorology volume*, *103*(3), 409 - 438. doi:
 951 10.1023/A:1014913624535
- 952 Patton, E. G., Sullivan, P. P., Kosović, B., Dudhia, J., Mahrt, L., Žagar, M., &
 953 Marić, T. (2019). On the influence of swell propagation angle on surface drag.
 954 *Journal of Applied Meteorology and Climatology*, *58*(5), 1039 - 1059. doi:
 955 10.1175/JAMC-D-18-0211.1
- 956 Phillips, O. M. (1985). Spectral and statistical properties of the equilibrium range in
 957 wind-generated gravity waves. *Journal of Fluid Mechanics*, *156*, 505-531. doi:
 958 10.1017/S0022112085002221
- 959 Porchetta, S., Temel, O., Muñoz Esparza, D., Reuder, J., Monbaliu, J., van Beeck,
 960 J., & van Lipzig, N. (2019). A new roughness length parameterization ac-
 961 counting for wind-wave (mis)alignment. *Atmospheric Chemistry and Physics*,
 962 *19*(10), 6681-6700. doi: 10.5194/acp-19-6681-2019
- 963 Porchetta, S., Temel, O., Warner, J., Muñoz-Esparza, D., Monbaliu, J., van Beeck,
 964 J., & van Lipzig, N. (2021). Evaluation of a roughness length parametrization
 965 accounting for wind-wave alignment in a coupled atmosphere-wave model.
 966 *Quarterly Journal of the Royal Meteorological Society*, *147*(735), 825-846. doi:
 967 <https://doi.org/10.1002/qj.3948>
- 968 Quinn, P. K., Thompson, E. J., Coffman, D. J., Baidar, S., Bariteau, L., Bates,
 969 T. S., ... Zuidema, P. (2021). Measurements from the RV *Ronald H. Brown*
 970 and related platforms as part of the Atlantic Tradewind Ocean-Atmosphere
 971 Mesoscale Interaction Campaign (ATOMIC). *Earth System Science Data*,
 972 *13*(4), 1759-1790. doi: 10.5194/essd-13-1759-2021
- 973 Rasche, N., & Ardhuin, F. (2013). A global wave parameter database for geophysical
 974 applications. Part 2: Model validation with improved source term parame-
 975 terization. *Ocean Modelling*, *70*, 174 - 188. (Ocean Surface Waves) doi:

- 10.1016/j.ocemod.2012.12.001
- 976 Rascle, N., Arduin, F., Queffelec, P., & Croizé-Fillon, D. (2008). A global
 977 wave parameter database for geophysical applications. part 1: Wave-
 978 current–turbulence interaction parameters for the open ocean based on
 979 traditional parameterizations. *Ocean Modelling*, *25*(3), 154–171. doi:
 980 <https://doi.org/10.1016/j.ocemod.2008.07.006>
 981
- 982 Renault, L., McWilliams, J. C., & Penven, P. (2017). Modulation of the Agulhas
 983 Current Retroreflection and Leakage by Oceanic Current Interaction with the
 984 Atmosphere in Coupled Simulations. *Journal of Physical Oceanography*, *47*(8),
 985 2077 - 2100. doi: 10.1175/JPO-D-16-0168.1
- 986 Renault, L., Molemaker, M. J., Gula, J., Masson, S., & McWilliams, J. C. (2016).
 987 Control and Stabilization of the Gulf Stream by Oceanic Current Interaction
 988 with the Atmosphere. *Journal of Physical Oceanography*, *46*(11), 3439 - 3453.
 989 doi: 10.1175/JPO-D-16-0115.1
- 990 Rieder, K. F., Smith, J. A., & Weller, R. A. (1994). Observed directional charac-
 991 teristics of the wind, wind stress, and surface waves on the open ocean. *Jour-
 992 nal of Geophysical Research: Oceans*, *99*(C11), 22589–22596. doi: [https://doi
 993 .org/10.1029/94JC02215](https://doi.org/10.1029/94JC02215)
- 994 Rogers, W. E., Babanin, A. V., & Wang, D. W. (2012). Observation-Consistent
 995 Input and Whitecapping Dissipation in a Model for Wind-Generated Surface
 996 Waves: Description and Simple Calculations. *Journal of Atmospheric and
 997 Oceanic Technology*, *29*(9), 1329 - 1346. doi: 10.1175/JTECH-D-11-00092.1
- 998 Sauvage, C., Lebeaupin Brossier, C., Bouin, M.-N., & Ducrocq, V. (2020). Charac-
 999 terization of the air–sea exchange mechanisms during a Mediterranean heavy
 1000 precipitation event using realistic sea state modelling. *Atmos. Chem. Phys.*,
 1001 *20*(3), 1675–1699. doi: 10.5194/acp-20-1675-2020
- 1002 Semedo, A., Sušelj, K., Rutgersson, A., & Sterl, A. (2011). A Global View on
 1003 the Wind Sea and Swell Climate and Variability from ERA-40. *Journal of
 1004 Climate*, *24*(5), 1461 - 1479. doi: 10.1175/2010JCLI3718.1
- 1005 Seo, H., Miller, A. J., & Roads, J. O. (2007). The Scripps Coupled
 1006 Ocean–Atmosphere Regional (SCOAR) Model, with Applications in the
 1007 Eastern Pacific Sector. *Journal of Climate*, *20*(3), 381–402. doi: 10.1175/
 1008 jcli4016.1
- 1009 Seo, H., Song, H., O’Neill, L. W., Mazloff, M. R., & Cornuelle, B. D. (2021). Im-
 1010 pacts of Ocean Currents on the South Indian Ocean Extratropical Storm Track
 1011 through the Relative Wind Effect. *Journal of Climate*, *34*(22), 9093 - 9113.
 1012 doi: 10.1175/JCLI-D-21-0142.1
- 1013 Seo, H., Subramanian, A. C., Song, H., & Chowdary, J. S. (2019). Coupled effects of
 1014 ocean current on wind stress in the bay of bengal: Eddy energetics and upper
 1015 ocean stratification. *Deep Sea Research Part II: Topical Studies in Oceanogra-
 1016 phy*, *168*, 104617. (Atmosphere-Ocean Dynamics of Bay of Bengal - Volume I)
 1017 doi: <https://doi.org/10.1016/j.dsr2.2019.07.005>
- 1018 Shchepetkin, A. F., & McWilliams, J. C. (2005). The regional oceanic modeling
 1019 system (ROMS): a split-explicit, free-surface, topography-following-coordinate
 1020 oceanic model. *Ocean Modelling*, *9*(4), 347–404. doi: [https://doi.org/10.1016/
 1021 j.ocemod.2004.08.002](https://doi.org/10.1016/j.ocemod.2004.08.002)
- 1022 Shin, H. H., Hong, S.-Y., & Dudhia, J. (2012). Impacts of the Lowest Model
 1023 Level Height on the Performance of Planetary Boundary Layer Parame-
 1024 terizations. *Monthly Weather Review*, *140*(2), 664 - 682. doi: 10.1175/
 1025 MWR-D-11-00027.1
- 1026 Skamarock, W. C., Klemp, J. B., Dudhia, J., Gill, D. O., Barker, D., Duda, M. G.,
 1027 & Powers, J. G. (2008). A Description of the Advanced Research WRF Ver-
 1028 sion 3. (NCAR/TN-475+STR). doi: 10.5065/D68S4MVH
- 1029 Smedman, A., Höögström, U., Sahlé, E., Drennan, W. M., Kahma, K. K., Petters-
 1030 son, H., & Zhang, F. (2009). Observational Study of Marine Atmospheric

- 1031 Boundary Layer Characteristics during Swell. *Journal of the Atmospheric*
1032 *Sciences*, 66(9), 2747 - 2763. doi: 10.1175/2009JAS2952.1
- 1033 Stevens, B., Bony, S., Farrell, D., Ament, F., Blyth, A., Fairall, C., ... Zöger, M.
1034 (2021). EUREC⁴A. *Earth System Science Data*, 13(8), 4067–4119. doi:
1035 10.5194/essd-13-4067-2021
- 1036 Sullivan, P. P., Edson, J. B., Hristov, T., & McWilliams, J. C. (2008). Large-eddy
1037 simulations and observations of atmospheric marine boundary layers above
1038 nonequilibrium surface waves. *Journal of the Atmospheric Sciences*, 65(4),
1039 1225 - 1245. doi: 10.1175/2007JAS2427.1
- 1040 Taylor, P. K., & Yelland, M. J. (2001, 02). The Dependence of Sea Surface Rough-
1041 ness on the Height and Steepness of the Waves. *Journal of Physical Oceanogra-*
1042 *phy*, 31(2), 572-590. doi: 10.1175/1520-0485(2001)031(0572:TDOSSR)2.0.CO;
1043 2
- 1044 The WAVEWATCH III Development Group. (2016). *User manual and system*
1045 *documentation of WAVEWATCH III version 5.16* (Tech. Rep. No. 329).
1046 College Park, MD, USA: NOAA/NWS/NCEP/MMAB. Retrieved from
1047 <https://polar.ncep.noaa.gov/waves/wavewatch/manual.v5.16.pdf>
- 1048 Thompson, E., Fairall, C., Pezoa, S., & Bariteau, L. (2021). ATOMIC ship nav-
1049 igation, meteorology, seawater, fluxes: Near-surface meteorology, air-sea
1050 fluxes, surface ocean waves, and near surface ocean parameters (tempera-
1051 ture, salinity, currents) and primary dataset of ship location and navigation
1052 estimated from in-situ and remote sensing instruments aboard NOAA Ship
1053 Ronald H. Brown in the North Atlantic Ocean, near Barbados: Atlantic
1054 Tradewind Ocean-Atmosphere Mesoscale Interaction Campaign 2020-01-
1055 09 to 2020-02-12 (NCEI Accession 0225427). *NOAA National Centers for*
1056 *Environmental Information. Dataset.* (Accessed 23 December 2021) doi:
1057 <https://doi.org/10.25921/etxb-ht19>
- 1058 Thomson, J. (2012). Wave Breaking Dissipation Observed with “SWIFT” Drifters.
1059 *Journal of Atmospheric and Oceanic Technology*, 29(12), 1866 - 1882. doi: 10
1060 .1175/JTECH-D-12-00018.1
- 1061 Thomson, J., Moulton, M., de Klerk, A., Talbert, J., Guerra, M., Kastner, S.,
1062 ... Nylund, S. (2019). A new version of the SWIFT platform for waves,
1063 currents, and turbulence in the ocean surface layer. In *2019 IEEE/OES*
1064 *Twelfth Current, Waves and Turbulence Measurement (CWTM)* (p. 1-7).
1065 doi: 10.1109/CWTM43797.2019.8955299
- 1066 Thomson, J., Thompson, E., Iyer, S., Drushka, K., & de Klerk, A. (2021). ATOMIC
1067 SWIFT drifters: Near-surface meteorology, air-sea fluxes, surface ocean waves,
1068 and near-surface ocean properties (turbulent dissipation rate, currents, tem-
1069 perature, salinity) estimated from in-situ and remote sensing instruments
1070 aboard six SWIFT drifters (Surface Wave Instrument Float with Tracking)
1071 launched and recovered for two different deployments from NOAA Ship
1072 Ronald H. Brown in the North Atlantic Ocean, near Barbados: Atlantic
1073 Tradewind Ocean-Atmosphere Mesoscale Interaction Campaign 2020-01-
1074 14 to 2020-02-11 (NCEI Accession 0225279). *NOAA National Centers for*
1075 *Environmental Information. Dataset.* (Accessed 23 December 2021) doi:
1076 <https://doi.org/10.25921/s5d7-tc07>
- 1077 Tolman, H. L., Balasubramanian, B., Burroughs, L. D., Chalikov, D. V., Chao,
1078 Y. Y., Chen, H. S., & Gerald, V. M. (2002). Development and Implemen-
1079 tation of Wind-Generated Ocean Surface Wave Modelsat NCEP. *Weather*
1080 *and Forecasting*, 17(2), 311 - 333. doi: 10.1175/1520-0434(2002)017(0311:
1081 DAIOWG)2.0.CO;2
- 1082 Villas Bôas, A. B., & Young, W. R. (2020). Directional diffusion of surface gravity
1083 wave action by ocean macroturbulence. *Journal of Fluid Mechanics*, 890, R3.
1084 doi: 10.1017/jfm.2020.116
- 1085 Wunsch, C. (1998). The Work Done by the Wind on the Oceanic General Circu-

- 1086 lation. *Journal of Physical Oceanography*, 28(11), 2332 - 2340. doi: 10.1175/
1087 1520-0485(1998)028<2332:TWDBTW>2.0.CO;2
- 1088 Zheng, Y., Alapaty, K., Herwehe, J. A., Genio, A. D. D., & Niyogi, D. (2016). Im-
1089 proving High-Resolution Weather Forecasts Using the Weather Research and
1090 Forecasting (WRF) Model with an Updated Kain-Fritsch Scheme. *Monthly*
1091 *Weather Review*, 144(3), 833 - 860. doi: 10.1175/MWR-D-15-0005.1
- 1092 Zou, Z., Song, J., Li, P., Huang, J., Zhang, J. A., Wan, Z., & Li, S. (2019). Effects
1093 of swell waves on atmospheric boundary layer turbulence: A low wind field
1094 study. *Journal of Geophysical Research: Oceans*, 124(8), 5671-5685. doi:
1095 <https://doi.org/10.1029/2019JC015153>